

# **Quantitative stability for Optimal Transport**

**- a complex geometric approach**

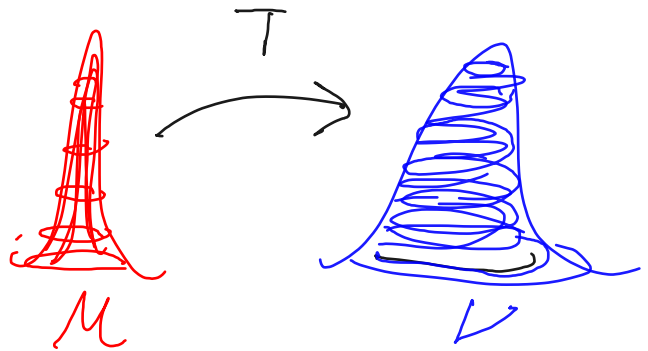
(Inra, 9/2-2022)

Robert Berman

## OT setup

Assume given two probability measures on  $\mathbb{R}^n$   
 $\mu$  and  $\nu$  which have

- compact support
- densities wrt  $dx$



A map  $T$  from  $\mathbb{R}^n$  to  $\mathbb{R}^n$  is called a *transport map* if

$$T_*\mu = \nu, \quad T \in L_{loc}^\infty$$

The *optimal transport map* minimizes

$$\min_T \int_{\mathbb{R}^n} |T(x) - x|^2 \mu =: \mathcal{W}_2(\mu, \nu)^2$$

It is unique and denoted by  $T_{\mu, \nu}$ .

## Stability

The optimal transport map  $T_{\mu,\nu}$  is stable wrt variations of the source  $\mu$  :

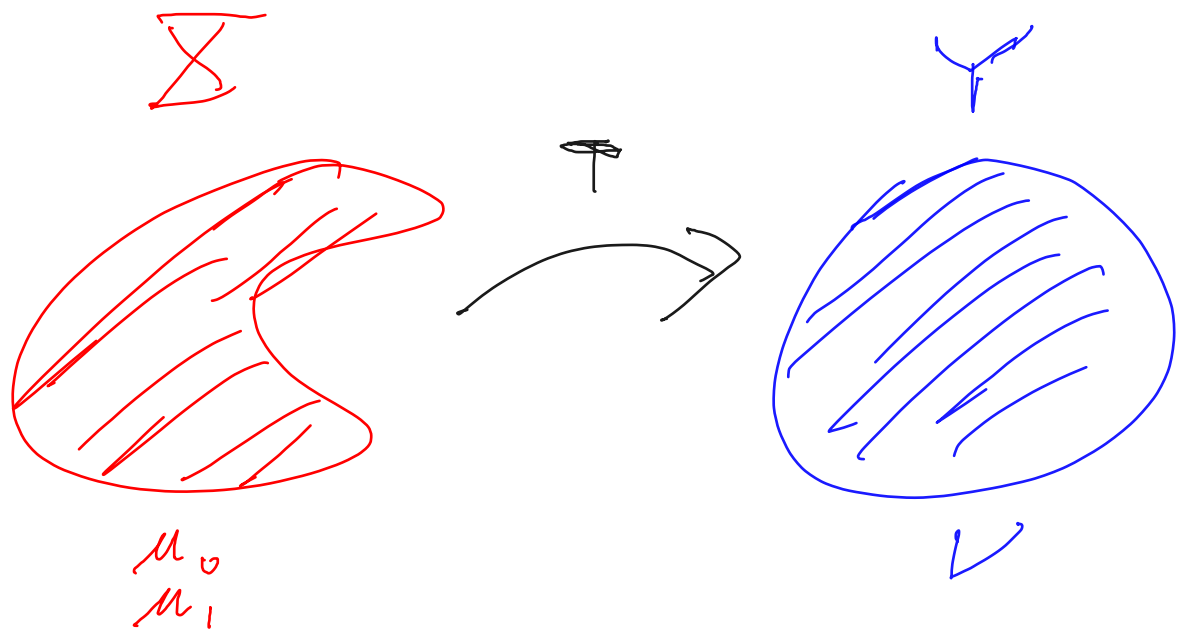
$$\mu_0 - \mu_1 \rightarrow 0 \implies T_{\mu_0,\nu} - T_{\mu_1,\nu} \rightarrow 0$$

(in the weak topology).

Can we make this stability *quantitative* (with a rate):

$$\|T_{\mu_0,\nu} - T_{\mu_1,\nu}\|_{L^2(X,dx)} \preceq \mathcal{W}_1(\mu_0, \mu_1)^\alpha$$

for some  $\alpha > 0$ ?



Throughout the talk we will assume

- “sources”  $\mu_0$  and  $\mu_1$  are supported in the same bounded domain  $X$
- The support of the “target”  $\nu$  is a convex body  $Y$
- $\inf_Y \frac{\nu}{dy} > 0$

## Result (regular case)

There exists a constant  $c_0$  such that

$$\|T_{\mu_0, \nu} - T_{\mu_1, \nu}\|_{L^2(X, dx)} \leq c_0 \mathcal{W}_1(\mu_0, \mu_1)^{1/2}$$

if  $\mu_0$  is “regular” (defined later on).

## The general case

There exists a constant  $c$  such that

$$\|T_{\mu_0, \nu} - T_{\mu_1, \nu}\|_{L^2(X, dx)} \leq c \mathcal{W}_1(\mu_0, \mu_1)^{1/2^n}$$

The constant  $c$  only depends on

- upper bounds on the the *diameters* of  $X$  and  $Y$
- positive lower bounds on the *volume* of  $Y$  and  $\delta(:= \sup_Y(\nu/dy))$

## Comparison with other results

- Previous result by Ambrosio 2011 (but variations wrt the target in a regular setting):
- Subsequent general result by Delalande-Méridot 2021: universal exponent independent of  $n$  (!)(but worse when  $n$  is small)
- ...

## Monge-Ampère formulation

**Recall:** there exists a *convex* function  $\phi$  on  $\mathbb{R}^n$  such that

$$T_{\mu,\nu} = \nabla\phi, \quad \overline{(\nabla\phi)(\mathbb{R}^n)} = Y$$

It is uniquely determined (mod  $\mathbb{R}$ ) by the transport condition:

$$(\nabla\phi)^{-1}\nu = \mu$$

Thus  $\phi$  solves the following *MA-equation* on  $\mathbb{R}^n$ :

$$MA_\nu(\phi) = \mu, \quad \overline{(\nabla\phi)(\mathbb{R}^n)} = Y$$



$$MA_\nu(\phi) = \mu, \quad \overline{(\nabla\phi)(\mathbb{R}^n)} = Y$$

- Now  $\mu$  need *not* be absolute continuous wrt  $dx$ .

- The quantitative stability now becomes

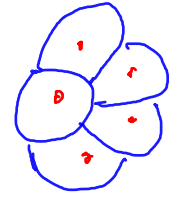
$$\|\nabla\phi_0 - \nabla\phi_1\|_{L^2(X, dx)} \preceq \mathcal{W}_1(\mu_0, \mu_1)^\alpha$$

- The regularity assumption on  $\mu_0$  is:

$$\nabla^2\phi_0 \geq \epsilon I, \quad \epsilon > 0$$

(in the weak sense)

## Application to discretizations



$$MA_\nu(\phi) = \mu, \quad \overline{(\nabla\phi)(\mathbb{R}^n)} = Y$$

Discretize  $\mu$  with a point cloud: *and dual tessellation  $C_i$*

$$\mu_h \rightarrow \mu \quad \mu_h = \sum_{i=1}^N f_i \delta_{x_i} \rightarrow \mu \text{ (weakly)} \quad f_i = \mu(C_i)$$

Denote by  $h$  the “spatial resolution”

$$h := \max_{i < N} \text{diam}(C_i)$$

for the cells  $C_i$  in a dual tessellation of  $X$  and by  $\phi_h$  the corresponding solution. Then

$$\|\nabla\phi - \nabla\phi_h\|_{L^2(X, dx)} \leq h^{1/2}$$

in the “regular case” and in general

$$\|\nabla\phi - \nabla\phi_h\|_{L^2(X, dx)} \leq h^{1/2^n}$$

## The key new analytic inequalities

Henceforth, consider only smooth convex functions  $\phi$  on  $\mathbb{R}^n$  such that

$$\overline{(\nabla\phi)(\mathbb{R}^n)} = Y \quad (:= \text{support of } \nu)$$

In the “regular case”,  $\nabla^2 \phi_0 \geq \epsilon I$ ,

$$\int_X |\nabla \phi_0 - \nabla \phi_1|^2 dx \leq c_0 \int (\phi_1 - \phi_0) (MA(\phi_0) - MA(\phi_1))$$

In general, need to raise the lhs to  $2^{n-1}$ .

- In the rhs can integrate over  $X$  or  $\mathbb{R}^n$
- OT stability follows from

$$\left| \int_{\mathbb{R}^n} (\phi_1 - \phi_0) (\mu_0 - \mu_1) \right| \leq \text{diam}(Y) \mathcal{W}_1(\mu_0, \mu_1)$$

using Kantorovich duality formula for  $\mathcal{W}_1(\mu_0, \mu_1)$

$$\|\nabla \phi_i\|_{L^\infty} \leq C$$

## Proof using complex geometry

Why?

- Can leverage the exterior algebra of complex forms ( = multi-linearity )
- Can “compactify”  $\mathbb{R}^n$  to avoid boundary terms in integration by parts

## Starting point

It is enough to consider the case when

$$\nu = 1_Y dx$$

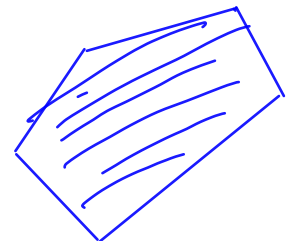
(not explained today...). Then

$$MA_\nu(\phi) = \det(\nabla^2 \phi) dx$$

Moreover, may assume:

Y is a convex polytope with rational vertices

(by approximation)



**Warm-up:**  $n = 1$

$$Y = [0, 1], \quad MA(\phi) = (\nabla^2 \phi) dx.$$

In this case the key inequality is an *equality*:

$$\int_{\mathbb{R}} |\nabla \phi_0 - \nabla \phi_1|^2 dx = \int_{\mathbb{R}} (\phi_1 - \phi_0) (\nabla^2 \phi_0 - \nabla^2 \phi_1) dx.$$

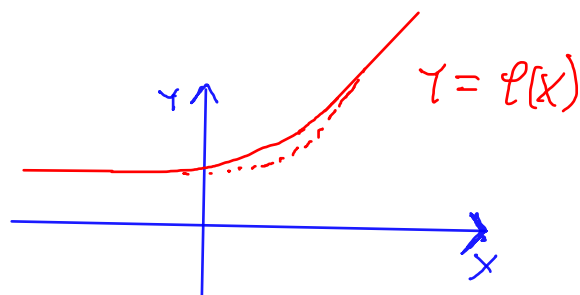
Indeed, integration by parts gives “boundary terms”:

$$(\phi_1 - \phi_0) (\nabla \phi_0 - \nabla \phi_1) (\pm\infty).$$

But, by assumption,

$$\overline{(\nabla \phi)(\mathbb{R}^n)} = [0, 1] \implies (\nabla \phi_0 - \nabla \phi_1) (\pm\infty) = 0$$

and  $(\phi_1 - \phi_0) = O(1)$ .



# The “regular case” when $n > 1$

We will focus on the “regular case”  $\nabla^2\phi_0 \geq \epsilon I$  and show

$$\|\nabla\phi_0 - \nabla\phi_1\|_{L^2(X, dx)}^2 \leq c_0 \int (\phi_0 - \phi_1) (MA(\phi_1) - MA(\phi_0))$$

Start with the rhs.

- Will leverage complex exterior algebra:

$$MA(\phi) := \det(\nabla^2\phi) dx \sim (\partial\bar{\partial}\phi)^n.$$

“Hessian”

- Will first proceed formally
- To simplify notation will consider  $\mathbb{R}^2$



# Caveat

To simplify the notation will leave out all factors of  $2, \pi$  etc!

$$a^2 - b^2 = (a-b)(a+b)$$

$$\begin{aligned} \parallel MA(\phi_1) - MA(\phi_0) &= (\partial\bar{\partial}\phi_1)^2 - (\partial\bar{\partial}\phi_0)^2 = \\ &= (\partial\bar{\partial}\phi_1 - \partial\bar{\partial}\phi_0) \wedge (\partial\bar{\partial}\phi_1 + \partial\bar{\partial}\phi_0). \end{aligned}$$

Hence,

$$\begin{aligned} \parallel \int_{\mathbb{R}^2} (\phi_0 - \phi_1) (MA(\phi_1) - MA(\phi_0)) &= \\ \int (\phi_0 - \phi_1) (\partial\bar{\partial}\phi_1 - \partial\bar{\partial}\phi_0) \wedge (\partial\bar{\partial}\phi_1 + \partial\bar{\partial}\phi_0) &= \\ \int \partial(\phi_0 - \phi_1) \wedge \bar{\partial}(\phi_0 - \phi_1) \wedge (\cancel{\partial\bar{\partial}\phi_1} + \partial\bar{\partial}\phi_0) &\geq \\ \int \partial(\phi_0 - \phi_1) \wedge \bar{\partial}(\phi_0 - \phi_1) \wedge \partial\bar{\partial}\phi_0 &\geq \\ \epsilon \|\nabla(\phi_0 - \phi_1)\|_{L^2(X, dx)}^2 & \end{aligned}$$

## Details on complexification

$$z := x + iy \in \mathbb{C}^n := \mathbb{R}^n + i\mathbb{R}^n$$

$$\bar{z} = x - iy$$

Decompose real exterior derivative  $d$  wrt  $dz_i$  and  $d\bar{z}_i$ :

$$d = \partial + \bar{\partial},$$

on  $\mathbb{R}^n \times \mathbb{R}^n$

$$\partial := \sum_{i=1}^n \frac{\partial}{\partial z_i} dz_i, \quad \bar{\partial} := \sum_{i=1}^n \frac{\partial}{\partial \bar{z}_i} d\bar{z}_i$$

Hence,

$$\partial \bar{\partial} \phi(z) = \sum_{i,j=1}^n \frac{\partial^2 \phi}{\partial z_j \partial \bar{z}_i} dz_j \wedge d\bar{z}_i$$

In particular, if  $\phi = \phi(x)$ , then

$$\partial \bar{\partial} \phi = \sum_{i,j=1}^n \frac{\partial^2 \phi}{\partial x_j \partial x_i} dz_j \wedge d\bar{z}_i$$

$$\left( \frac{\partial^2 \phi}{\partial x^2} \right)$$

As a consequence,

$$\frac{\partial^2 \phi}{\partial x^2} = 0$$

$$\partial\bar{\partial}\phi = \sum_{i=1}^n \frac{\partial^2\phi}{\partial x_j \partial \bar{x}_j} dz_i \wedge d\bar{z}_j \implies$$

$$\underline{(\partial\bar{\partial}\phi)^n} = \det \left( \frac{\partial^2\phi}{\partial x_j \partial \bar{x}_j} \right) \underline{dx} \wedge \underline{dy} = \underline{MA(\phi)} \wedge \underline{dy} \quad \mathbb{C}^n$$

Moreover, if  $u = u(x)$  and  $\phi_0 = |x|^2/2$ , then

$$\underline{\partial u \wedge \bar{\partial} u} \wedge (\partial\bar{\partial}\phi_0)^{n-1} = \underline{(|\nabla u|^2 dx)} \wedge \underline{dy}$$

$$u = \ell_0 - \ell_1$$

This justifies the formal proof, except the integration by parts.

It is handled by a compactification argument:

$$\int_{\mathbb{R}^n} (\phi_0 - \phi_1) (MA(\phi_1) - MA(\phi_0)) = \\ = \int_{X_Y} (\phi_0 - \phi_1) \left( (\partial\bar{\partial}\phi_1)^n - (\partial\bar{\partial}\phi_0)^n \right)$$

on a certain *compact* complex manifold  $X_Y$  (without boundary!) such that

$(\partial\bar{\partial}\phi_i)$  and  $\phi_0 - \phi_1$   
extend to  $X_Y$ .

## Compactification

**Recall:** starting with  $x \in \mathbb{R}^n$ , we consider  $z = x + iy \in \mathbb{C}^n$ .

**Step 1:** “compactify” the  $y$ -direction by replacing

$$i\mathbb{R}^n \rightsquigarrow i(\mathbb{R}/2\pi\mathbb{Z})^n.$$

In complex terms: set

$$z_i = \log w_i \in \mathbb{C}, \quad w \in (\mathbb{C}^*)^n$$

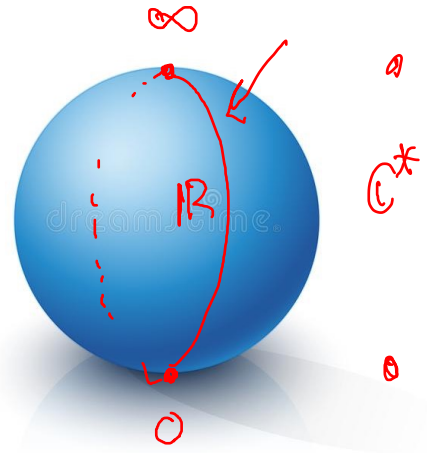
Hence,

$$x_i = \log |w_i|$$

and

$$\int_{\mathbb{R}^n} (\phi_0 - \phi_1) (MA(\phi_1) - MA(\phi_0)) = \int_{(\mathbb{C}^*)^n} (\phi_0 - \phi_1) \left( (\partial\bar{\partial}\phi_1)^n - (\partial\bar{\partial}\phi_0)^n \right)$$

using that  $\int_{\mathbb{R}/2\pi\mathbb{Z}} dy = 1/(2\pi)^n$ .



**Step 2:** compactify  $(\mathbb{C}^*)^n$

**Motivation** ( $n = 1$ ):

$$\mathbb{C}^* \hookrightarrow \mathbb{P}_{\mathbb{C}}^1 (:= \mathbb{C}^* \cup \{0, \infty\}),$$

where  $\mathbb{P}_{\mathbb{C}}^1$  is a complex compact manifold (the Riemann sphere).

If

$$(\nabla \phi)(\mathbb{R}) = [0, 1],$$

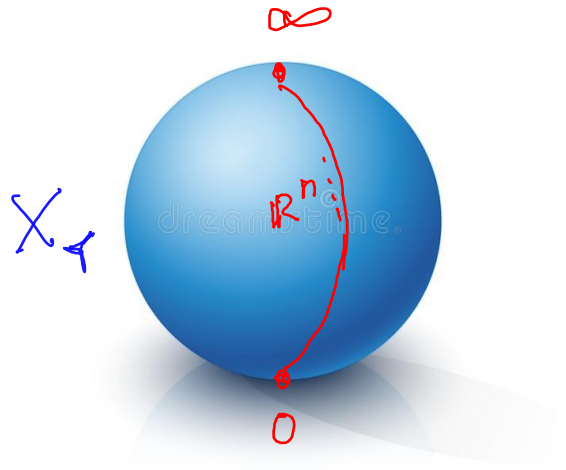
then

$$1 = \int_{\mathbb{R}} \nabla^2 \phi dx = \frac{i}{2\pi} \int_{\mathbb{C}^*} \partial \bar{\partial} \phi = \frac{i}{2\pi} \int_{\mathbb{P}_{\mathbb{C}}^1} \partial \bar{\partial} \phi$$

Moreover,

$$\int_{\mathbb{R}} |\nabla(\phi_1 - \phi_0)|^2 dx = \frac{i}{2\pi} \int_{\mathbb{P}_{\mathbb{C}}^1} \partial(\phi_1 - \phi_0) \wedge \bar{\partial}(\phi_1 - \phi_0)$$





In higher dimensions, if

$$(\nabla \phi)(\mathbb{R}^n) = Y$$

for a convex bounded rational *polytope*  $Y$ . Then

$$(\mathbb{C}^*)^n \hookrightarrow X_Y \text{ (dense image)}$$

for a compact complex manifold (without boundary!) such that

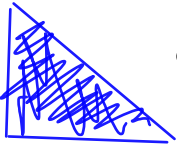
$$\partial \bar{\partial} \phi \text{ extends from } (\mathbb{C}^*)^n \text{ to } X_Y$$

and likewise for differences  $\phi_0 - \phi_1$

**Ex:**



- for  $Y = [0, 1]^2 \subset \mathbb{R}^2$  get  $X_Y = \mathbb{P}_{\mathbb{C}}^1 \times \mathbb{P}_{\mathbb{C}}^1$ .



- for  $Y$  a simplex in  $\mathbb{R}^n$  get  $X_Y = \mathbb{P}_{\mathbb{C}}^n$

In general,  $X_Y$  is a toric variety, i.e. an equivariant compactification of the complex torus  $\mathbb{C}^{*n}$ .

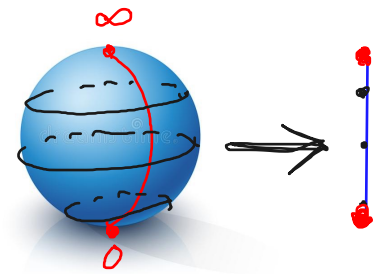
This proves the regular case.

Questions?

General case from Block's ineq.  
(Cauchy-Schwarz)

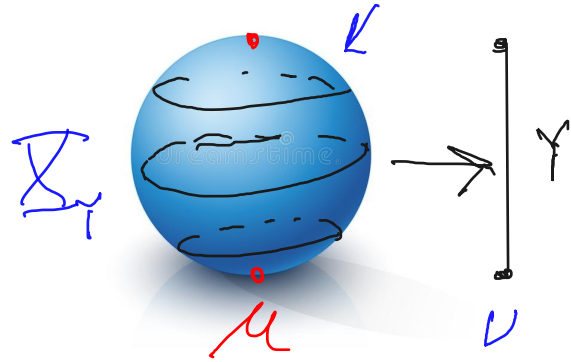
## Symplectic interpretation

$$\begin{array}{ccccc}
 \mathbb{T} & \rightarrow & X_Y & \twoheadrightarrow & Y \\
 & & \uparrow & & \uparrow \\
 \nabla\phi : & & \mathbb{R}^n & \twoheadrightarrow & \mathring{Y}
 \end{array}$$



- $\dim X_Y = 2n$
- The fibers of  $X_Y$  over  $\mathring{Y}$  are  $n$ -dimensional tori  $T^n$
- Some of the torus-dimensions shrink over the boundary of  $Y$ .

$$\omega = i \partial \bar{\partial} \phi$$



- The torus  $T^n$  acts on  $X_Y$
- Can identify  $\mu$  with a  $T^n$ -invariant measure on  $X_Y$
- Finding  $\nabla \phi$  amounts to finding a symplectic two-form  $\omega$  on  $X_Y$  such that

$$\omega^n / n! = \mu.$$

- Then  $\nabla \phi$  lifts to the moment map for the  $T^n$ -action on  $(X, \omega)$ :

$$X_Y \rightarrow Y \subset \text{Lie}(T^n)^* = \mathbb{R}^n \quad (\text{action/angle coord.})$$

- The moment map pushes forward  $\left(\frac{\omega^n}{n!}\right)$  to  $dy$  on  $Y$ :

$$(\nabla \phi)_* \left( \frac{\omega^n}{n!} \right) = (\nabla \phi)_* \mu = 1_Y dy,$$

$$\omega = i \partial \bar{\partial} \phi$$

- In complex notation:

$$\omega = i\partial\bar{\partial}\phi$$