# Approximations of displacement interpolations by entropic interpolations

Christian Léonard

Université Paris Ouest

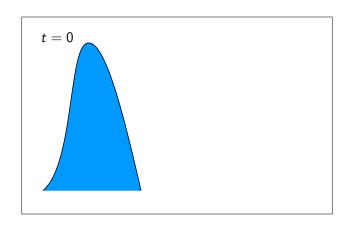
Mokaplan. 10 décembre 2015

- ullet  ${\cal X}$  : Riemannian manifold (state space)
- ullet  $\mathcal{P}(\mathcal{X})$  : set of all probability measures on  $\mathcal{X}$
- $\mu_0, \mu_1 \in \mathcal{P}(\mathcal{X})$
- ullet interpolate between  $\mu_0$  and  $\mu_1$

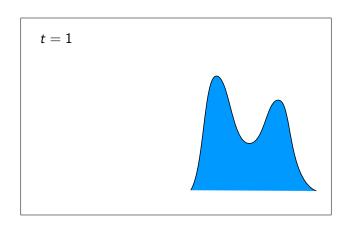
• Standard affine interpolation between 
$$\mu_0$$
 and  $\mu_1$   $\mu_t^{\mathrm{aff}} := (1-t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \leq t \leq 1$ 



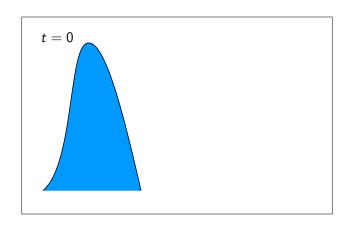
• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\mathrm{aff}} := (1-t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



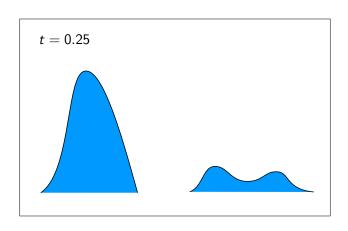
• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\text{aff}} := (1 - t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



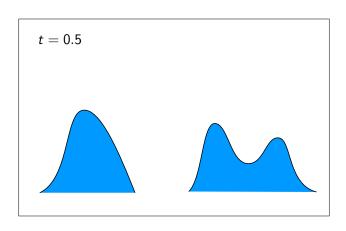
• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\mathrm{aff}} := (1-t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



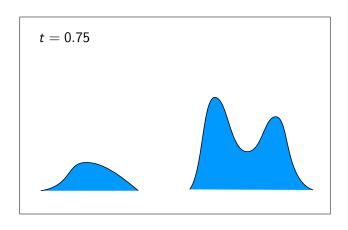
• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\text{aff}} := (1 - t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



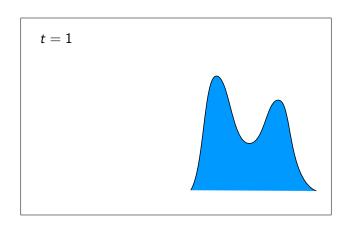
• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\text{aff}} := (1 - t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\rm aff} := (1-t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 



• Standard affine interpolation between  $\mu_0$  and  $\mu_1$   $\mu_t^{\text{aff}} := (1 - t)\mu_0 + t\mu_1 \in \mathcal{P}(\mathcal{X}), \ 0 \le t \le 1$ 

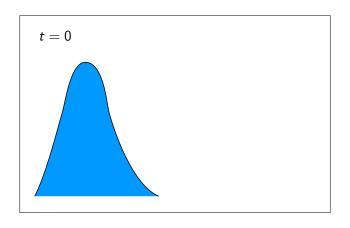


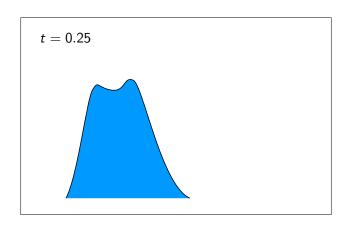
Affine interpolations require mass transference with infinite speed

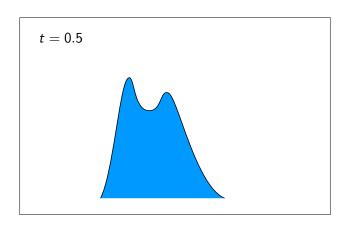


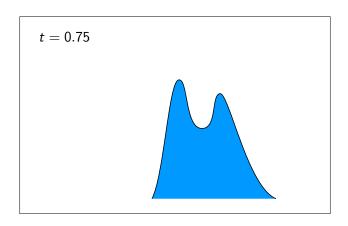
- ullet Denial of the geometry of  ${\mathcal X}$
- We need interpolations built upon trans-portation, not tele-portation

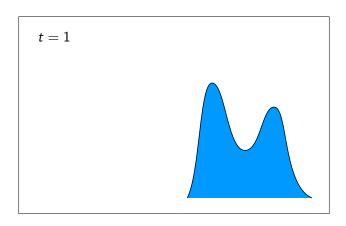
• We seek interpolations of this type	

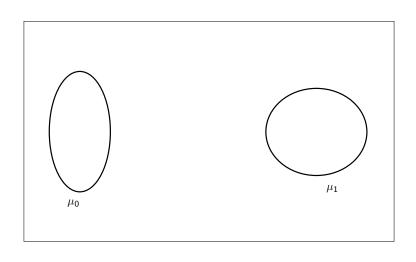


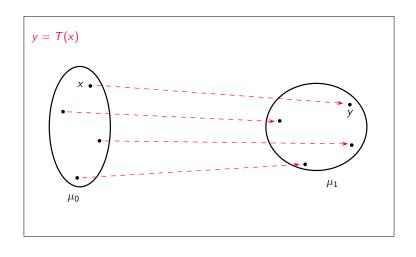


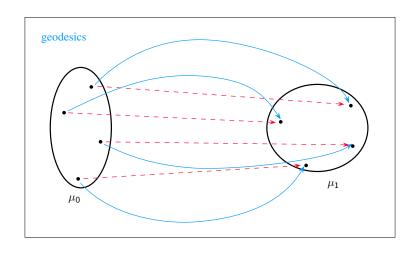


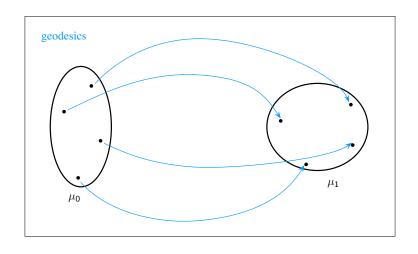


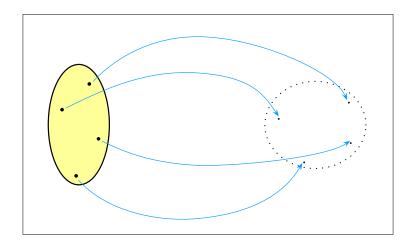


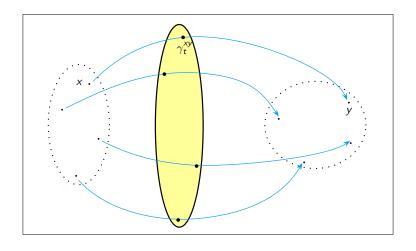


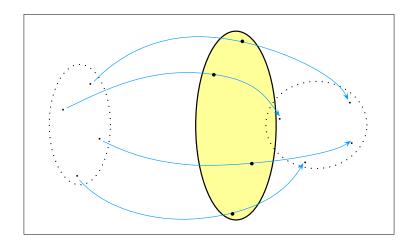


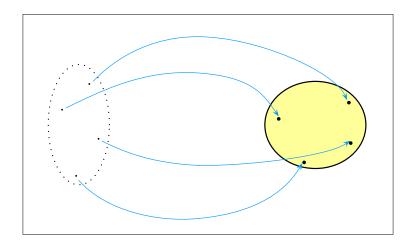






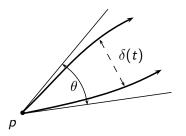




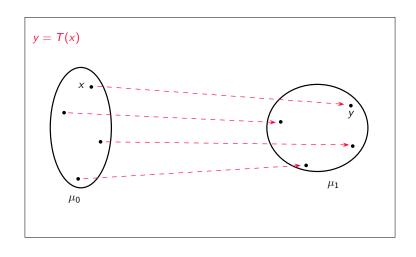


#### Curvature

- geodesics and curvature are intimately linked
- several geodesics give information on the curvature



$$\delta(t) = \sqrt{2(1-\cos\theta)} \ t \left(1 - \frac{\sigma_{p}(\mathsf{S})\cos^{2}(\theta/2)}{6}t^{2} + O(t^{4})\right)$$



#### Respect geometry

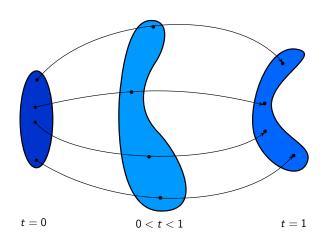
- we have already used geodesics
- how to choose y = T(x) such that interpolations encrypt curvature as best as possible?
- no shock
- perform optimal transport

#### Monge's problem

$$\int_{\mathcal{X}} d^2(x, T(x)) \, \mu_0(dx) \mapsto \min; \qquad T: \quad T_\# \mu_0 = \mu_1$$

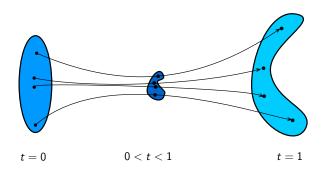
• d : Riemannian distance

# Lazy gas experiment



Positive curvature

# Lazy gas experiment



Negative curvature

# Curvature and displacement interpolations

#### Relative entropy

$$H(p|r) := \int \log(dp/dr) dp$$
,  $p, r$ : probability measures

#### Convexity of the entropy along displacement interpolations

The following assertions are equivalent

- Ric  $\geq K$
- along any  $[\mu_0, \mu_1]^{\mathrm{disp}} = (\mu_t)_{0 \le t \le 1}, \; \frac{d^2}{dt^2} H(\mu_t | \mathrm{vol}) \ge K W_2^2(\mu_0, \mu_1)$
- von Renesse-Sturm (04)
- W<sub>2</sub> is the Wasserstein distance
- starting point of the Lott-Sturm-Villani theory

# Schrödinger's thought experiment

Consider a huge collection of non-interacting identical Brownian particles. If the density profile of the system at time t=0 is approximately  $\mu_0\in\mathcal{P}(\mathbb{R}^3)$ , you expect it to evolve along the heat flow:

$$\begin{cases} \nu_t = \nu_0 e^{t\Delta/2}, & 0 \le t \le 1 \\ \nu_0 = \mu_0 \end{cases}$$

where  $\Delta$  is the Laplace operator.

Suppose that you observe the density profile of the system at time t=1 to be approximately  $\mu_1 \in \mathcal{P}(\mathbb{R}^3)$  with  $\mu_1$  different from the expected  $\nu_1$ . Probability of this rare event  $\simeq \exp(-\mathit{CN}_{\mathrm{Avogadro}})$ .

#### Schrödinger's question (1931)

Conditionally on this very rare event, what is the *most likely path*  $(\mu_t)_{0 \le t \le 1} \in \mathcal{P}(\mathbb{R}^3)^{[0,1]}$  of the evolving profile of the particle system?

# Schrödinger's problem

- $\bullet$   $\mathcal{X}$  : Riemannian manifold
- $\Omega := \{ \text{paths} \} \subset \mathcal{X}^{[0,1]}$
- $P \in \mathcal{P}(\Omega)$  and  $(P_t)_{0 \leq t \leq 1} \in \mathcal{P}(\mathcal{X})^{[0,1]}$
- ullet  $R\in \mathcal{P}(\Omega)$  : Wiener measure (Brownian motion)

#### Schrödinger's problem

$$H(P|R) \to \min; \quad P \in \mathcal{P}(\Omega) : P_0 = \mu_0, P_1 = \mu_1$$
 (S<sub>dyn</sub>)

 $\mu_0, \mu_1 \in \mathcal{P}(\mathcal{X})$  are the initial and final prescribed profiles

### Definition. *R*-entropic interpolation

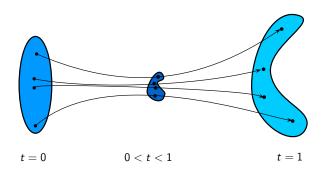
$$[\mu_0, \mu_1]^R := (P_t)_{0 \le t \le 1}$$
 with  $P$  the unique solution of  $(S_{\mathrm{dyn}})$ .

It is the answer to Schrödinger's question

#### Lazy gas experiments

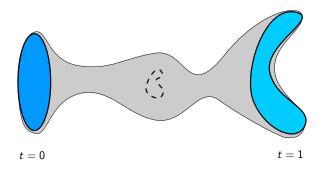
- Lazy gas experiment at zero temperature (Monge)
  - Zero temperature
  - Displacement interpolations
  - Optimal transport
- Lazy gas experiment at positive temperature (Schrödinger)
  - Positive temperature
  - Entropic interpolations
  - Minimal entropy

# Lazy gas experiments



Negative curvature Zero temperature

# Lazy gas experiments



Negative curvature Positive temperature

## Cooling down

#### Aim

Drifting from Schrödinger problem to an optimal transport problem

#### To decrease temperature:

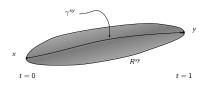
- slow down the particles of the heat bath
- more generally, decrease fluctuations

#### Slowed down reference measures

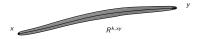
- ullet  $(B_t)_{t\geq 0}$  : Brownian motion on the Riemannian manifold  ${\mathcal X}$
- R: law of  $(B_t)_{0 \le t \le 1}$
- $R^k$ : law of  $(B_{t/k})_{0 \le t \le 1}$
- $k \to \infty$

# Cooling down

k = 1:



k = 10:



 $k = \infty$ :



## Cooling down

- $N \to \infty, \ k=1$  : the whole particle system performs a rare event to travel from  $\mu_0$  to  $\mu_1$ 
  - cooperative behavior
  - lacktriangleright Gibbs conditioning principle (thermodynamical limit:  $N o\infty$ )
- $N=1,\ k\to\infty$  : each individual particle faces a hard task and must travel along an approximate geodesic
  - individual behavior
  - ▶ large deviation principle (cooling down limit:  $k \to \infty$ )

### Cooling down principle

The cooled down sequence  $(R^k)_{k\geq 1}$  encodes some geometry

•  $N \to \infty, k \to \infty$ : these two behaviors superpose

### Results

#### Results 1

- displacement interpolations feel curvature
- entropic interpolations also feel curvature

#### Results 2

- entropic interpolations converge to displacement interpolations
- entropic interpolations regularize displacement interpolations

F-convergence

### Results

### Results 2 (continued)

The same kind of results hold in other settings

- (a) discrete graphs
- (b) Finsler manifolds
- (c) interpolations with varying mass
- (a) graphs: random walk
- (b) Finsler: jump process in a manifold, (work in progress)
- (c) varying mass: branching process, (work in progress)

### Results

#### Results 3

Schrödinger's problem is an analogue of Hamilton's least action principle. It allows for *dynamical* theories of

- diffusion processes
- random walks on graphs
- stochastic Newton equation
- acceleration is related to curvature

### Remainder of the talk

Let us give some details about Results 2:

entropic interpolations converge to displacement interpolations

### Notation

- $\mathcal{X} = \{\text{states}\}$
- $\Omega = \{ \text{paths} \} \subset \mathcal{X}^{[0,1]}, \quad \omega = (\omega_t)_{0 \le t \le 1} \in \Omega$
- $\bullet \ \, \textit{$X_t:\omega\in\Omega\mapsto\omega_t\in\mathcal{X},$} \ \, 0\leq t\leq 1 \qquad \text{(canonical process)}$
- $P \in \mathcal{P}(\Omega)$
- $P_t(dz) := [(X_t)_{\#}P](dz) = P(X_t \in dz) \in \mathcal{P}(\mathcal{X})$

## Particle system

- $\bullet \ \Omega = \{\mathrm{paths}\} \subset \mathcal{X}^{[0,1]}$
- $R \in \mathcal{P}(\Omega)$ : reference Markov measure
- $\bullet$   $R^x := R(\cdot \mid X_0 = x), x \in \mathcal{X}$
- $(Z^i)_{1 \le i \le N} \in \Omega^N$ , independent dynamical particles  $\text{Law}(Z^i) = R^{x^i}, \quad 1 \le i \le N$
- $\bullet$   $N \to \infty$

#### Assume that

$$L_0^N = \frac{1}{N} \sum_{1 \le i \le N}^N \delta_{x^i} \underset{N \to \infty}{\longrightarrow} \mu_0 \in \mathcal{P}(\mathcal{X}), \quad (t=0)$$

### **Empirical measures**

$$L^{N} := \frac{1}{N} \sum_{1 \leq i \leq N}^{N} \delta_{Z^{i}} \in \mathcal{P}(\Omega)$$
  
$$L_{t}^{N} := \frac{1}{N} \sum_{1 \leq i \leq N}^{N} \delta_{Z^{i}} \in \mathcal{P}(\mathcal{X}), \quad 0 \leq t \leq 1$$

### Law of large numbers

- $L^N(d\omega) \underset{N \to \infty}{\rightarrow} R^{\mu_0}(d\omega) := \int_{\mathcal{X}} R^{\kappa}(d\omega) \, \mu_0(d\kappa) \in \mathcal{P}(\Omega)$
- $\bullet \ (L_t^N)_{0 \le t \le 1} \underset{N \to \infty}{\to} (R_t^{\mu_0})_{0 \le t \le 1} \quad \in \mathcal{P}(\mathcal{X})^{[0,1]}$
- $L_1^N(dy) \underset{N \to \infty}{\to} R_1^{\mu_0}(dy) \in \mathcal{P}(\mathcal{X})$

### Schrödinger's question (1931)

- N large
- suppose that you observe:  $L_1^{\it N} \simeq \mu_1,\, \mu_1 
  eq R_1^{\mu_0} \quad (t=1,\, {\sf rare \, event})$
- question: "conditionally on this rare event, what is the most likely path  $(L_t^N)_{0 \le t \le 1}$ ?"

Sanov's theorem:

$$\Pr(L^N \in A) \underset{N \to \infty}{\asymp} \exp\Big(-N \inf_{P \in A, P_0 = \mu_0} H(P|R)\Big), \quad A \subset \mathcal{P}(\Omega)$$

### Relative entropy

$$H(p|r) := \int \log (dp/dr) \ dp \in [0,\infty]$$

• 
$$\Pr(L^N \in A \mid L_1^N \simeq \mu_1) \underset{N \to \infty}{\simeq}$$
  
 $\exp\left(-N\left\{\inf_{P \in A, P_0 = \mu_0, P_1 = \mu_1} H(P|R) - \inf_{P_0 = \mu_0, P_1 = \mu_1} H(P|R)\right\}\right)$ 

• answer to Schrödinger's question: the most likely path is close to the time-marginal flow  $(P_t)_{0 \le t \le 1}$  of the unique solution P of

### Dynamical Schrödinger problem

$$H(P|R) o \min; \qquad P \in \mathcal{P}(\Omega) : P_0 = \mu_0, P_1 = \mu_1$$
 (S<sub>dyn</sub>)

H. Föllmer

### Definition (entropic interpolation)

Let P be the solution of  $(S_{dyn})$ . Then,

$$\mu_t := P_t, \quad 0 \le t \le 1$$

is the *R*-entropic interpolation between  $\mu_0$  and  $\mu_1$  in  $\mathcal{P}(\mathcal{X})$ .

- $P_{01} = (X_0, X_1)_{\#} P \in \mathcal{P}(\mathcal{X}^2)$
- $P^{xy} = P(\cdot \mid X_0 = x, X_1 = y) \in \mathcal{P}(\Omega)$ : bridge
- $P(\cdot) = \int_{\mathcal{X}^2} P^{xy}(\cdot) P_{01}(dxdy) \in \mathcal{P}(\Omega)$

#### Result

If it exists, the unique solution P of  $(S_{\mathrm{dyn}})$  satisfies

- $\bullet P^{xy} = R^{xy}, \forall x, y$
- $P(\cdot) = \int_{\mathcal{X}^2} R^{xy}(\cdot) \pi(dxdy)$

where  $P_{01} = \pi \in \mathcal{P}(\mathcal{X}^2)$  is the unique solution of (S) below

•  $\inf(S) = H(P|R) = H(P_{01}|R_{01})$ 

### Schrödinger's problem

$$H(\pi|R_{01}) \to \min; \quad \pi \in \mathcal{P}(\mathcal{X}^2) : \pi_0 = \mu_0, \pi_1 = \mu_1$$
 (S)

• 
$$H(P|R) = H(P_{01}|R_{01}) + \int_{\mathcal{X}^2} H(P^{xy}|R^{xy}) P_{01}(dxdy)$$

#### Result

Assume: R is m-stationary Markov,  $m \otimes m \ll R_{01} \ll m \otimes m$ ,  $H(\mu_0|m), H(\mu_1|m) < \infty, \dots$ 

Then,  $(S_{dyn})$  and (S) admit a solution.

- long history: Schrödinger, Bernstein, Fortet, Beurling, Csiszár,
   Rüschendorf & Thomsen, Föllmer & Gantert, L.
- and also: Jamison, Zambrini, Dai Pra, Wakolbinger, Pavon,
   Mikami, Rœlly, Thieullen, . . .

- suppose that in addition the heat bath is pretty cold
- cooling down is (mostly) slowing down the particles

$$\bullet \ R = \operatorname{Law}((B_t)_{0 \le t \le 1}), \quad R^k := \operatorname{Law}((B_{k^{-1}t})_{0 \le t \le 1}), \quad k \to \infty$$

$$H(P|R^{k})/k \to \min; \qquad P \in \mathcal{P}(\Omega) : P_0 = \mu_0, P_1 = \mu_1 \qquad (S_{\text{dyn}}^{k})$$

$$H(\pi|R_{01}^{k})/k \to \min; \quad \pi \in \mathcal{P}(\mathcal{X}^{2}) : \pi_{0} = \mu_{0}, \pi_{1} = \mu_{1}$$
 (S<sup>k</sup>)

- $dX_t = \sqrt{1/k} dB_t$ ,  $R^k$ -a.s.
- $R_{01}^k(dxdy) = (2\pi/k)^{-d/2} \exp(-k|y-x|^2/2) dxdy$ • with  $P^k$  solution of  $(S_{dyn}^k)$

$$\inf(S_{\text{dyn}}^{k}) = \inf(S^{k}) = \int_{\mathbb{R}^{2}} \frac{1}{2} |y - x|^{2} P_{01}^{k}(dxdy) + o_{k \to \infty}(1)$$

this suggests that

•  $dX_t = dB_t$ , R-a.s.

$$\mathsf{''lim}_{k o \infty}(\mathsf{S}^k) = (\mathsf{MK}_2)$$
"

### Cooled down Schrödinger problem

$$H(\pi|R_{01}^k)/k \to \min; \quad \pi \in \mathcal{P}(\mathcal{X}^2) : \pi_0 = \mu_0, \pi_1 = \mu_1$$
 (S<sup>k</sup>)

#### Monge-Kantorovich problem

$$\int_{\mathcal{X}^2} \frac{1}{2} |y - x|^2 \, \pi(dxdy) \to \min; \quad \pi \in \mathcal{P}(\mathcal{X}^2) : \pi_0 = \mu_0, \pi_1 = \mu_1 \quad (\mathsf{MK}_2)$$

#### Theorem

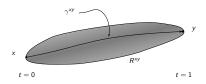
- $\Gamma$   $\lim_{k\to\infty} (\mathsf{S}^k) = (\mathsf{MK}_2)$
- $\lim_{k\to\infty}\inf(\mathsf{S}^k)=\inf(\mathsf{MK}_2):=W_2^2(\mu_0,\mu_1)/2$
- " $\lim_{k\to\infty} \pi^k = \pi$ ": solution of (MK<sub>2</sub>)

Mikami (2004), L. (2012)

### Cooling particles down brings geometry (Brownian case)

 $\lim_{k \to \infty} R^{k, xy} = \delta_{\gamma^{xy}}, \quad \gamma^{xy}$  : constant speed geodesic

• k = 1:



• k = 10:



•  $k = \infty$ :



### Convergence schema

$$P^{k}(d\omega) = \int_{\mathcal{X}^{2}} R^{k,xy}(d\omega) \quad \pi^{k}(dxdy)$$

$$\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow$$

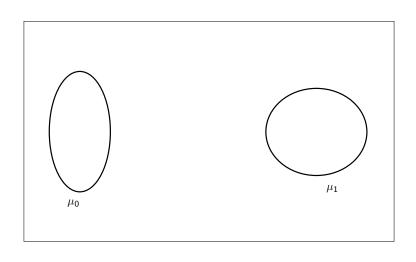
$$P(d\omega) = \int_{\mathcal{X}^{2}} \delta_{\gamma^{xy}}(d\omega) \quad \pi(dxdy)$$

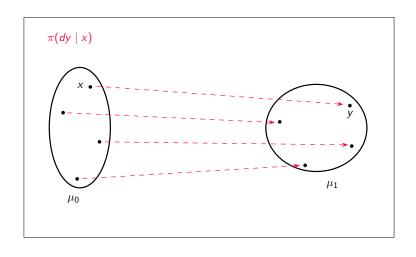
### Entropic interpolations converge to displacement interpolations

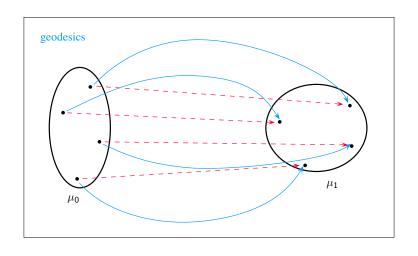
$$\begin{array}{rcl} \mu_t^k(dz) & = & \int_{\mathcal{X}^2} & R_t^{k,xy}(dz) & \pi^k(dxdy), & & 0 \leq t \leq 1 \\ \downarrow & & \downarrow & & \downarrow \\ \mu_t(dz) & = & \int_{\mathcal{X}^2} & \delta_{\gamma_t^{xy}}(dz) & \pi(dxdy), & & 0 \leq t \leq 1 \end{array}$$

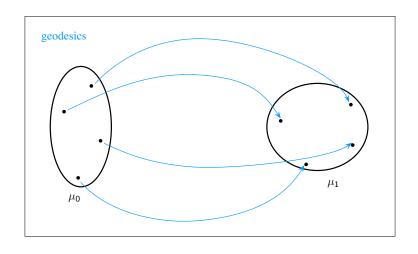
### McCann's displacement interpolation

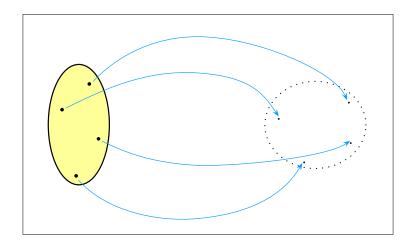
$$[\mu_0, \mu_1]^{\text{disp}} := (\mu_t)_{0 \le t \le 1}$$

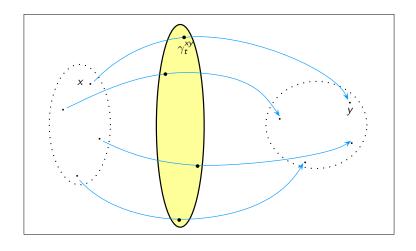


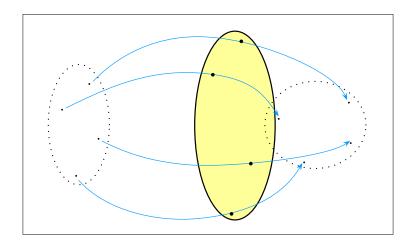


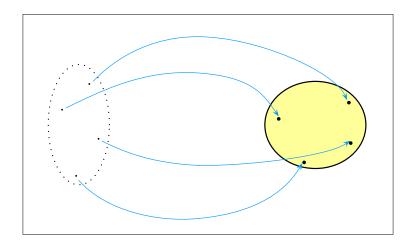












## Doubly indexed large deviation principle

• choose  $(R^k)_{k>1}$  such that it satisfies some LDP

## LDP for $(R^k)_{k>1}$

- $R^{k,x} \underset{k \to \infty}{\approx} \exp(-a_k[C + \iota_{\{X_0 = x\}}]), \quad \forall x$
- $a_k \to \infty$ ,  $C: \Omega \to [0, \infty]$ , coercive
- $\{C = 0\}$  is the limiting support of "all geodesics"

### Example (slow Brownian motion)

$$a_k = k$$
,  $C(\omega) = \int_0^1 \frac{1}{2} |\dot{\omega}_t|^2 dt$ ,  $\lim_{k \to \infty} R^{k,xy} = \delta_{\gamma^{xy}}$ 

### **Γ-convergence**

### Dynamical cooled down Schrödinger problem

$$H(P|R^k)/a_k \to \min; \quad P \in \mathcal{P}(\Omega) : P_0 = \mu_0, P_1 = \mu_1^k$$
  $(S_{\mathrm{dyn}}^k)$ 

### Dynamical Monge-Kantorovich problem

$$E_PC o \min; \quad P \in \mathcal{P}(\Omega) : P_0 = \mu_0, P_1 = \mu_1$$
 (MK<sub>dyn</sub>)

#### **Theorem**

- there exists  $\mu_1^k o \mu_1$  such that  $\Gamma ext{-lim}_{k o\infty}(\mathsf{S}^k_{ ext{dyn}}) = (\mathsf{MK}_{ ext{dyn}})$ 
  - $\lim_{k\to\infty}\inf(\mathsf{S}_{\rm dyn}^k)=\inf(\mathsf{MK}_{\rm dyn})$
  - " $\lim_{k\to\infty} P^k = P$ :" solution of  $(MK_{dyn})$

## Γ-convergence

ullet if  $P^k o P$  we get the following schema

### Convergence schema

$$P^{k}(d\omega) = \int_{\mathcal{X}^{2}} R^{k,xy}(d\omega) \pi^{k}(dxdy)$$

$$\downarrow \qquad \qquad \downarrow \qquad \qquad \downarrow$$

$$P(d\omega) = \int_{\mathcal{X}^{2}} G^{xy}(d\omega) \pi(dxdy)$$

## Entropic interpolations converge to displacement interpolation

$$egin{array}{lll} \mu_t^k( extit{d}z) &=& \int_{\mathcal{X}^2} & R_t^{k, ext{xy}}( extit{d}z) & \pi^k( extit{d}x extit{d}y), & 0 \leq t \leq 1 \ \downarrow & & \downarrow & \downarrow \ \mu_t( extit{d}z) &=& \int_{\mathcal{X}^2} & G_t^{ extit{xy}}( extit{d}z) & \pi( extit{d}x extit{d}y), & 0 \leq t \leq 1 \end{array}$$

## Definition (displacement interpolation)

$$[\mu_0, \mu_1]^{\text{disp}} := (\mu_t)_{0 \le t \le 1}$$

# $L^2$ -type displacement interpolations on a vector space

•  $\mathcal{X} = \mathbb{R}^n$ 

$$ullet R^{k,x}: \quad Z^{k,x}_t = x + k^{-1} \sum_{j=1}^{\lfloor kt \rfloor} V_j, \quad (V_j)_{j \geq 1} \text{ i.i.d. , } \mathbb{E}V = 0$$

### Mogulskii's theorem

$$C(\omega) = \int_0^1 L_V(\dot{\omega}_t) dt$$

- $L_V(v) = \sup_p \{p \cdot v H_V(p)\}; \quad H_V(p) = \log \mathbb{E} \exp(p \cdot V)$
- $L_V(v) = \sup_{\rho \in \mathcal{P}} \{\rho \cdot v H_V(\rho)\}, \quad H_V(\rho) = \log \mathbb{E} \exp(\rho \cdot v)$ •  $L_V(v) = O_{v \to 0}(|v|^2)$

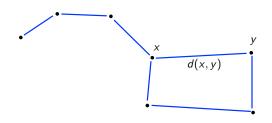
## Contraction principle

# $c(x, y) = \inf\{C(\omega); \omega \in \Omega : \omega_0 = x, \omega_1 = y\}$

•  $c(x, y) = L_V(y - x)$ 

## Interpolations on a discrete graph

ullet metric graph  $(\mathcal{X},\sim,d)$ 



 $x \sim y$  means that (x, y) is an edge

- ullet length  $\ell(\omega) := \sum_{0 \leq t \leq 1} \mathbf{1}_{\{\omega_t = 
  eq \omega_t\}} d(\omega_{t^-}, \omega_t)$
- intrinsic distance  $d(x,y) = \inf \{\ell(\omega) : \omega \in \Omega, \omega_0 = x, \omega_1 = y\}$

## Interpolations on a discrete graph

- to recover d:
  - ► slow down the walk
  - condition at t = 0 and t = 1
- reference walk:  $R \in \mathcal{P}(\Omega)$  with jump kernel  $J_x(dy) = \sum_{y:y \sim x} J_x(y) \, \delta_y$

Lazy random walks  $R^k$ 

$$J_x^k(dy) = \sum_{y:y \sim x} k^{-d(x,y)} J_x(y) \delta_y$$

## Interpolations on a discrete graph

### Geodesics

$$\Gamma^{xy} := \{ \omega \in \Omega; \omega_0 = x, \omega_1 = y, \ell(\omega) = d(x, y) \}$$
  
$$\Gamma := \cup_{x,y} \Gamma^{xy}$$

### Convergence of bridges

$$\lim_{k \to \infty} R^{k,xy} = G^{xy} \in \mathcal{P}(\Gamma^{xy})$$

• 
$$G := \mathbf{1}_{\Gamma} e^{\int_0^1 J_{X_t}(\mathcal{X}) dt} R$$

### Convergence of the interpolations

$$a_k = \log k$$
,  $C = \ell$ ,  $c = d$ 

• 
$$\lim_{k\to\infty}\inf(\mathsf{S}^k)=W_1(\mu_0,\mu_1)$$

# $L^1$ -type interpolations on a diffuse length space

•  $(\mathcal{X}, d)$ : diffuse metric space

## Definition (diffuse metric space)

 $(\mathcal{X},d)$  is diffuse if there exists a Borel measure m on  $\mathcal{X}$  such that

- $\sup_{x} m(B_{x}^{1}) < \infty$
- $m(B_x^{\epsilon}) > 0$ ,  $\forall x \in \mathcal{X}, \forall \epsilon > 0$
- $\bullet \ B_{\mathsf{x}}^{\epsilon} := \{ y \in \mathcal{X} : \mathsf{d}(\mathsf{x}, \mathsf{y}) < \epsilon \}$

# L¹-type interpolations on a diffuse length space

### Reference processes

- $R \leftrightarrow J_{\times}(dy) = \mathbf{1}_{B^1_{\circ}}(y) \, m(dy)$
- $\bullet \ R^k \leftrightarrow J_x^k(dy) = e^{-1} \mathbf{1}_{S_x^{1/k}}(y) \ m(dy)$
- $S_x^{\epsilon} := B_x^{\epsilon} \setminus B_x^{\epsilon \epsilon^2}$

### Convergence of the entropic interpolations

$$a_k = k$$
,  $C = \text{length}$ ,  $c = d$ 

- $H(P|R^k) = H(P|R) E_P \log(dR^k/dR)$
- $-\log(dR^k/dR)/k = \#\{t: X_{t^-} \neq X_t\}/k + O_{k\to\infty}(1/k)$ =  $\operatorname{length}(X) + O_{k\to\infty}(1/k)$
- work in progress with Luca Tamanini

## References about entropic interpolations

- T. Mikami. Monge's problem with a quadratic cost by the zero-noise limit of h-path processes. PTRF, (2004)
- 2 L. From the Schrödinger problem to the Monge-Kantorovich problem. JFA, (2012)
- L. A survey of the Schrödinger problem and some of its connections with optimal transport. DCDS (A), (2014)
- L. Lazy random walks and optimal transport on graphs.
   AoP, (to appear)

