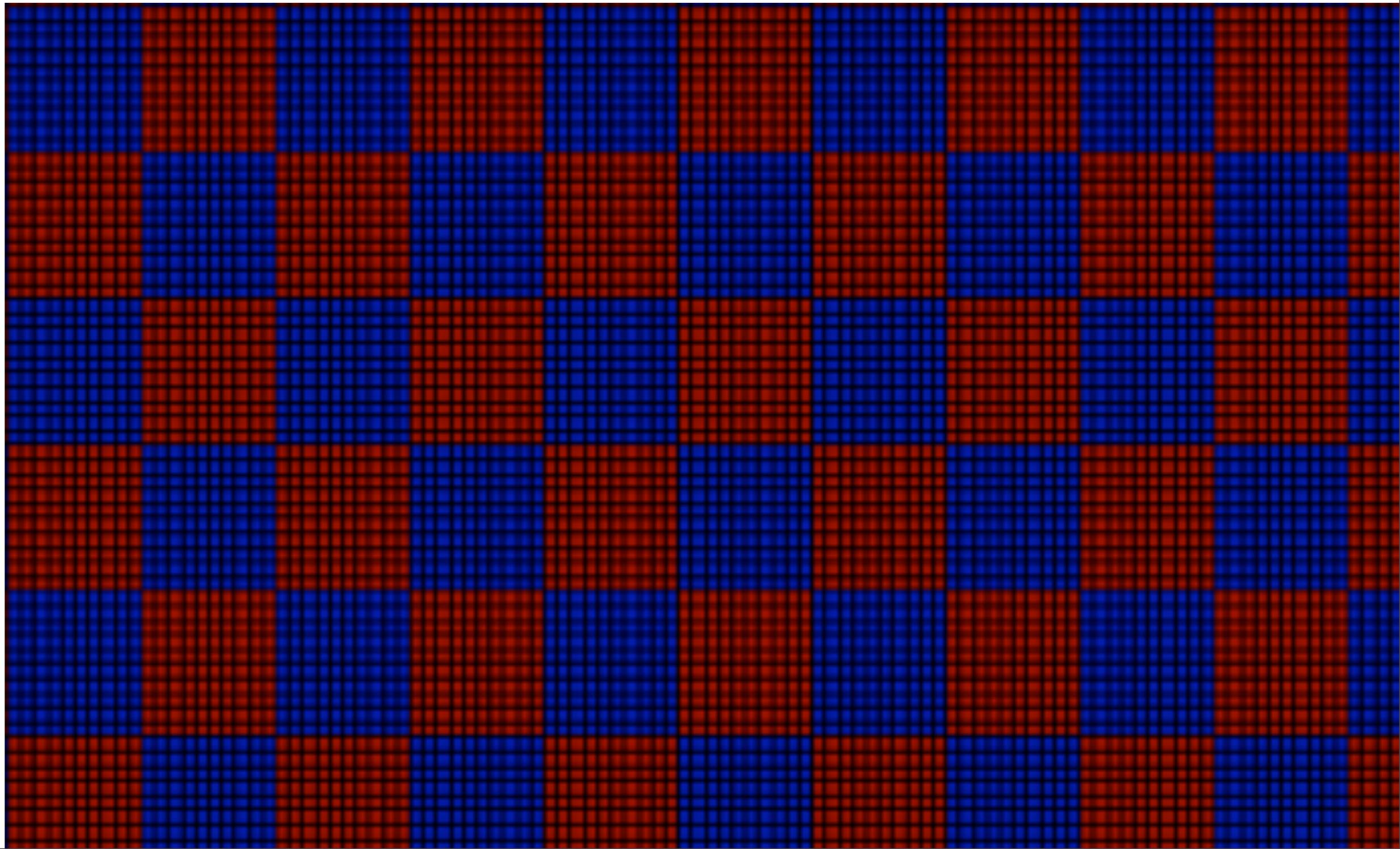


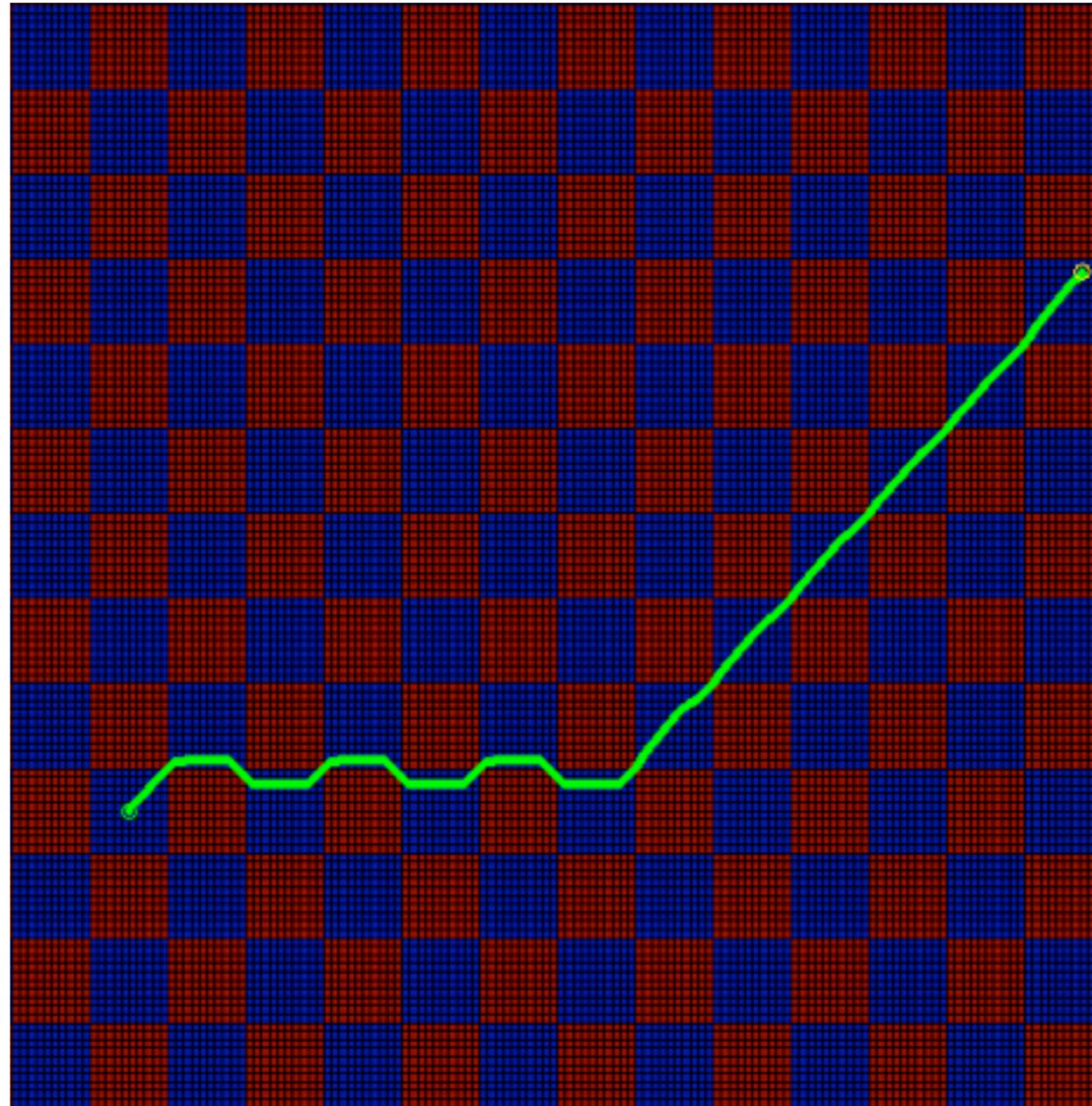
Homogenization of Metric Hamilton- Jacobi equations

Adam Oberman,
Ryo Takei
Alex Vladimirkysy

Optimal Path Checkerboard

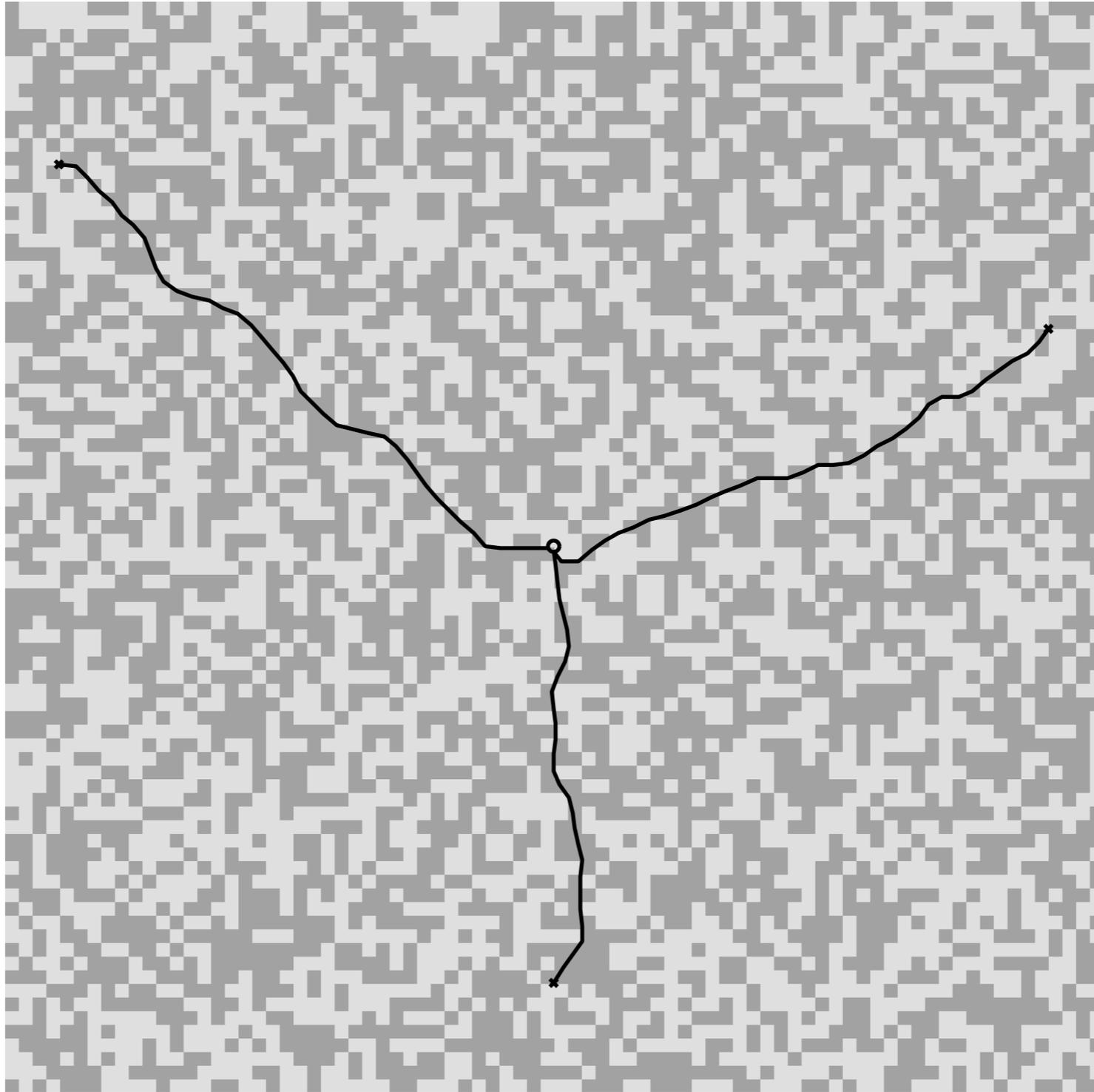


Optimal Path in Checkerboard

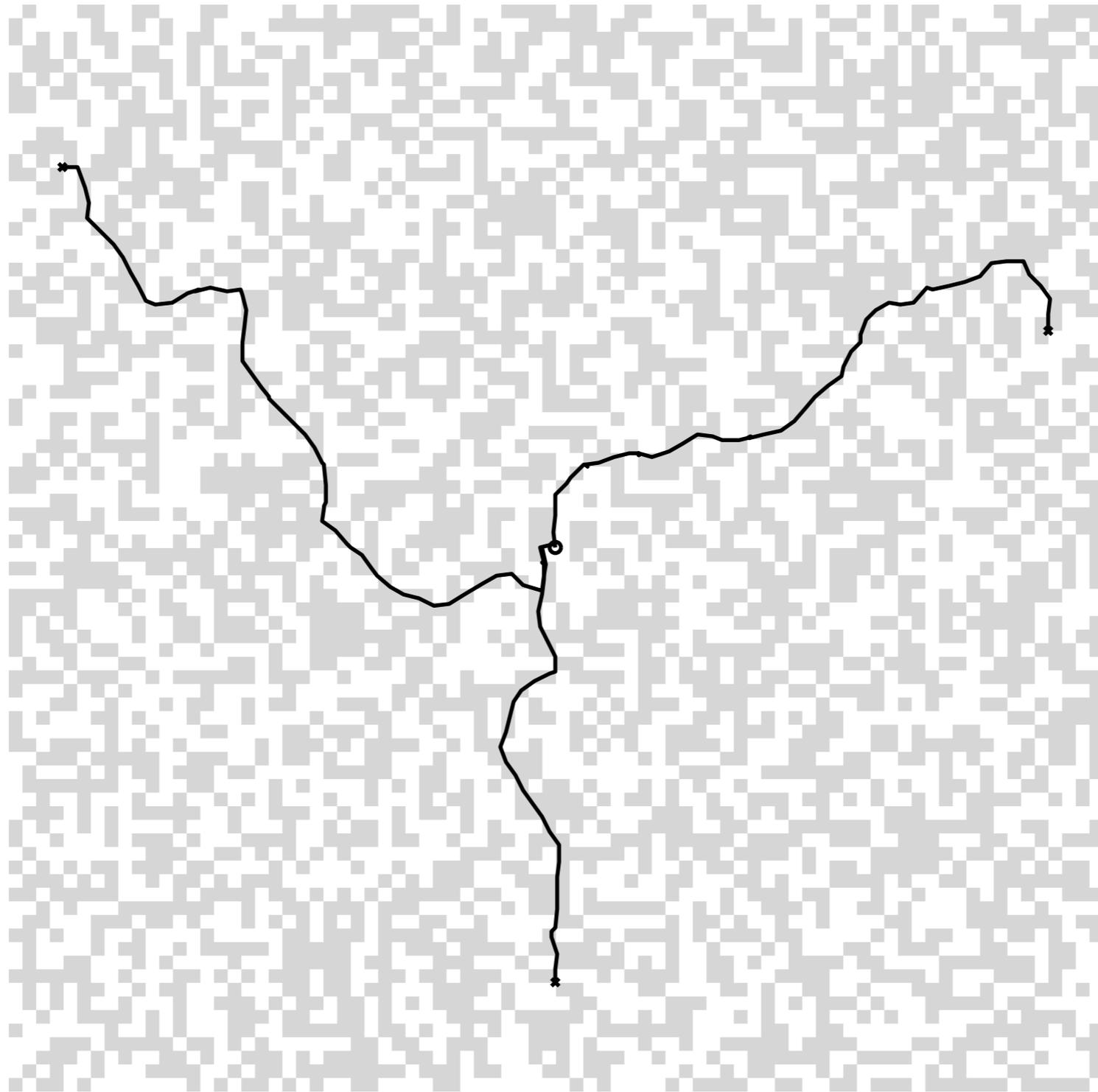


Optimal Paths in Random Media

$c = 1, 2$



Optimal Paths in Random Media



$c = 1, 10$

Optimal Paths via Dynamic Programming

- Finding optimal paths via direct search is too costly.
- Finding optimal paths via Euler-Lagrange equation is not informative.
- Cell problem requires one problem for each direction.
- Instead, use Hamilton-Jacobi equation for function whose level sets are arrival times from the origin.

$$c(x) |\nabla T| = 1$$

$$T(0) = 0$$

Hamilton-Jacobi Eqn for General Speeds

$$v_n(x) = \max_{\|\alpha\| \leq 1} \{n \cdot f(x, \alpha)\}$$

$f(x, \alpha)$ particle speed in direction α

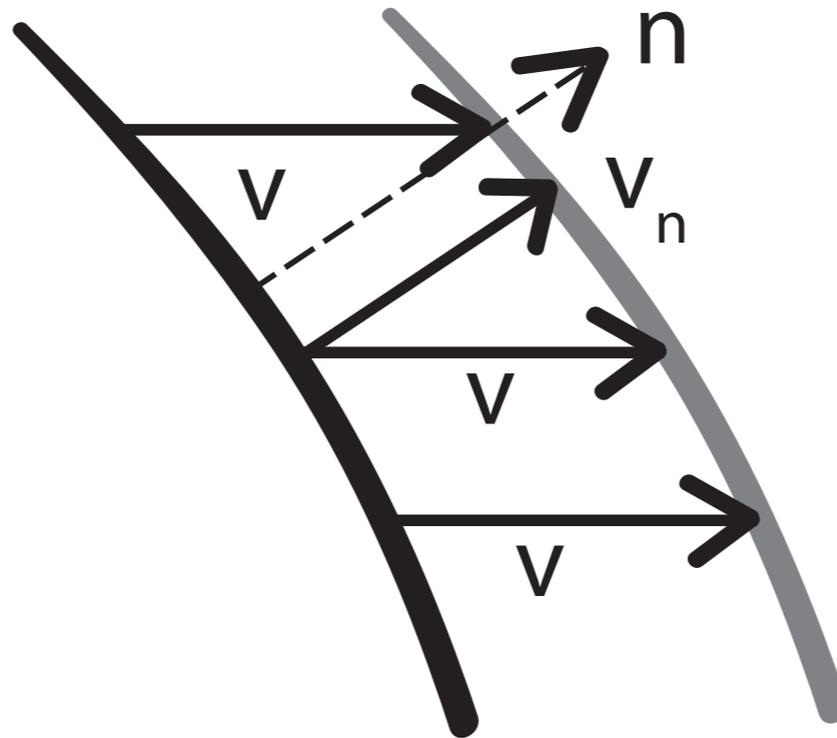
Then the function $T(x)$ whose level sets $\{T(x) = s\}$ are the position of the front at time s solves the Hamilton-Jacobi equation

$$H(\nabla T(x), x) = 1$$

where H is given by

$$H(p, x) := \sup_{\|\alpha\| \leq 1} \{p \cdot f(x, \alpha)\}$$

Front Speeds vs. Particle Speeds



- In the anisotropic case, normal velocity of fronts is not the same as particle velocity.
- Suppose particles move horizontally with speed 1 , and have no vertical speed allowed. Then front moves with normal speed given by projection of the normal onto the particle speed.

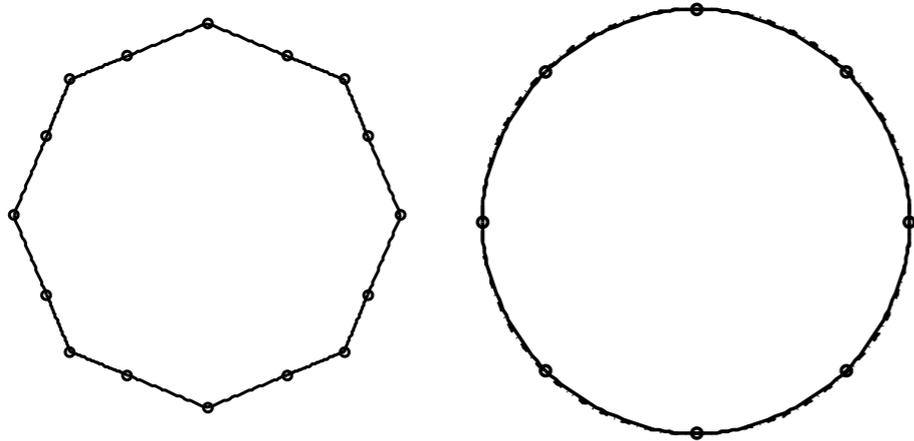
Controlled ODEs for paths

Let x denote the position, and α denote the direction of motion.

Write $\dot{x}(s, \alpha(s)) := \frac{d}{ds}x(s, \alpha(s))$. The admissible paths $x(t, \alpha(t))$ satisfy

$$\text{(ODE)} \quad \dot{x}(s, \alpha(s)) = f(x(s), \alpha(s))$$

where $\alpha(\cdot) \in \mathcal{A}$ is the *control*.



vectogram $V_c \subset \mathbb{R}^n$.

$$V_c = \{c(x, \alpha)\alpha \mid |\alpha| \leq 1\}.$$

$$T_c(x_1, x_2) = \inf_{x(\cdot) \text{ admissible}} \{t \mid x(0) = x_1, x(t) = x_2\},$$

this defines a metric on the ambient space

Optimal Paths via Lagrangian

Can Homogenize the Lagrangian:

$$\bar{L}(q) = \lim_{T \rightarrow \infty} \inf \frac{1}{T} \inf_{\phi \in H_0^1(0,T)} \int_0^T L(qt + \phi(t), q + \dot{\phi}(t)) dt$$

In this case, the minimization of over curves is performed for each value of q , and the Hamiltonian $\bar{H}(p)$ is recovered via the Legendre transform.

Main Result

Theorem 1. *Let $H(p, x)$ be a metric Hamiltonian which is periodic on the unit cube. The $H(p, x)$ defines a distance on \mathbb{R}^n . We can write*

$$H(p, x) = \|p\|_{b(x)^*}$$

where $b(x, p)$ is corresponding metric cost function. Then $H(p, x)$ homogenizes to $\bar{H}(p)$ which is a homogeneous metric Hamiltonian. Furthermore

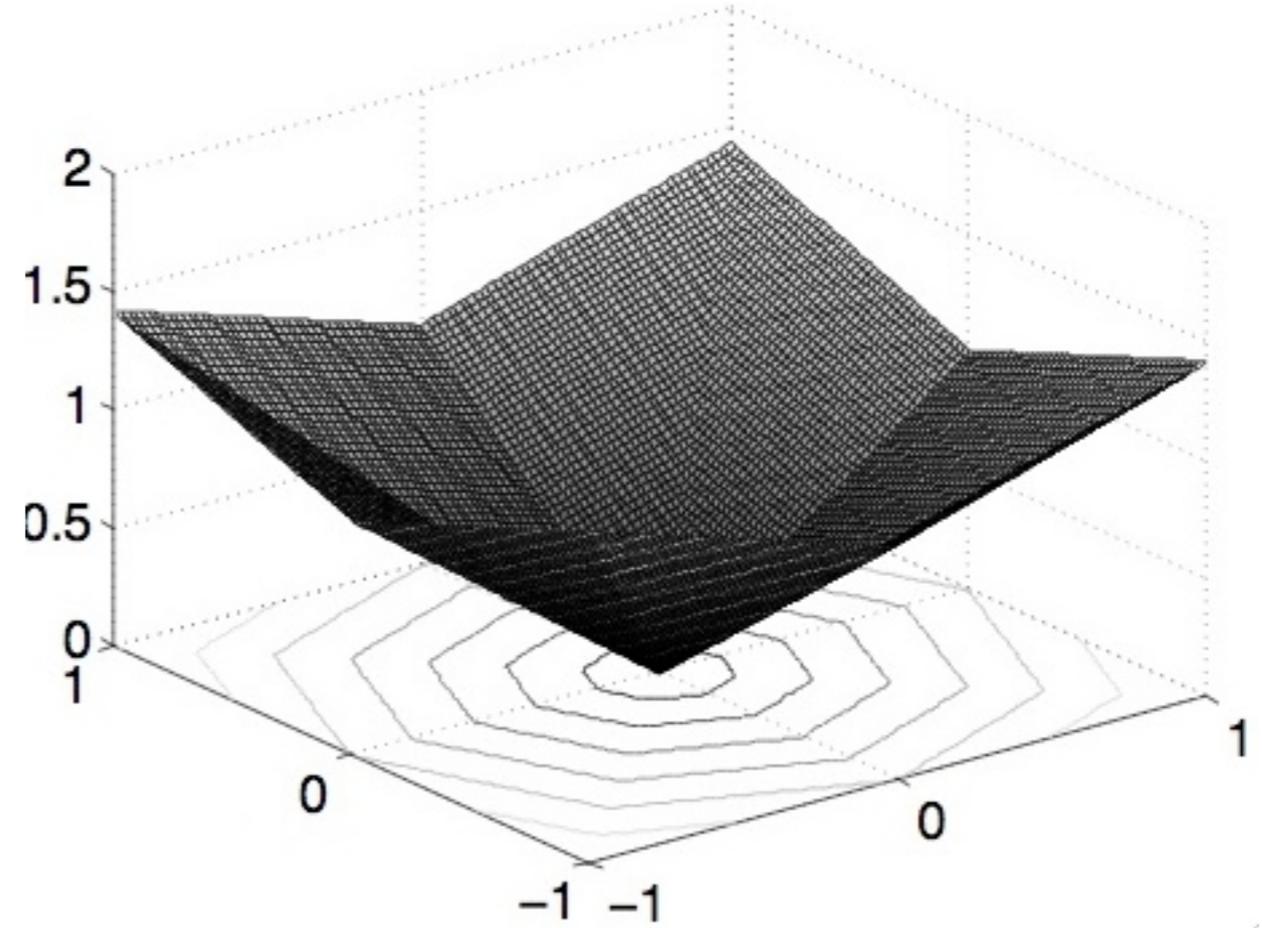
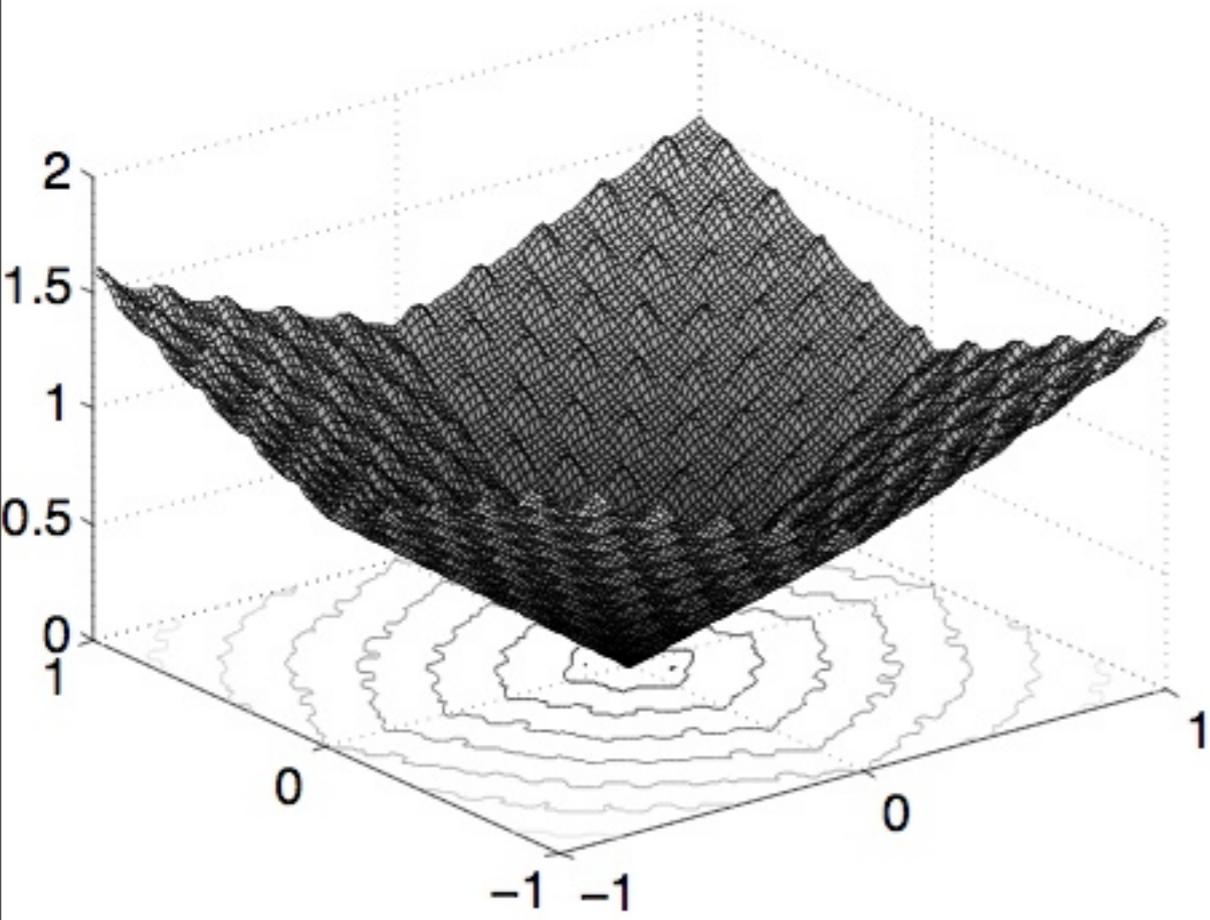
$$(17) \quad \bar{H}(p) = \|p\|_{\bar{b}^*}$$

which is the dual norm (14) of the \bar{b} norm. The values can be obtained from

$$(18) \quad \bar{b}(p) = \frac{1}{\bar{c}(p)} = \frac{T(p)}{|p|} = \lim_{\epsilon \rightarrow 0} \frac{T^\epsilon(p)}{|p|}$$

where T^ϵ is the solution of (HJ^ϵ) .

Illustration of Main Result



Computations

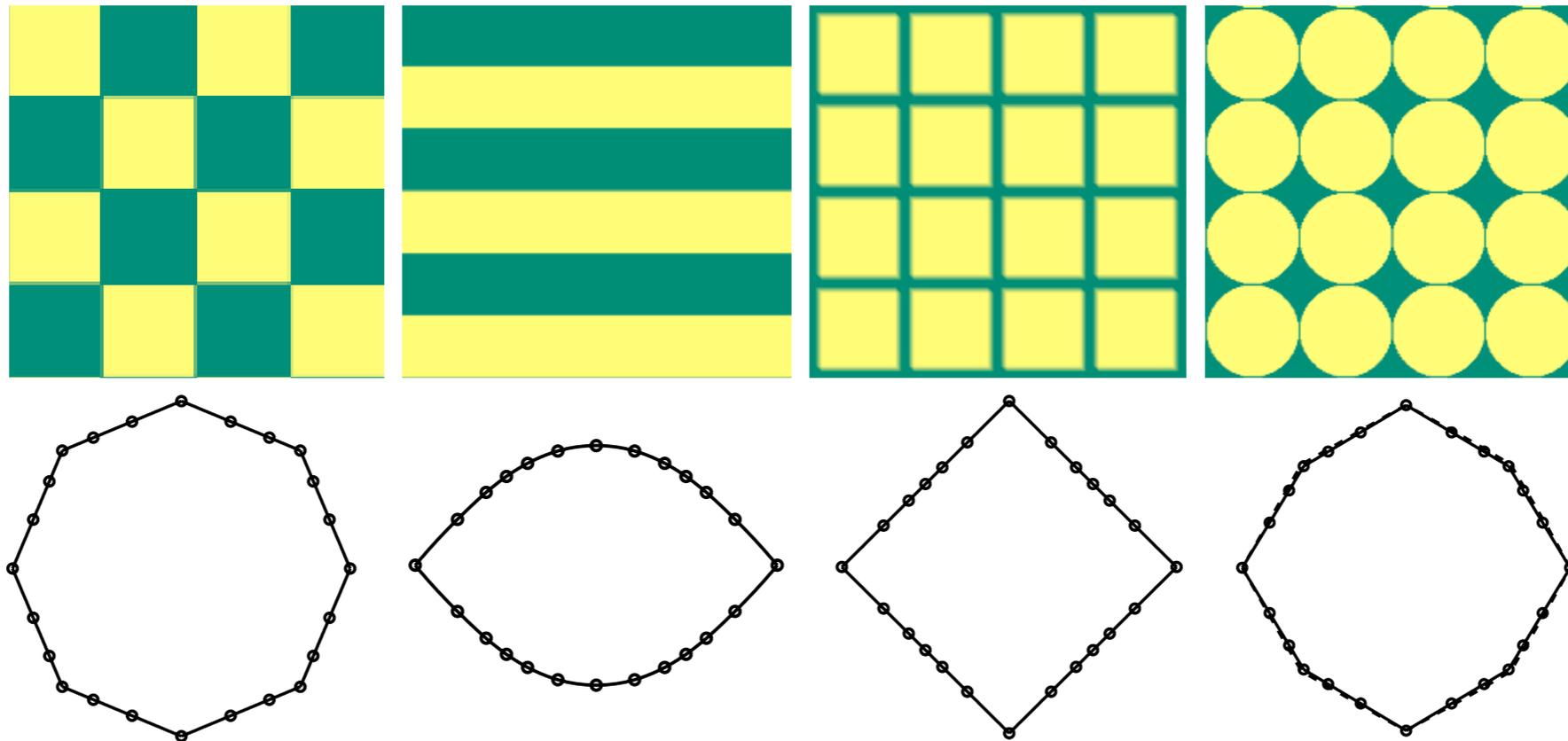


Figure 8: Period domains and computed vectograms: checkerboard, stripes, squares, and circles.

These are validations of analytic results

Open Problem: random case

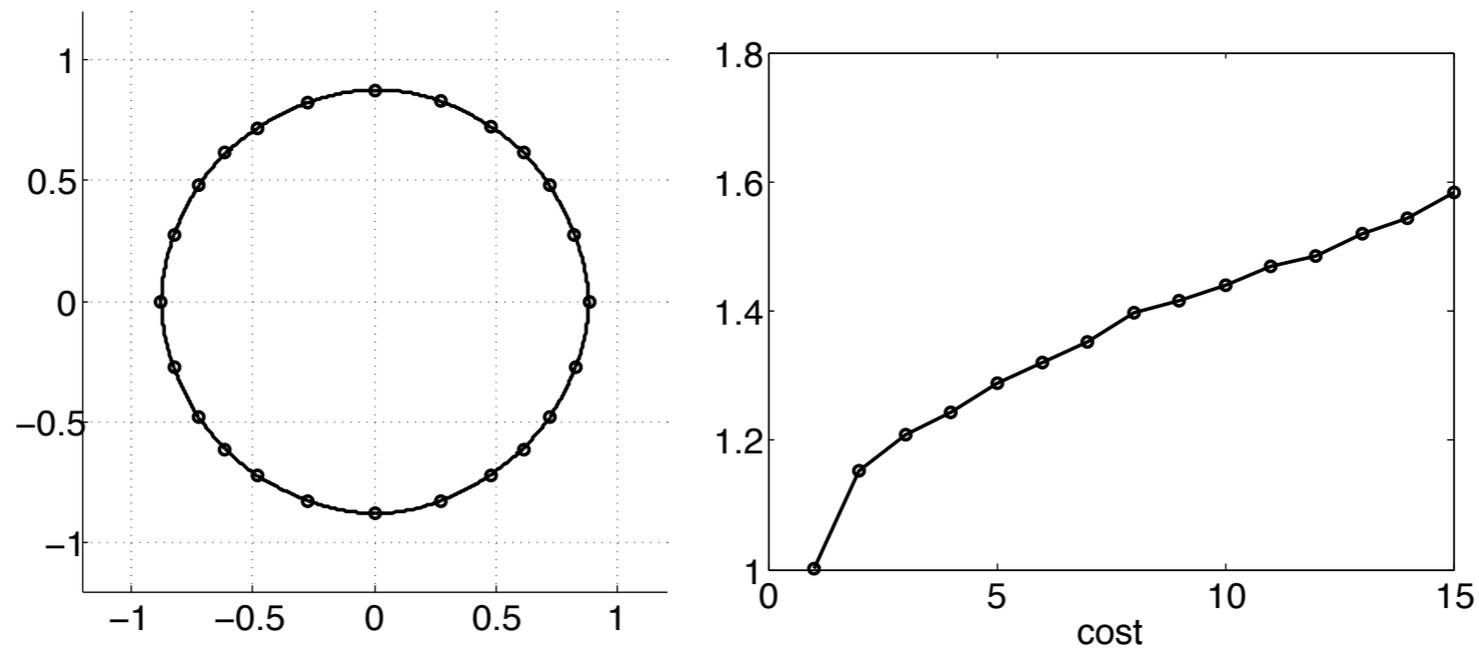
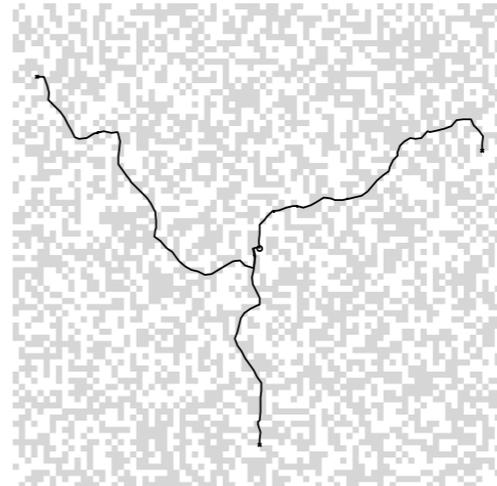
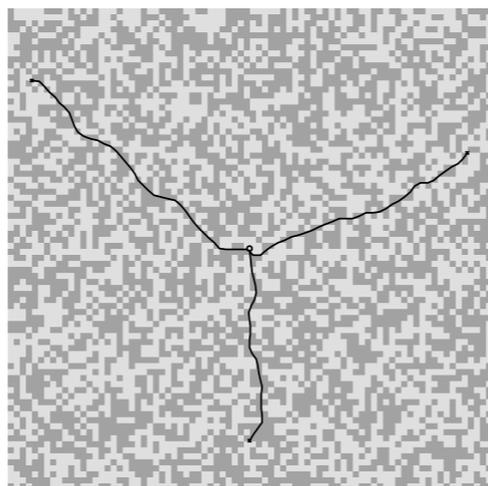


Figure 14: Illustration of the homogenized speed/cost in the random case. Left: computed vectogram, averaged over several trials, for $c = 1, 1/2$ with probability $1/2$. Right: computed homogenized cost \bar{b} as a function of the random cost $b = 1$, or $b = b_0 = 1/c_0$ with probability $1/2$.



More Open Problems

- Application: wave propagation in inhomogeneous medium.
- Anisotropic eikonal equation. Similar results?
- Add an external velocity field to the equation.
 - Now measures the combination of growth with advection. Model for combustion.
 - Expect to homogenize to either: “shifted circle” (non-Finsler metric) or “non-controllable”