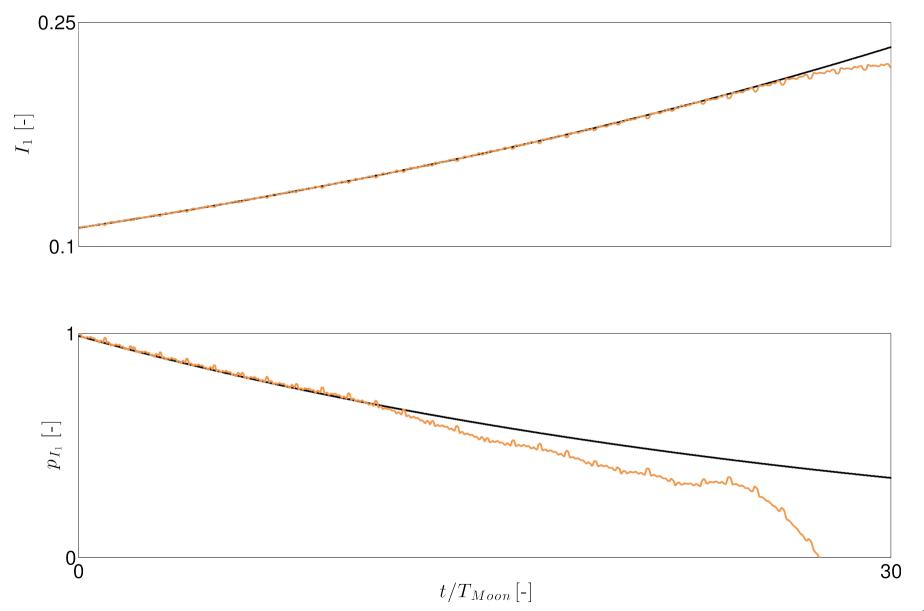
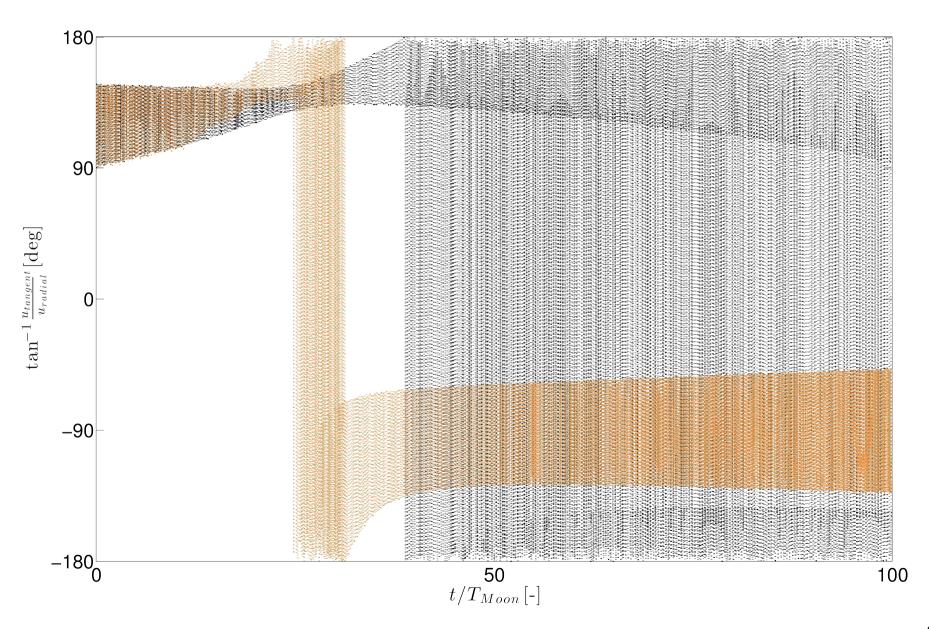
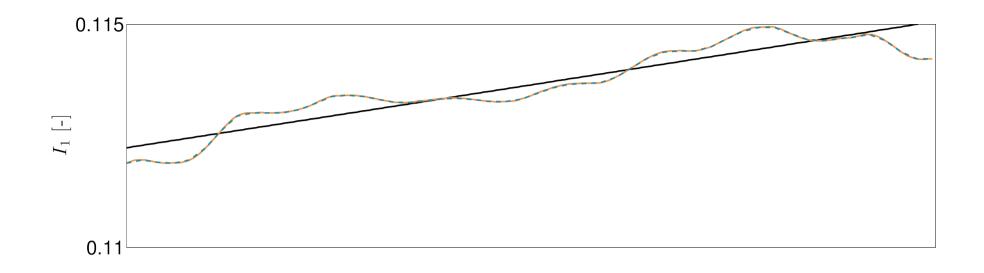
# Adjoints drift faster slow states

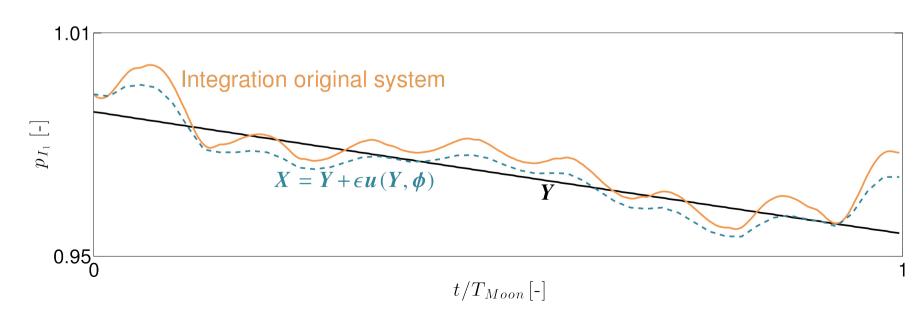


# The perturbation is the trigger, the control yields the drift



## The near-identity transformation of $p_I$ is not adequate





# The term $\omega'(I)p_{\phi}$ has to be included in the transformation

Previous transformation:

$$u(\mathbf{Y}, \boldsymbol{\phi}) = -i \sum_{0 < |\mathbf{k}| \le N} \frac{\Delta f_{\mathbf{k}}}{\mathbf{k} \cdot \boldsymbol{\omega}(\mathbf{Y})} \exp(i\mathbf{k} \cdot \boldsymbol{\phi})$$

Where  $\Delta f_k$  are Fourier coefficients of  $\Delta f \doteqdot f(Y, \phi, \epsilon) - \overline{f}(Y)$  and  $Y = \{J, p_J, p_\alpha = 0, p_\beta\}$ 

# The term $\omega'(I)p_{\phi}$ has to be included in the transformation

Previous transformation:

$$\boldsymbol{u}\left(\boldsymbol{Y},\boldsymbol{\phi}\right) = -i\sum_{0<|\boldsymbol{k}|\leq N} \frac{\Delta f_{\boldsymbol{k}}}{\boldsymbol{k}\cdot\boldsymbol{\omega}(\boldsymbol{Y})} \exp\left(i\boldsymbol{k}\cdot\boldsymbol{\phi}\right)$$

Where  $\Delta f_k$  are Fourier coefficients of  $\Delta f \doteqdot f(Y, \phi, \epsilon) - \overline{f}(Y)$  and  $Y = \{J, p_J, p_\alpha = 0, p_\beta\}$ 

The differential equation of  $p_I$  is:

$$\dot{p}_{I} = -\epsilon \frac{\partial K}{\partial I} - \underbrace{\frac{\partial |\omega|}{\partial I} p_{\alpha}}_{O(\epsilon)}$$

# The term $\omega'(I)p_{\phi}$ has to be included in the transformation

Previous transformation:

$$u(Y, \phi) = -i \sum_{0 < |\mathbf{k}| \le N} \frac{\Delta f_{\mathbf{k}}}{\mathbf{k} \cdot \omega(Y)} \exp(i\mathbf{k} \cdot \phi)$$

Where  $\Delta f_k$  are Fourier coefficients of  $\Delta f \doteqdot f(Y, \phi, \epsilon) - \overline{f}(Y)$  and  $Y = \{J, p_J, p_\alpha = 0, p_\beta\}$ 

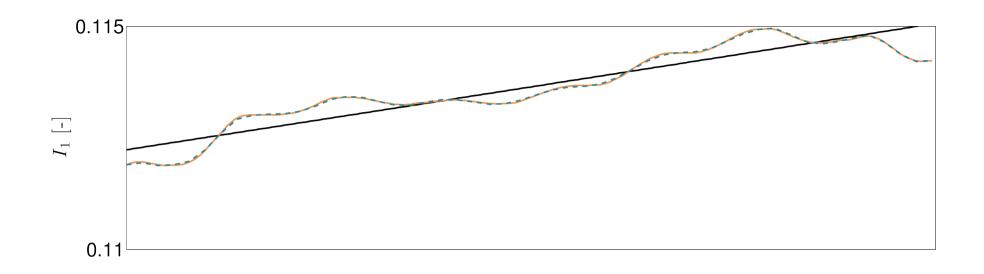
The differential equation of  $p_I$  is:

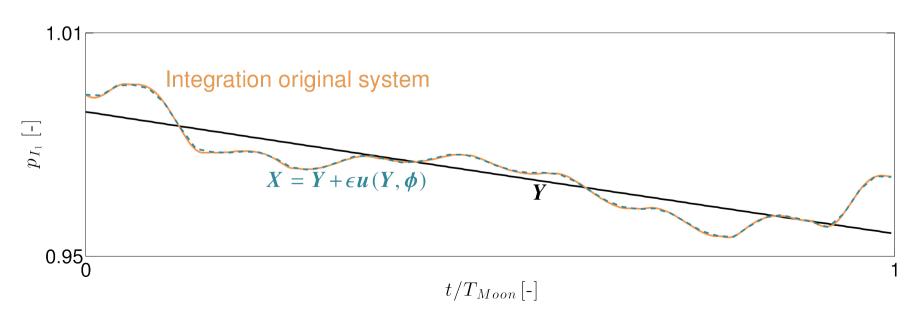
$$\dot{\mathbf{p}_{I}} = -\epsilon \frac{\partial K}{\partial \mathbf{I}} - \underbrace{\frac{\partial |\boldsymbol{\omega}|}{\partial \mathbf{I}} p_{\alpha}}_{O(\epsilon)}$$

The expansion should be carried out by using:

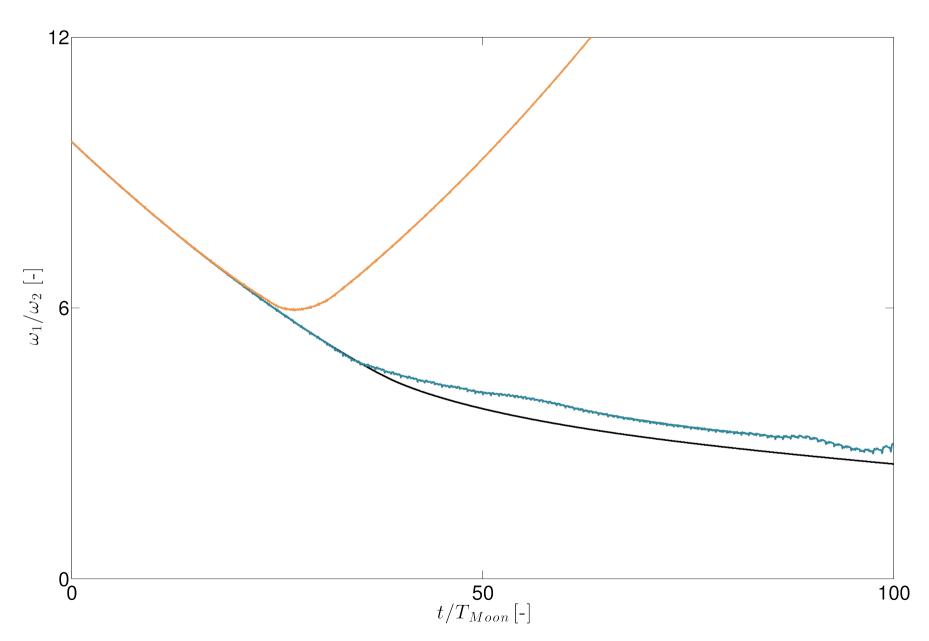
$$\dot{p_I} = -\epsilon \left( \frac{\partial K}{\partial I} - \frac{\partial |\omega|}{\partial I} h(J, p_J, \beta, \phi) \right)$$

# The term $\omega'(I)p_\phi$ has to be included in the transformation





# Beneficial effect of the enhanced transformation of $p_I$



#### Conclusions

Simple averaging is sufficient for real-life problems in astrodynamics, e.g., orbit raising

Initial conditions should undergo a near-identity transformation to reduce the drift

Key role of the transformation of the adjoints of slow variables

Effects of main resonances cannot be neglected in relevant astrodynamics applications

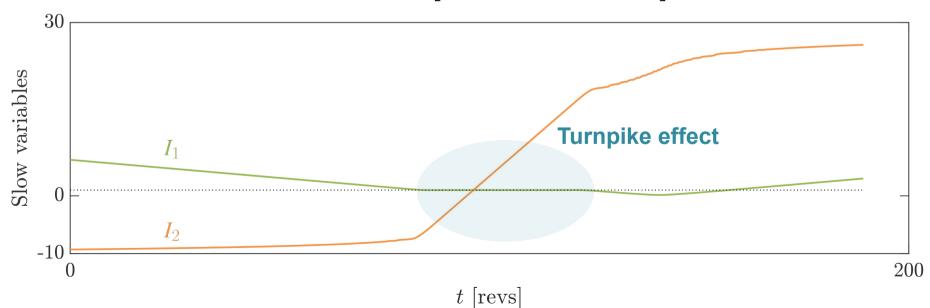
The averaged system might be enriched by using a composite expansion

## Way forward: exploitation of resonances

$$\min_{\substack{||\boldsymbol{u}||\leq 1}} t_f \quad \text{subject to}$$

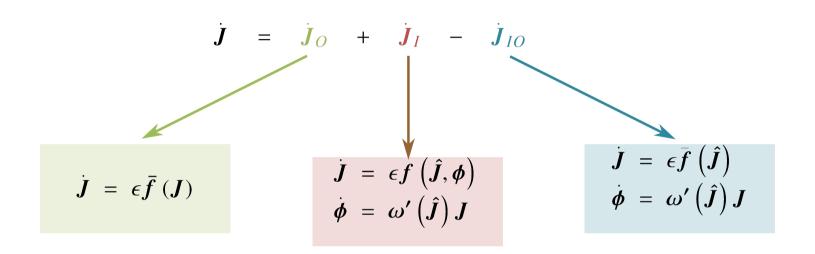
$$\begin{array}{rcl} \dot{\boldsymbol{I}}_1 &=& \epsilon \, u_1 \\ \dot{\boldsymbol{I}}_2 &=& \epsilon \, f(\phi_1,\phi_2) \, u_2 \\ \dot{\phi}_1 &=& \boldsymbol{I}_1 \\ \dot{\phi}_2 &=& 1 \end{array}$$

where 
$$f(\phi_1, \phi_2) = \left[0.1 + \left(\frac{\cos \phi_1 + \cos \phi_2}{2}\right)^2\right]^{-1}$$



39

### Way forward: composite expansion to model captures

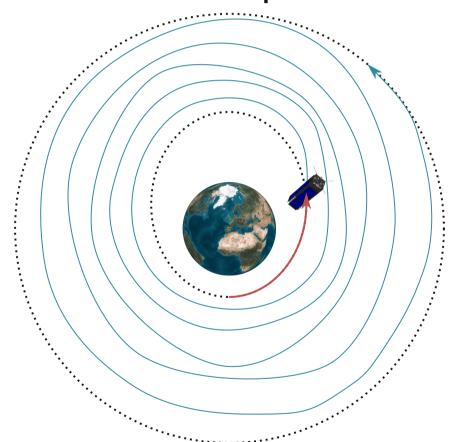


 $\exists \mathbf{k} \in \mathbb{Z}^0$  such that  $\mathbf{k} \cdot \hat{\mathbf{J}} = 0$ 

J.A. SANDERS, F. VERHULST, J. MURDOCK, *Averaging Methods in Nonlinear Dynamical Systems*, Springer, 2007.



# Two-Frequency Averaging of Optimal Control Problems with Application to Time-Optimal Orbital Transfer



L. Dell'Elce, J.B. Caillau, and J.B. Pomet