

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

based on recent works with L. Chizat, T. Gallouët, G. Peyré, B. Schmitzer.

Ceremade, Université Paris-Dauphine
INRIA team Mokaplan
Defi Imag'in CNRS: Cavalieri

Paris, Inria, October, 2016

A somewhat surprising result

Theorem

Solutions $u(t) \in C^\infty(S_1, \mathbb{R})$ to the Camassa-Holm equation

$$\partial_t u - \frac{1}{4} \partial_{txx} u + 3 \partial_x u u - \frac{1}{2} \partial_{xx} u \partial_x u - \frac{1}{4} \partial_{xxx} u u = 0 \quad (1)$$

are particular solutions of an incompressible Euler equation on $\mathbb{R}^2 \setminus \{0\}$ for a density $\rho(r, \theta) = \frac{1}{r^3} dr d\theta = \frac{1}{r^4} \text{Leb}$

$$\begin{cases} \dot{v} + \nabla_v v = -\nabla p, \\ \nabla \cdot (\rho v) = 0. \end{cases} \quad (2)$$

Outline

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

- 1 Arnold's point of view
- 2 Ebin-Marsden's approach and Brenier's
- 3 The Wasserstein-Fisher-Rao metric
- 4 The Camassa-Holm equation as an incompressible Euler equation

Contents

- 1 Arnold's point of view
- 2 Ebin-Marsden's approach and Brenier's
- 3 The Wasserstein-Fisher-Rao metric
- 4 The Camassa-Holm equation as an incompressible Euler equation

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Arnold's geometric point of view

Sur la géométrie différentielle des groupes de Lie de dimension infinie et ses applications à l'hydrodynamique des fluides parfaits, Ann. Inst. Fourier, 1966.

Proposition

The incompressible Euler equation is the geodesic flow of the right-invariant metric L^2 on $\text{SDiff}(M)$.

- An intrinsic point of view by Ebin and Marsden, *Groups of diffeomorphisms and the motion of an incompressible fluid*, Ann. of Math., 1970. Short time existence results for smooth initial conditions.
- An extrinsic point of view by Brenier, relaxation of the variational problem, optimal transport, polar factorization.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Arnold's remark continued

- Let (M, g) be a Riemannian manifold without boundary.
- Treat $\text{SDiff}(M)$ as an infinite dimensional Riemannian submanifold of $\text{Diff}(M)$.
- Consider the metric L^2 on $\text{SDiff}(M)$, it is right-invariant.

Proof.

Notations: $\varphi \in \text{SDiff}(M)$, $X \in T_\varphi \text{SDiff}(M)$. Since $\varphi_*(\text{vol}) = \text{vol}$,

$$G(\varphi)(X, X) = \int_M |X|^2 \, \text{dvol} = \int_M |X \circ \varphi^{-1}|^2 \, \text{dvol}$$



Right-invariant metric on a Lie group

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Definition (Right-invariant metric)

Let $g_1, g_2 \in G$ be two group elements, the distance between g_1 and g_2 can be defined by:

$$d^2(g_1, g_2) = \inf_{g(t)} \left\{ \int_0^1 \|v(t)\|_{\mathfrak{g}}^2 dt \mid g(0) = g_0 \text{ and } g(1) = g_1 \right\}$$

where $\partial_t g(t)g(t)^{-1} = v(t) \in \mathfrak{g}$, with \mathfrak{g} the Lie algebra.

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Right-invariance simply means:

$$d^2(g_1 g, g_2 g) = d(g_1, g_2).$$

It comes from:

$$\partial_t(g(t)g_0)(g(t)g_0)^{-1} = \partial_t g(t)g_0 g_0^{-1} g(t)^{-1} = \partial_t g(t)g(t)^{-1}.$$

Euler-Arnold-Poincaré equation

Compute the Euler-Lagrange equation of the distance functional:

$$\frac{\partial L}{\partial g} - \frac{d}{dt} \frac{\partial L}{\partial \dot{g}} = 0$$

In the case of $\int_0^1 L(g, \dot{g}) dt = \int_0^1 \|u\|^2 dt$,
Euler-Poincaré-Arnold equation

$$\begin{cases} \dot{g} = u \circ g \\ \dot{u} + \text{ad}_u^* u = 0 \end{cases} \quad (3)$$

where ad_u^* is the (metric) adjoint of $\text{ad}_u v = [v, u]$.

Proof.

Compute variations of $v(t)$ in terms of $u(t) = \delta g(t)g(t)^{-1}$. Find that admissible variations on \mathfrak{g} can be written as:
 $\delta v(t) = \dot{u} - \text{ad}_v u$ for any u vanishing at 0 and 1. □

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Euler-Arnold-Poincaré equation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Consider the Hilbert scalar product on vector fields on \mathbb{R}^d
 $\langle u, u \rangle = \langle u, Lu \rangle_{L^2}$. Denoting $m = Lu$,

$$\partial_t m + Dm \cdot u + Du^T \cdot m + \operatorname{div}(u)m = 0. \quad (4)$$

For example, the L^2 metric, $L = \operatorname{Id}$ gives:

$$\partial_t u + Du \cdot u + Du^T \cdot u + \operatorname{div}(u)u = 0. \quad (5)$$

On the group of volume preserving diffeomorphisms of (M, μ) with the L^2 metric:

Euler's equation for ideal fluid where $\operatorname{div}(u) = 0$

$$\partial_t u + \nabla_u u = -\nabla p,$$

(use $\operatorname{div}(u) = 0$ and write the term $Du^T \cdot u$ as a gradient as $\frac{1}{2} \nabla \langle u, u \rangle$)

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Fluid dynamics examples of Euler-Arnold equations

- Incompressible Euler equation.
- Korteweg-de-Vries equation.
- **Camassa-Holm equation 1981/1993.** *An integrable shallow water equation with peaked solitons*

Consider $\text{Diff}(S_1)$ endowed with the H^1 right-invariant metric $\|v\|_{L^2}^2 + \frac{1}{4}\|\partial_x v\|_{L^2}^2$. One has

$$\begin{cases} \partial_t u - \frac{1}{4}\partial_{txx} u u + 3\partial_x u u - \frac{1}{2}\partial_{xx} u \partial_x u - \frac{1}{4}\partial_{xxx} u u = 0 \\ \partial_t \varphi(t, x) = u(t, \varphi(t, x)). \end{cases} \quad (6)$$

- Model for waves in shallow water.
- Completely integrable system (bi-Hamiltonian).
- Exhibits particular solutions named as peakons. (geodesics as collective Hamiltonian).
- Blow-up of solutions which gives a model for wave breaking.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Contents

- 1 Arnold's point of view
- 2 Ebin-Marsden's approach and Brenier's
- 3 The Wasserstein-Fisher-Rao metric
- 4 The Camassa-Holm equation as an incompressible Euler equation

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Ebin-Marsden analytical framework

Rewrite the metric in Lagrangian coordinates φ and a tangent vector X_φ and realize that it is smooth...

- The right-invariant H^{div} metric:

$$G_\varphi(X_\varphi, X_\varphi) = \int_M a^2 |X_\varphi \circ \varphi^{-1}|^2 + b^2 \operatorname{div}(X_\varphi \circ \varphi^{-1})^2 d\mu. \quad (7)$$

can be written

$$G_\varphi(X_\varphi, X_\varphi) = \int_M a^2 |X_\varphi|^2 \operatorname{Jac}(\varphi) + b^2 (\operatorname{Tr}(DX_\varphi \cdot [D\varphi]^{-1}))^2 \operatorname{Jac}(\varphi) d\mu.$$

Smooth metric on an infinite dimensional Riemannian manifold.

Consequences:

- Geodesic equations is a simple ODE (No need for a Riemannian connection)
- Gauss lemma on H^s for $s > d/2 + 2$
- Geodesics are minimizing within H^s topology.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

The case of incompressible Euler

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Need to deal with the projection on $\text{SDiff}(M)$...

Geodesic equation on $\text{SDiff}(M)$ as a submanifold of $\text{Diff}(M)$:

$$\ddot{\varphi} = -\nabla p \circ \varphi. \quad (8)$$

where $-\nabla p = A(\varphi, \dot{\varphi})$ is a smooth function of $\varphi, \dot{\varphi}$.

More explicitly,

$$A(\varphi, \dot{\varphi}) = (\nabla \Delta)_{\varphi}^{-1} [\text{div}, \nabla_{\nu}]_{\varphi} \dot{\varphi}. \quad (9)$$

where, if L is a differential operator on functions,

$$L_{\varphi}(f) := L(f \circ \varphi^{-1}) \circ \varphi. \quad (10)$$

If $\varphi(t)$ is a smooth curve in H^s s.t. $\dot{\varphi} = w \circ \varphi$, then

$$\frac{d}{ds} L_{\varphi} = [L_{\varphi}, \nabla_w].$$

About Brenier's approach to incompressible Euler

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Variational approach to geodesics on $\text{SDiff}(M)$ isometrically embedded in a Hilbert space.

- Smooth solutions of Euler are minimizing (for $t \in [0, 1]$) if $\nabla^2 p$ is bounded in L^∞ (by π).
- In general, relaxation of the boundary value problem as (infinite) multimarginal optimal transport.
- Polar factorization as a nonlinear extension of the pressure.
- Benamou-Brenier's dynamic formulation.

A geometric picture: Otto's Riemannian submersion

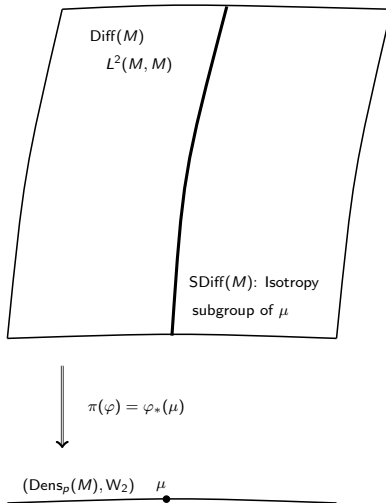


Figure – A Riemannian submersion: $\text{SDiff}(M)$ as a Riemannian submanifold of $L^2(M, M)$: Incompressible Euler equation on $\text{SDiff}(M)$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Reminders: Riemannian submersion

Let (M, g_M) and (N, g_N) be two Riemannian manifolds and $f : M \mapsto N$ a differentiable mapping.

Definition

The map f is a Riemannian submersion if f is a submersion and for any $x \in M$, the map $df_x : \text{Ker}(df_x)^\perp \mapsto T_{f(x)}N$ is an isometry.

- $\text{Vert}_{f(x)} := \text{Ker}(df(x))$ is the vertical space.
- $\text{Hor}_{f(x)} \stackrel{\text{def.}}{=} \text{Ker}(df(x))^\perp$ is the horizontal space.
- Geodesics on N can be lifted "horizontally" to geodesics on M .

Theorem (O'Neill's formula)

Let f be a Riemannian submersion and X, Y be two orthonormal vector fields on M with horizontal lifts \tilde{X} and \tilde{Y} , then

$$K_N(X, Y) = K_M(\tilde{X}, \tilde{Y}) + \frac{3}{4} \|\text{vert}([\tilde{X}, \tilde{Y}])\|_M^2, \quad (11)$$

where K denotes the sectional curvature and vert the orthogonal projection on the vertical space.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

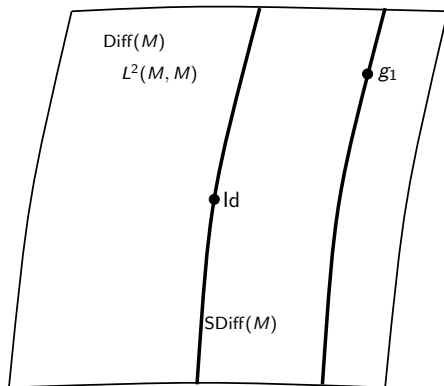
Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

A pre-formulation of the polar factorization



$$\pi(\varphi) = \varphi_*(\mu)$$

$(\text{Dens}_p(M), W_2)$ μ

Figure – A Riemannian submersion: $\text{SDiff}(M)$ as a Riemannian submanifold of $L^2(M, M)$: Incompressible Euler equation on $\text{SDiff}(M)$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Reminders: Static Formulation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Monge formulation (1781)

Let $\mu, \nu \in \mathcal{P}_+(M)$,

$$\text{Minimize } \int_{\Omega} c(x, \varphi(x)) d\mu \quad (12)$$

among the map s.t. $\varphi_*(\mu) = \nu$.

- 1 ill posed problem, the constraint may not be satisfied.
- 2 the constraint can hardly be made weakly closed.

→ Relaxation of the Monge problem.

Reminders: Static Formulation

Kantorovich formulation (1942)

Let $\mu, \nu \in \mathcal{P}_+(\Omega)$, define D by

$$D(\mu, \nu) = \inf_{\gamma \in \mathcal{P}(\Omega^2)} \left\{ \int_{\Omega^2} c(x, y) d\gamma(x, y) : \pi_*^1 \gamma = \mu \text{ and } \pi_*^2 \gamma = \nu \right\}$$

- 1 Existence result: c lower semi-continuous and bounded from below.
- 2 Also valid in Polish spaces.
- 3 If $c(x, y) = \frac{1}{p} |x - y|^p$, $D^{1/p}$ is the Wasserstein distance denoted by W_p .

Linear optimization problem and associated numerical methods. Recently introduced, entropic regularization. (C. Léonard, M. Cuturi)

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Reminders: Dynamic formulation (Benamou-Brenier)

For geodesic costs, for instance $c(x, y) = \frac{1}{2}|x - y|^2$

$$\inf \mathcal{E}(v) = \frac{1}{2} \int_0^1 \int_{\Omega} |v(x)|^2 \rho(x) \, dx \, dt, \quad (13)$$

s.t.

$$\begin{cases} \dot{\rho} + \nabla \cdot (v\rho) = 0 \\ \rho(0) = \mu_0 \text{ and } \rho(1) = \mu_1. \end{cases} \quad (14)$$

Convex reformulation: Change of variable: momentum $m = \rho v$,

$$\inf \mathcal{E}(m) = \frac{1}{2} \int_0^1 \int_{\Omega} \frac{|m(x)|^2}{\rho(x)} \, dx \, dt, \quad (15)$$

s.t.

$$\begin{cases} \dot{\rho} + \nabla \cdot m = 0 \\ \rho(0) = \mu_0 \text{ and } \rho(1) = \mu_1. \end{cases} \quad (16)$$

where $(\rho, m) \in \mathcal{M}([0, 1] \times \Omega, \mathbb{R} \times \mathbb{R}^d)$.

Existence of minimizers: Fenchel-Rockafellar.

Numerics: First-order splitting algorithm: Douglas-Rachford.

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Starting point and initial motivation

- Extend the Wasserstein L^2 distance to positive Radon measures.
- Develop associated numerical algorithms.

Possible applications: Imaging, machine learning, gradient flows, ...

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Bibliography before (june) 2015

Taking into account locally the change of mass:

Two directions: Static and dynamic.

- Static, Partial Optimal Transport [FG10]
- Static, Hanin 1992, Benamou and Brenier 2001.
- Dynamic, Numerics, Metamorphoses [MRSS15]
- Dynamic, Numerics, Growth model [LM13]
- Dynamic and static, [PR13, PR14]
- ...

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Bibliography after june 2015

More than 300 pages on the same model!

Starting point: Dynamic formulation

- Dynamic, Numerics, Imaging [CSPV15] (40 pages)
- Dynamic, Geometry and Static [CSPV15] (40 pages)
- Dynamic, Gradient flow [KMV15] (50 pages)
- Dynamic, Gradient flow [LMS15b] (40 pages)
- Static and more [LMS15a] (100 pages)
- Static relaxation of OT, Numerics [FZM⁺15] (20 pages)

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Contents

- 1 Arnold's point of view
- 2 Ebin-Marsden's approach and Brenier's
- 3 The Wasserstein-Fisher-Rao metric
- 4 The Camassa-Holm equation as an incompressible Euler equation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

An extension of Benamou-Brenier formulation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Add a source term in the constraint: (weak sense)

$$\dot{\rho} = -\nabla \cdot (\rho v) + \alpha \rho,$$

where α can be understood as the growth rate.

$$\begin{aligned} \text{WF}(m, \alpha)^2 &= \frac{1}{2} \int_0^1 \int_{\Omega} |v(x, t)|^2 \rho(x, t) \, dx \, dt \\ &\quad + \frac{\delta^2}{2} \int_0^1 \int_{\Omega} \alpha(x, t)^2 \rho(x, t) \, dx \, dt. \end{aligned}$$

Remark: very natural and not studied before...

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Convex reformulation

Add a source term in the constraint: (weak sense)

$$\dot{\rho} = -\nabla \cdot m + \mu.$$

$$\text{WF}(m, \mu)^2 = \frac{1}{2} \int_0^1 \int_{\Omega} \frac{|m(x, t)|^2}{\rho(x, t)} dx dt + \frac{\delta^2}{2} \int_0^1 \int_{\Omega} \frac{\mu(x, t)^2}{\rho(x, t)} dx dt.$$

- Fisher-Rao metric: Hessian of the Boltzmann entropy/
Kullback-Leibler divergence and reparametrization invariant.
Wasserstein metric on the space of variances in 1D.
- Convex and 1-homogeneous: convex analysis (existence and more)
- Numerics: First-order splitting algorithm: Douglas-Rachford.
- Code available at
<https://github.com/lchizat/optimal-transport/>

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

A general framework

Definition (Infinitesimal cost)

An infinitesimal cost is $f : \Omega \times \mathbb{R} \times \mathbb{R}^d \times \mathbb{R} \rightarrow \mathbb{R}_+ \cup \{+\infty\}$ such that for all $x \in \Omega$, $f(x, \cdot, \cdot, \cdot)$ is convex, positively 1-homogeneous, lower semicontinuous and satisfies

$$f(x, \rho, m, \mu) \begin{cases} = 0 & \text{if } (m, \mu) = (0, 0) \text{ and } \rho \geq 0 \\ > 0 & \text{if } |m| \text{ or } |\mu| > 0 \\ = +\infty & \text{if } \rho < 0. \end{cases}$$

Definition (Dynamic problem)

For $(\rho, m, \mu) \in \mathcal{M}([0, 1] \times \Omega)^{1+d+1}$, let

$$J(\rho, m, \mu) \stackrel{\text{def.}}{=} \int_0^1 \int_{\Omega} f(x, \frac{d\rho}{d\lambda}, \frac{dm}{d\lambda}, \frac{d\mu}{d\lambda}) d\lambda(t, x) \quad (17)$$

The dynamic problem is, for $\rho_0, \rho_1 \in \mathcal{M}_+(\Omega)$,

$$C(\rho_0, \rho_1) \stackrel{\text{def.}}{=} \inf_{(\rho, \omega, \zeta) \in \mathcal{CE}_0^1(\rho_0, \rho_1)} J(\rho, \omega, \zeta). \quad (18)$$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Existence of minimizers

Proposition (Fenchel-Rockafellar)

Let $B(x)$ be the polar set of $f(x, \cdot, \cdot, \cdot)$ for all $x \in \Omega$ and assume it is a lower semicontinuous set-valued function. Then the minimum of (18) is attained and it holds

$$C_D(\rho_0, \rho_1) = \sup_{\varphi \in K} \int_{\Omega} \varphi(1, \cdot) d\rho_1 - \int_{\Omega} \varphi(0, \cdot) d\rho_0 \quad (19)$$

with $K \stackrel{\text{def.}}{=} \{ \varphi \in C^1([0, 1] \times \Omega) : (\partial_t \varphi, \nabla \varphi, \varphi) \in B(x), \forall (t, x) \in [0, 1] \times \Omega \}$.

$$\text{WF}(x, y, z) = \begin{cases} \frac{|y|^2 + \delta^2 z^2}{2x} & \text{if } x > 0, \\ 0 & \text{if } (x, |y|, z) = (0, 0, 0) \\ +\infty & \text{otherwise} \end{cases}$$

and the corresponding Hamilton-Jacobi equation is

$$\partial_t \varphi + \frac{1}{2} \left(|\nabla \varphi|^2 + \frac{\varphi^2}{\delta^2} \right) \leq 0.$$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Numerical simulations

Initial and final densities: gray and blue curves.

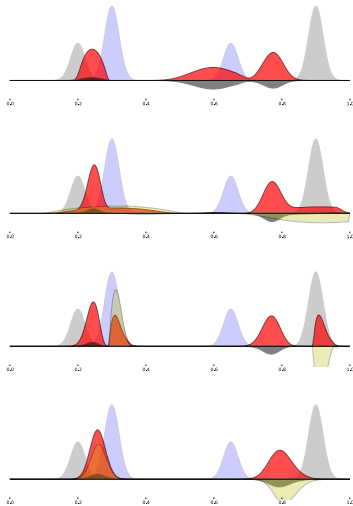


Figure – 1st row: Standard W_2 , 2nd row: Non-homogeneous L^2 , 3rd row: Partial OT, 4th row: Wasserstein-Fisher-Rao

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Numerical simulations

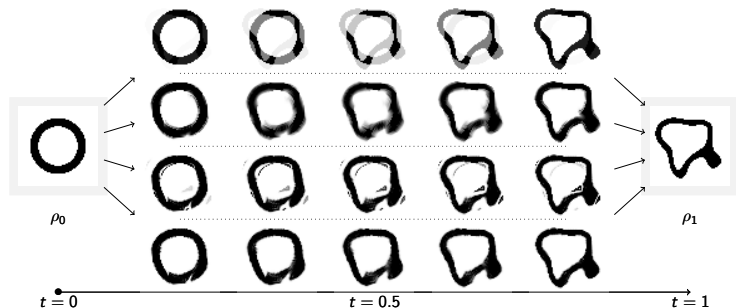


Figure – Geodesics between ρ_0 and ρ_1 for (1st row) Hellinger, (2nd row) W_2 , (3rd row) partial OT, (4th row) WF.

An Interpolating Distance between Optimal Transport and Fisher-Rao, L. Chizat, B. Schmitzer, G. Peyré, and F.-X. Vialard, FoCM, 2016.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

From dynamic to static

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Group action

Mass can be moved and changed: consider $m(t)\delta_{x(t)}$.

Infinitesimal action

$$\dot{\rho} = -\nabla \cdot (v\rho) + \mu \Leftrightarrow \begin{cases} \dot{x}(t) = v(x(t)) \\ \dot{m}(t) = \mu(x(t)) \end{cases}$$

A cone metric

$$\text{WF}^2(x, m) = \frac{1}{2} \left(m \, dx^2 + \frac{dm^2}{m} \right),$$

Riemannian cone

Definition

Let (M, g) be a Riemannian manifold. The cone over (M, g) is the Riemannian manifold $(M \times \mathbb{R}_+^*, r^2g + dr^2)$.

- Change of variable: $WF^2 = \frac{1}{2}r^2g + 2dr^2$.
- Non complete metric space: add the vertex $M \times \{0\}$.
- The distance:

$$d((x_1, m_1), (x_2, m_2))^2 = m_2 + m_1 - 2\sqrt{m_1 m_2} \cos\left(\frac{1}{2}d_M(x_1, x_2) \wedge \pi\right). \quad (20)$$

- Curvature tensor: $R(\tilde{X}, e) = 0$ and $R(\tilde{X}, \tilde{Y})\tilde{Z} = (R_g(X, Y)Z - g(Y, Z)X + g(X, Z)Y, 0)$.
- $M = \mathbb{R}$ then $(x, m) \mapsto \sqrt{m}e^{ix/2} \in \mathbb{C}$ local isometry.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Geometry of a cone

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Corollary

If (Ω, g) has sectional curvature greater than 1, then
 $(\Omega \times \mathbb{R}_+^*, m g + \frac{1}{4m} dm^2)$ has non-negative sectional curvature.
For X, Y two orthonormal vector fields on Ω ,

$$K(\tilde{X}, \tilde{Y}) = (K_g(X, Y) - 1) \quad (21)$$

where K and K_g denote respectively the sectional curvatures of
 $\Omega \times \mathbb{R}_+^*$ and Ω .

Generalization of Otto's Riemannian submersion

Idea of a left group action:

$$\begin{aligned}\pi &: (\text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*)) \times \text{Dens}(M) \mapsto \text{Dens}(M) \\ \pi((\varphi, \lambda), \rho) &:= \varphi_*(\lambda^2 \rho)\end{aligned}$$

Group law:

$$(\varphi_1, \lambda_1) \cdot (\varphi_2, \lambda_2) = (\varphi_1 \circ \varphi_2, (\lambda_1 \circ \varphi_2) \lambda_2) \quad (22)$$

Theorem

Let $\rho_0 \in \text{Dens}(M)$ and $\pi_0 : \text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*) \mapsto \text{Dens}(M)$ defined by $\pi_0(\varphi, \lambda) := \varphi_*(\lambda^2 \rho_0)$. It is a Riemannian submersion

$$(\text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*), L^2(M, M \times \mathbb{R}_+^*)) \xrightarrow{\pi_0} (\text{Dens}(M), \text{WF})$$

(where $M \times \mathbb{R}_+^*$ is endowed with the cone metric).

O'Neill's formula: sectional curvature of $(\text{Dens}(\Omega), \text{WF})$.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Horizontal lift

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Proposition (Horizontal lift)

Let $\rho \in \text{Dens}^s(\Omega)$ be a smooth density and $X_\rho \in H^s(\Omega, \mathbb{R})$ be a tangent vector at the density ρ . The horizontal lift at $(\text{Id}, 1)$ of X_ρ is given by $(\frac{1}{2}\nabla\Phi, \Phi)$ where Φ is the solution to the elliptic partial differential equation:

$$-\text{div}(\rho\nabla\Phi) + 2\Phi\rho = X_\rho. \quad (23)$$

By elliptic regularity, the unique solution Φ belongs to $H^{s+1}(M)$.

Geometric consequence

The sectional curvature of $\text{Dens}(\Omega)$ at point ρ is:

$$K(\rho)(X_1, X_2) = \int_{\Omega} k(x, 1)(Z_1(x), Z_2(x))w(Z_1(x), Z_2(x))\rho(x) d\nu(x) + \frac{3}{4} \|[Z_1, Z_2]^V\|^2 \quad (24)$$

where

$$w(Z_1(x), Z_2(x)) = g(x)(Z_1(x), Z_1(x))g(x)(Z_2(x), Z_2(x)) - g(x)(Z_1(x), Z_2(x))^2$$

and $[Z_1, Z_2]^V$ denotes the vertical projection of $[Z_1, Z_2]$ at identity and $\|\cdot\|$ denotes the norm at identity.

Corollary

Let (Ω, g) be a compact Riemannian manifold of sectional curvature bounded below by 1, then the sectional curvature of $(\text{Dens}(\Omega), \text{WF})$ is non-negative.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Consequences

Monge formulation

$$WF(\rho_0, \rho_1) = \inf_{(\varphi, \lambda)} \{ \|(\varphi, \lambda) - (Id, 1)\|_{L^2(\rho_0)} : \varphi_*(\lambda^2 \rho_0) = \rho_1 \} \quad (25)$$

Under existence and smoothness of the minimizer, there exists a function $p \in C^\infty(M, \mathbb{R})$ such that

$$(\varphi(x), \lambda(x)) = \exp_x^{C(M)} \left(\frac{1}{2} \nabla p(x), p(x) \right), \quad (26)$$

Equivalent to Monge-Ampère equation

With $z \stackrel{\text{def.}}{=} \log(1 + p)$ one has

$$(1 + |\nabla z|^2) e^{2z} \rho_0 = \det(D\varphi) \rho_1 \circ \varphi \quad (27)$$

and

$$\varphi(x) = \exp_x^M \left(\arctan \left(\frac{1}{2} |\nabla z| \right) \frac{\nabla z(x)}{|\nabla z(x)|} \right).$$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

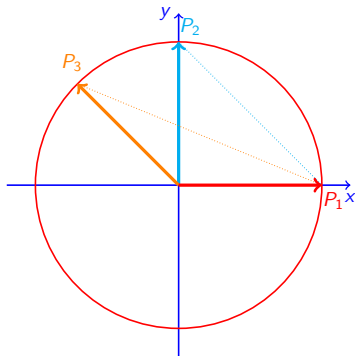
Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Distance between Diracs



$$\frac{1}{4} WF(m_1 \delta_{x_1}, m_2 \delta_{x_2})^2 = m_2 + m_1 - 2\sqrt{m_1 m_2} \cos\left(\frac{1}{2} d_\Omega(x_1, x_2) \wedge \pi/2\right).$$

Proof: prove that an explicit geodesic is a critical point of the convex functional.

Properties: positively 1-homogeneous and convex in (m_1, m_2) .

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

General Kantorovich formulation

Definition (Cost function)

A *cost function* is

$$c : \begin{array}{ll} (\Omega \times [0, +\infty[)^2 & \rightarrow [0, +\infty] \\ (x_1, m_1), (x_2, m_2) & \mapsto c(x_1, m_1, x_2, m_2) \end{array}$$

which is l.s.c. and positively 1-homogeneous and convex in (m_0, m_1) .

Example: $c_d(x_1, m_1, x_2, m_2) \stackrel{\text{def.}}{=} C_D(m_1\delta_{x_1}, m_2\delta_{x_2})$ (if l.s.c.)

Definition (Semi-couplings)

Let $\rho_1, \rho_2 \in \mathcal{M}_+(\Omega)$, the set of semi-couplings is

$$\Gamma(\rho_1, \rho_2) := \left\{ (\gamma_1, \gamma_2) \in (\mathcal{M}_+(M^2))^2 : (\text{Proj}_1)_*\gamma_1 = \rho_1, (\text{Proj}_2)_*\gamma_2 = \rho_2 \right\},$$

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Kantorovich formulation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

The functional is

$$J_K(\gamma_0, \gamma_1) \stackrel{\text{def.}}{=} \int_{\Omega^2} c\left(x, \frac{d\gamma_0}{d\gamma}, y, \frac{d\gamma_1}{d\gamma}\right) d\gamma(x, y), \quad (28)$$

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

The new Kantorovich problem is

$$C_K(\rho_0, \rho_1) \stackrel{\text{def.}}{=} \inf_{(\gamma_0, \gamma_1) \in \Gamma(\rho_0, \rho_1)} J_K(\gamma_0, \gamma_1). \quad (29)$$

- 1 If c is a *cost function* then a minimizer for $C_K(\rho_0, \rho_1)$ exists.
- 2 If $c^{1/p}$ is a metric on the cone, $C_K^{1/p}$ is a metric on $\mathcal{M}_+(\Omega)$.
- 3 If, in addition, $c(x, 1, y, 1)$ is continuous in the space variable and $c(x, 1, x, 0) < \infty$ for some x , C_K weak* continuous on $\mathcal{M}_+(\Omega)^2$.

Kantorovich formulation

Recall

$$\frac{1}{4} c_d^2(x_1, m_1, x_2, m_2) = m_2 + m_1 - 2\sqrt{m_1 m_2} \cos\left(\frac{1}{2} d_\Omega(x_1, x_2) \wedge \pi/2\right).$$

then

$$WF(\rho_1, \rho_2)^2 = \inf_{(\gamma_1, \gamma_2) \in \Gamma(\rho_1, \rho_2)} \int_{M^2} c_d^2\left(\left(x, \frac{d\gamma_1}{d\gamma}\right), \left(y, \frac{d\gamma_2}{d\gamma}\right)\right) d\gamma(x, y),$$

Theorem (Dual formulation)

$$WF^2(\rho_0, \rho_1) = \sup_{(\phi, \psi) \in C(M)^2} \int_M \phi(x) d\rho_0 + \int_M \psi(y) d\rho_1$$

subject to $\forall (x, y) \in M^2,$

$$\begin{cases} \phi(x) \leq 1, & \psi(y) \leq 1, \\ (1 - \phi(x))(1 - \psi(y)) \geq \cos^2(|x - y|/2 \wedge \pi/2) \end{cases}$$

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

A relaxed static OT formulation

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Define

$$KL(\gamma, \nu) = \int \frac{d\gamma}{d\nu} \log \left(\frac{d\gamma}{d\nu} \right) d\nu + |\nu| - |\gamma|$$

The corresponding primal formulation

$$WF^2(\rho_1, \rho_2) = KL(\text{Proj}_*^1 \gamma, \rho_1) + KL(\text{Proj}_*^2 \gamma, \rho_2) - \int_{\Omega^2} \gamma(x, y) \log(\cos^2(d(x, y)/2 \wedge \pi/2)) dx dy$$

New algorithm

Scaling Algorithms for Unbalanced Transport Problems, L. Chizat, G. Peyré, B. Schmitzer, F.-X. Vialard.

- Use of entropic regularization.

$$WF^2(\rho_1, \rho_2) = KL(\text{Proj}_*^1 \gamma, \rho_1) + KL(\text{Proj}_*^2 \gamma, \rho_2) - \int_{\Omega^2} \gamma(x, y) \log(\cos^2(d(x, y)/2 \wedge \pi/2)) dx dy + \varepsilon KL(\gamma, \mu_0).$$

- Alternate projection algorithm (contraction for a Hilbert type metric).
- Applications to color transfer, Fréchet-Karcher mean (barycenters).
- Simulations for gradient flows.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Contents

- 1 Arnold's point of view
- 2 Ebin-Marsden's approach and Brenier's
- 3 The Wasserstein-Fisher-Rao metric
- 4 The Camassa-Holm equation as an incompressible Euler equation

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

The Riemannian submersion

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

The Riemannian submersion is given by the push forward

$$\begin{aligned}\pi : \text{Diff}(M) &\rightarrow \text{Dens}_\rho(M) \\ \pi(\varphi) &= \varphi_*(\rho_0)\end{aligned}$$

between $(\text{Diff}(M), L^2(M, M))$ and $(\text{Dens}_\rho(M), W_2)$. The vertical space is

$$\text{Vert}_\varphi = \{v \circ \varphi; v \in \text{Vect}(M) \text{ s.t. } \text{div}(\rho v) = 0\}, \quad (30)$$

and the horizontal space

$$\text{Hor}_\varphi = \{\nabla p \circ \varphi; p \in C^\infty(M, \mathbb{R})\}. \quad (31)$$

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

The incompressible Euler equation

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

The usual Eulerian formulation of the equation:

$$\begin{cases} \partial_t v(t, x) + v(t, x) \cdot \nabla v(t, x) = -\nabla p(t, x), & t > 0, x \in M, \\ v(0, x) = v_0(x), \end{cases} \quad (32)$$

As a geodesic equation of the Riemannian submanifold $\text{SDiff}(M)$,
a Lagrangian formulation

$$\ddot{\phi} = -\nabla p \circ \phi. \quad (33)$$

and under the condition $\phi \in \text{SDiff}(M)$.

The Riemannian submersion for WFR

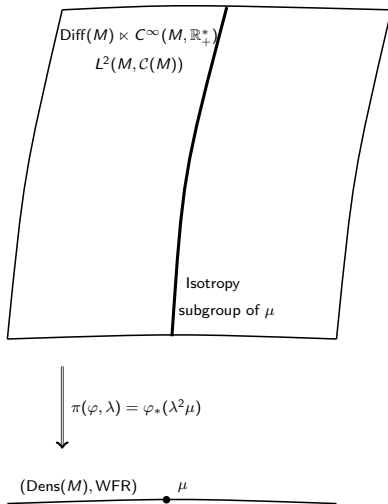


Figure – The same picture in our case: what is the corresponding equation to Euler?

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

The isotropy subgroup for unbalanced optimal transport

Recall that

$$\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \lambda) \in \text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*) : \varphi_*(\lambda^2 \rho_0) = \rho_0\}$$

$$\pi_0^{-1}(\{\rho_0\}) = \{(\varphi, \sqrt{\text{Jac}(\varphi)}) \in \text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*) : \varphi \in \text{Diff}(M)\}$$

The vertical space is

$$\text{Vert}_{(\varphi, \lambda)} = \{(v, \alpha) \circ (\varphi, \lambda) ; \text{div}(\rho v) = 2\alpha\rho\} , \quad (34)$$

where $(v, \alpha) \in \text{Vect}(M) \times C^\infty(M, \mathbb{R})$. The horizontal space is

$$\text{Hor}_{(\varphi, \lambda)} = \left\{ \left(\frac{1}{2} \nabla p, p \right) \circ (\varphi, \lambda) ; p \in C^\infty(M, \mathbb{R}) \right\} . \quad (35)$$

The induced metric is

$$G(v, \text{div } v) = \int_M |v|^2 d\mu + \frac{1}{4} \int_M |\text{div } v|^2 d\mu . \quad (36)$$

The H^{div} right-invariant metric on the group of diffeomorphisms.

An isometric embedding

We have

$$\text{inj} : (\text{Diff}(M), H^{\text{div}}) \hookrightarrow L^2(M, \mathcal{C}(M)) \quad (37)$$

The geodesic equations can be written as

$$\begin{cases} \frac{D}{Dt} \dot{\varphi} + 2 \frac{\dot{\lambda}}{\lambda} \dot{\varphi} = -\nabla^g P \circ \varphi \\ \ddot{\lambda} r - \lambda r g(\dot{\varphi}, \dot{\varphi}) = -2\lambda r P \circ \varphi. \end{cases} \quad (38)$$

Corollary (Michor and Mumford)

The distance on $\text{Diff}(M)$ with the right-invariant metric H^{Div} is non degenerate.

Proof.

Segments are length minimizing in the ambient space. \square

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Consequences

- Smooth geodesics are length minimizing for a short enough time under mild conditions (generalization of Brenier's proof).
- Using Gauss-Codazzi formula, it generalizes a curvature formula by Khesin et al. obtained on $\text{Diff}(S_1)$.

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Generalisation of Brenier's proof

Theorem

Let $(\varphi(t), r(t))$ be a smooth solution to the geodesic equations on the time interval $[t_0, t_1]$. If $(t_1 - t_0)^2 \langle w, \nabla^2 \Psi_{P(t)}(x, r) w \rangle < \pi^2 \|w\|^2$ holds for all $t \in [t_0, t_1]$ and $(x, r) \in \mathcal{C}(M)$ and $w \in T_{(x,r)}\mathcal{C}(M)$, then for every smooth curve $(\varphi_0(t), r_0(t)) \in \pi_0^{-1}(\{\mu\})$ satisfying $(\varphi_0(t_i), r_0(t_i)) = (\varphi(t_i), r_0(t_i))$ for $i = 0, 1$ and the condition (*), one has

$$\int_{t_0}^{t_1} \|(\dot{\varphi}, \dot{r})\|^2 dt \leq \int_{t_0}^{t_1} \|(\dot{\varphi}_0, \dot{r}_0)\|^2 dt, \quad (39)$$

with equality if and only if the two paths coincide on $[t_0, t_1]$.

The condition (*) is:

- 1 If the sectional curvature of $\mathcal{C}(M)$ can assume both signs, there exists $\delta > 0$ such that the curve $(\varphi_0(t), r_0(t))$ has to belong to a δ -neighborhood of $(\varphi(t), r(t))$, namely

$$d_{\mathcal{C}(M)}((\varphi_0(t, x), r_0(t, x)), (\varphi(t, x), r(t, x))) \leq \delta$$

for all $(x, t) \in M \times [t_0, t_1]$ where $d_{\mathcal{C}(M)}$ is the distance on the cone.

- 2 If $\mathcal{C}(M)$ has non positive sectional curvature, then, for every δ as above, there exists a short enough time interval on which the geodesic will be length minimizing.
- 3 If $M = S_n(r)$ the Euclidean sphere in \mathbb{R}^{n+1} of radius $r \leq 1$, the result is valid for every path $(\dot{\varphi}, \dot{r})$.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Consequences

From unbalanced
optimal transport to
the Camassa-Holm
equation

François-Xavier
Vialard

Arnold's point of view

Ebin-Marsden's
approach and
Brenier's

The Wasserstein-
Fisher-Rao
metric

The Camassa-Holm
equation as an
incompressible Euler
equation

Corollary

When $M = S_1$ and $a \leq 2b$, smooth solutions to the Camassa-Holm equation (here given in 1D)

$$a^2 \partial_t u - b^2 \partial_{txx} u + 3a^2 \partial_x u u - 2b^2 \partial_{xx} u \partial_x u - b^2 \partial_{xxx} u u = 0. \quad (40)$$

are length minimizing for short times.

Toward the incompressible Euler equation

Why? Liero, Mielke, Savaré derived the new metric using a minimization problem using probability densities on the cone.

Question

Understand $\text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^)$ as a subgroup of $\text{Diff}(\mathcal{C}(M))$?*

Proposition

The space of half-densities on M is a trivial principal fibre bundle over M which can be written $M \times \mathbb{R}_+^$ once a reference density is chosen.*

The automorphism group $\text{Aut}(\mathcal{C}(M))$ can be identified with $\text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^)$. One has $(\varphi, \lambda) : (x, r) \mapsto (\varphi(x), \lambda(x)r)$.*

Recall that $\psi \in \text{Aut}(\mathcal{C}(M))$ if $\psi \in \text{Diff}(\mathcal{C}(M))$ and $\forall \lambda \in \mathbb{R}_+^*$ one has $\psi(\lambda \cdot (x, r)) = \lambda \cdot \psi(x, r)$ where $\lambda \cdot (x, r) \stackrel{\text{def.}}{=} (x, \lambda r)$.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

CH as an incompressible Euler equation

The geodesic equation on $\text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*)$ can be extended to $\text{Aut}(\mathcal{C}(M))$ as

$$\frac{D}{Dt}(\dot{\varphi}, \dot{\lambda}r) = -\nabla \Psi_P \circ (\varphi, \lambda r), \quad (41)$$

where $\Psi_P(x, r) \stackrel{\text{def.}}{=} r^2 P(x)$.

Question

Does there exist a density $\tilde{\mu}$ on the cone such that $\text{inj}(\text{Diff}(M)) \subset \text{SDiff}_{\tilde{\mu}}(\mathcal{C}(M))$? (answer: yes)

Proof.

The measure $\tilde{\mu} \stackrel{\text{def.}}{=} r^{-3} dr d\mu$ where μ denotes the volume form on M . □

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Theorem

The solutions of the Camassa-Holm equation can be lifted to solutions to the incompressible Euler equation on the cone for a density which is non integrable at the cone point.

In Lagrangian coordinates, the correspondence is given by $\mathcal{M} : \varphi \mapsto (\varphi, \sqrt{\text{Jac}(\varphi)})$.

Case when $M = S_1$, $\mathcal{M}(\varphi) = \sqrt{\varphi'} e^{i\varphi}$ then the CH equation is

$$\begin{cases} \partial_t u - \frac{1}{4} \partial_{txx} u u + 3 \partial_x u u - \frac{1}{2} \partial_{xx} u \partial_x u - \frac{1}{4} \partial_{xxx} u u = 0 \\ \partial_t \varphi(t, x) = u(t, \varphi(t, x)). \end{cases} \quad (42)$$

The cone is $\mathcal{C}(M) = \mathbb{R}^2 \setminus \{0\}$, the density is $\frac{1}{r^4}$ Leb.

A corresponding polar factorization:

Proposition

Let $(\phi, \lambda) \in \text{Aut}(\mathcal{C}(M))$ be an element of the automorphism group of the half-densities bundle and ρ_0 . Denote by $C^\infty(\mathcal{C}(M))^{\mathbb{R}_+^*}$ functions $f : \mathcal{C}(M) \mapsto \mathbb{R}$ of the form $f(x, r) = r^2 p(x)$.

There exists a couple $(\varphi, \Psi) \in \text{Diff}(M) \times (C^\infty(\mathcal{C}(M)))^{\mathbb{R}_+^*}$ such that $\log(1 + p)$ is $-\log(\cos^2(d(x, y) \wedge \frac{\pi}{2}))$ -convex and

$$(\phi, \lambda) = \exp^{\mathcal{C}(M)}(\nabla \Psi) \circ (\varphi, \sqrt{\text{Jac}(\varphi)}), \quad (43)$$

where $(\varphi, \sqrt{\text{Jac}(\varphi)})$ is the natural lift of φ in $\text{Aut}(\mathcal{C}(M))$.

A new geometric picture

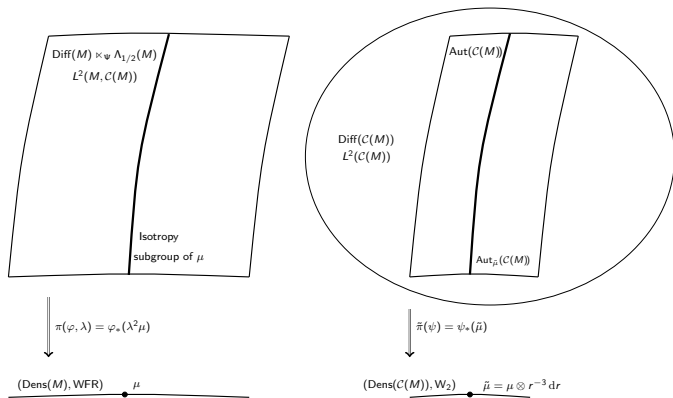


Figure – The group $\text{Aut}(C(M)) = \text{Diff}(M) \times C^\infty(M, \mathbb{R}_+^*)$ is totally geodesic in $(\text{Diff}(C(M)), L^2(C(M)))$ and there is a Riemannian submersion of the automorphism group of the cone endowed with the L^2 metric to the space of positive densities on M endowed with the Wasserstein-Fisher-Rao metric.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Perspectives

- Study the relaxation of geodesics for CH (uniqueness of the pressure, how the angle of the cone affects the results...)
- Develop numerical approaches following Mériqot et al.
- Treat other fluid dynamic equations ?

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view






Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

Figure – CH equation after the Madelung transform

References I

-  L. Chizat, B. Schmitzer, G. Peyré, and F.-X. Vialard, *An Interpolating Distance between Optimal Transport and Fisher-Rao*, ArXiv e-prints (2015).
-  A. Figalli and N. Gigli, *A new transportation distance between non-negative measures, with applications to gradients flows with dirichlet boundary conditions*, Journal de mathématiques pures et appliquées **94** (2010), no. 2, 107–130.
-  C. Frogner, C. Zhang, H. Mobahi, M. Araya-Polo, and T. Poggio, *Learning with a wasserstein loss*, Preprint 1506.05439, Arxiv, 2015.
-  S. Kondratyev, L. Monsaingeon, and D. Vorotnikov, *A new optimal transport distance on the space of finite Radon measures*, Tech. report, Pre-print, 2015.
-  D. Lombardi and E. Maitre, *Eulerian models and algorithms for unbalanced optimal transport*, <hal-00976501v3>, 2013.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard






Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation

References II

-  M. Liero, A. Mielke, and G. Savaré, *Optimal Entropy-Transport problems and a new Hellinger-Kantorovich distance between positive measures*, ArXiv e-prints (2015).
-  _____, *Optimal transport in competition with reaction: the Hellinger-Kantorovich distance and geodesic curves*, ArXiv e-prints (2015).
-  J. Maas, M. Rumpf, C. Schönlieb, and S. Simon, *A generalized model for optimal transport of images including dissipation and density modulation*, arXiv:1504.01988, 2015.
-  B. Piccoli and F. Rossi, *On properties of the Generalized Wasserstein distance*, arXiv:1304.7014, 2013.
-  _____, *Generalized Wasserstein distance and its application to transport equations with source*, Archive for Rational Mechanics and Analysis **211** (2014), no. 1, 335–358.

From unbalanced optimal transport to the Camassa-Holm equation

François-Xavier Vialard

Arnold's point of view

Ebin-Marsden's approach and Brenier's

The Wasserstein-Fisher-Rao metric

The Camassa-Holm equation as an incompressible Euler equation