

# Large Time Step scheme for Hyperbolic Systems

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*Inria*



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# 1. The CFL limitation

## 1.1 Problem position

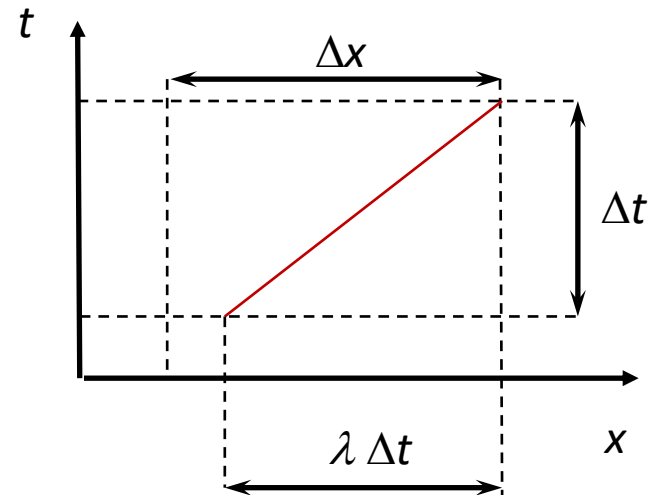
$$\partial_t \mathbf{u} + \partial_x \mathbf{f} = 0$$

$$u_i^{n+1} = u_i^n + \frac{\Delta t}{\Delta x_i} \left( \mathbf{f}_{i-1/2}^{n+1/2} - \mathbf{f}_{i+1/2}^{n+1/2} \right)$$

For stability

$$\nu = |\lambda| \frac{\Delta t}{\Delta x_i} \leq 1$$

For accuracy : keep  $\nu$  close to 1

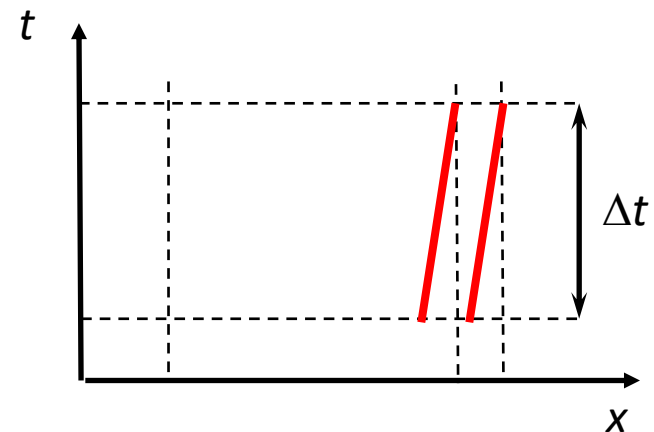
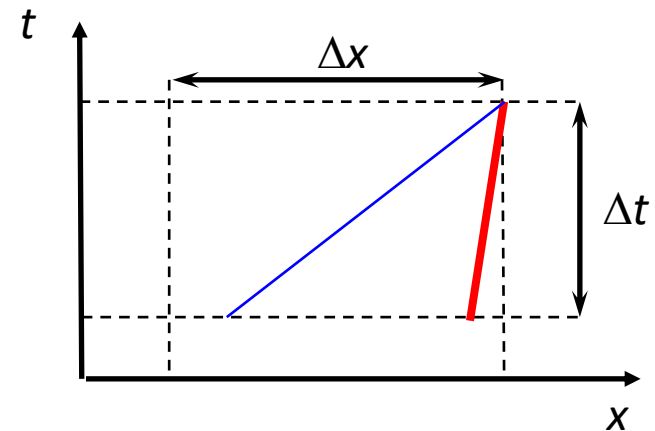


# 1. The CFL limitation

## 1.1 Problem position

Achieving  $\nu > 1$  or  $\nu \gg 1$  is desirable in at least 2 situations

- Preserving the slower waves is more important than preserving the fast ones (passive transport)
- A few narrow cells induce a locally much larger  $\nu$  than elsewhere (classical in industrial meshes)

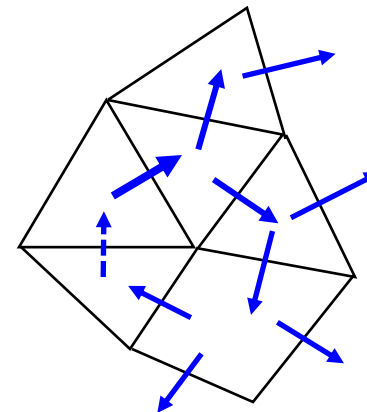
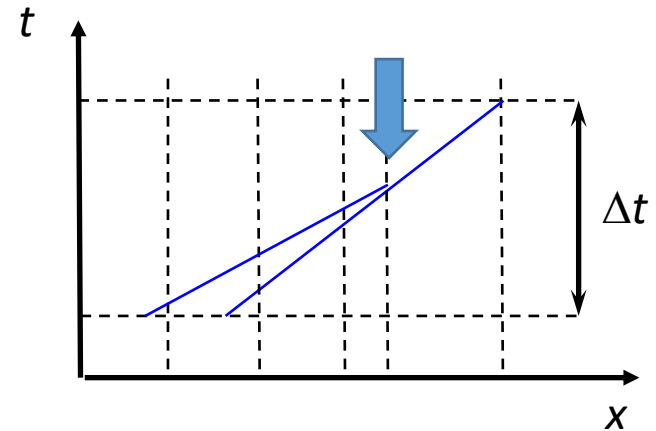


# 1. The CFL limitation

## 1.2 Most common approaches

Allow the waves to cross more than 1 cell within  $\Delta t$  : wave tracking [1] ,  
time line interpolation [2]

- Easy for 1D linear laws
- Wave interactions when the laws are non-linear
- Tricky to implement in 2D



[1] LeVeque (1982, 1988), Harten (1986), Murillo & al. (2006), Morales & Murillo (2012)

[2] Collins & al (1986), Guinot (2000)

### 2.1 The « naive » approach

Only an *averaged* flux is needed

$$\mathbf{f}_{i+1/2}^{n+1/2} = \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} \mathbf{f}(x_{i+1/2}, t) dt$$

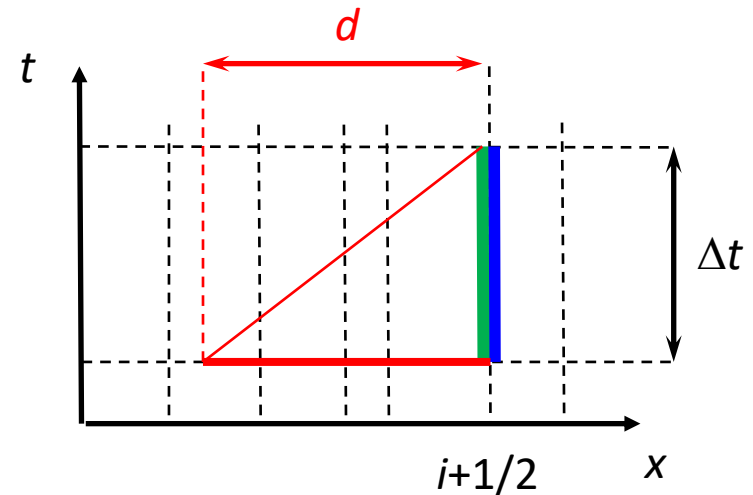
Approximation #1 : linearise wrt  $\mathbf{u}$

$$\mathbf{f}_{i+1/2}^{n+1/2} \approx \mathbf{f} \left( \frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} \mathbf{u}(x_{i+1/2}, t) dt \right)$$

Approximation #2 : convert time average to space average (characteristic form)

$$\frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} \mathbf{u}(x_{i+1/2}, t) dt \approx \frac{1}{d} \int_{x_{i+1/2}-d}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx$$

$$u_i^{n+1} = u_i^n + \frac{\Delta t}{\Delta x_i} (\mathbf{f}_{i-1/2}^{n+1/2} - \mathbf{f}_{i+1/2}^{n+1/2})$$



## 2.1 The « naive » approach

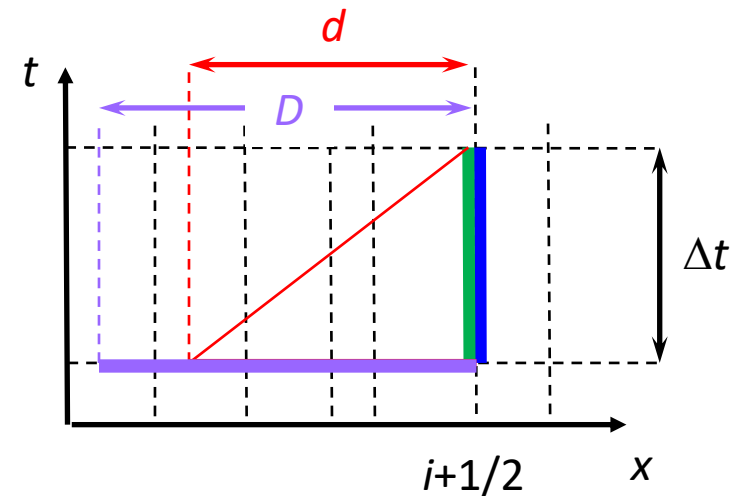
$$u_i^{n+1} = u_i^n + \frac{\Delta t}{\Delta x_i} (f_{i-1/2}^{n+1/2} - f_{i+1/2}^{n+1/2})$$

$$\frac{1}{\Delta t} \int_{t_n}^{t_{n+1}} \mathbf{u}(x_{i+1/2}, t) dt \approx \frac{1}{d} \int_{x_{i+1/2}-d}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx$$

BUT : computing  $d$  is not always straightforward (nonlinear laws)

Approximation #3 (our contribution 1):  
Do not use the exact  $d$  but a pre-defined  $D$

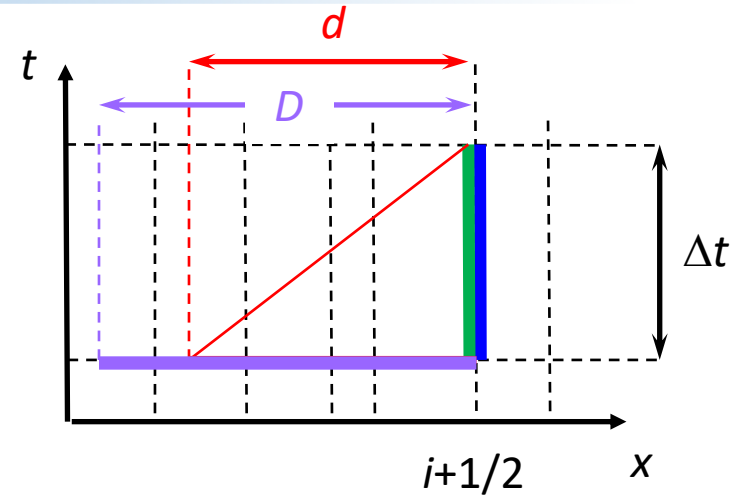
$$\frac{1}{d} \int_{x_{i+1/2}-d}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx \approx \frac{1}{D} \int_{x_{i+1/2}-D}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx$$



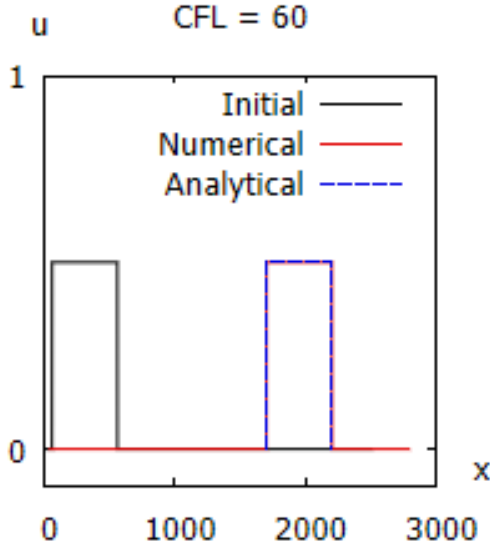
### 2.1 The « naive » approach

What happens if  $D \neq d$  ?

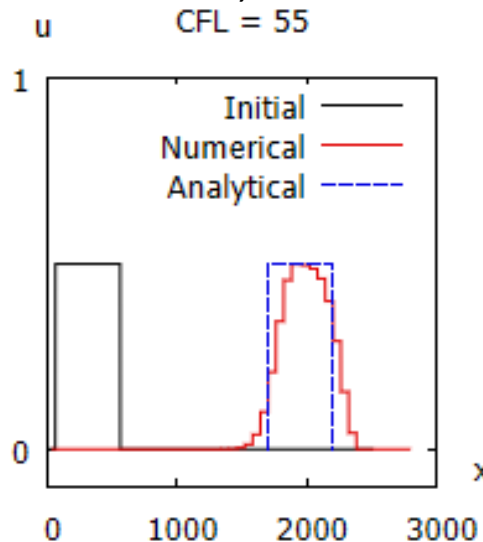
Linear advection of a hat function



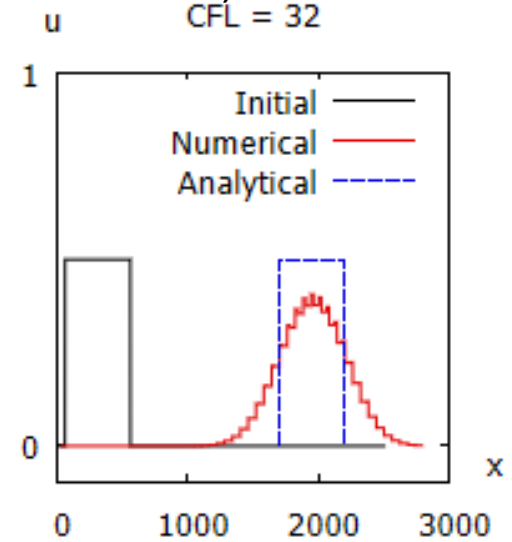
$D = 60, d = 60$   
CFL = 60



$D = 60, d = 55$   
CFL = 55



$D = 60, d = 32$   
CFL = 32





## 2. Our LTS solution (1D)

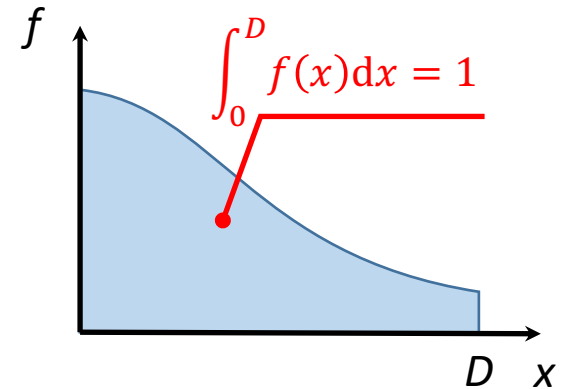
### 2.2 Convolution in 1D

The method performs poorly when  $D \neq d$  even for a small (relative) difference between the two

$$\frac{1}{d} \int_{x_{i+1/2}-d}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx \approx \frac{1}{D} \int_{x_{i+1/2}-D}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx$$



Approximation #4 (our contribution 2): give less weight to those cells that are farther away from the interface (weighted averaging/convolution)



$$\frac{1}{d} \int_{x_{i+1/2}-d}^{x_{i+1/2}} \mathbf{u}(x, t_n) dx \approx \int_{x_{i+1/2}-D}^{x_{i+1/2}} f(x_{i+1/2} - x) \mathbf{u}(x, t_n) dx = f * \mathbf{u}$$

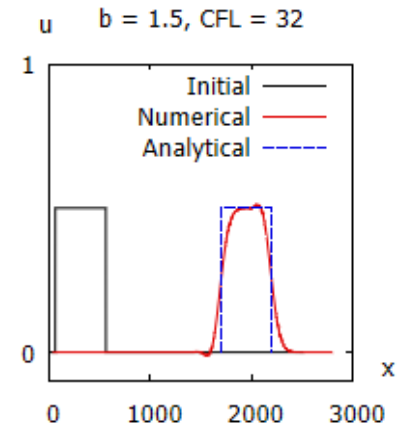
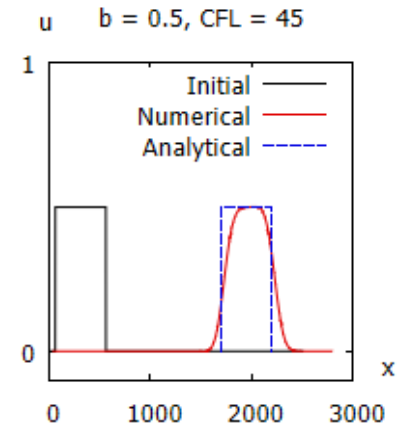
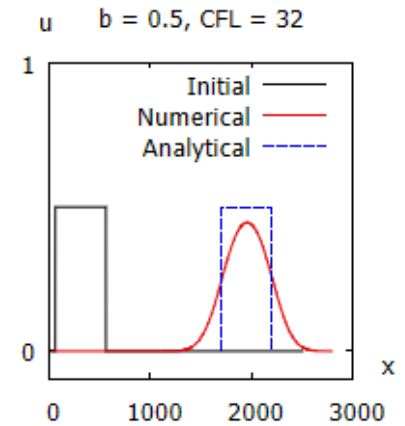


## 2. Our LTS solution (1D)

### 2.2 Convolution in 1D

Results for a power law convolution kernel  
( $D = 60$ )

$$f(x) = a \left(\frac{x}{D}\right)^b \qquad a = \frac{b + 1}{D}$$

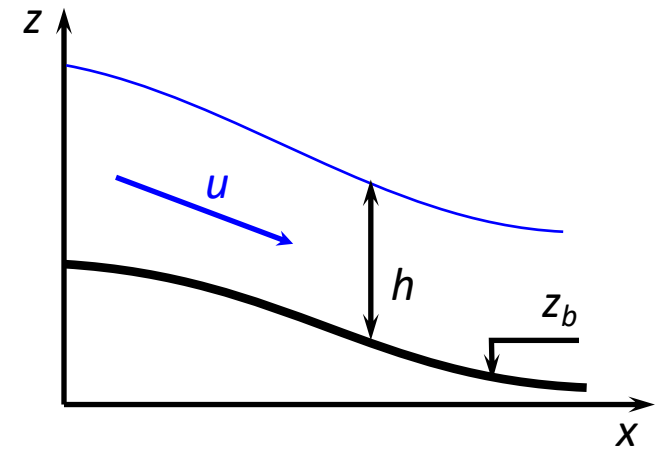


## 2.3 Shallow water equations &amp; bottom source term

$$\partial_t \mathbf{u} + \partial_x \mathbf{f} = \mathbf{s} \quad \mathbf{u} = \begin{bmatrix} h \\ hu \end{bmatrix}$$

$$\partial_t \mathbf{u} + \mathbf{A} \partial_x \mathbf{u} = \mathbf{s} \quad \mathbf{f} = \begin{bmatrix} hu \\ hu^2 + gh^2/2 \end{bmatrix}$$

$$\mathbf{s} = \begin{bmatrix} 0 \\ (S_0 - S_f)gh \end{bmatrix}$$



EVR approach : reconstructing  $z$  alone is sufficient

$$\mathbf{u} = \frac{h}{2} \mathbf{k}_1 + \frac{h}{2} \mathbf{k}_2 = \frac{h}{2} \begin{bmatrix} 1 \\ u - c \end{bmatrix} + \frac{h}{2} \begin{bmatrix} 1 \\ u + c \end{bmatrix}$$

$$\bar{\mathbf{u}} \approx \frac{\bar{h}}{2} \mathbf{k}_1 + \frac{\bar{h}}{2} \mathbf{k}_2$$

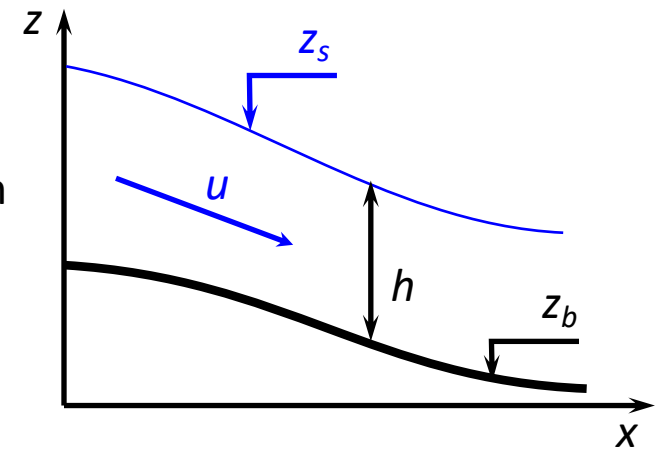
## 2.3 Shallow water equations &amp; bottom source term

We proved that static conditions are preserved provided that

- Averaging is done for the free surface elevation  $z_s$  instead of the depth  $h$

$$\partial_t \mathbf{u}' + \partial_x \mathbf{f} = \mathbf{s} \quad \mathbf{u}' = \begin{bmatrix} z_s \\ hu \end{bmatrix}$$

- The bottom elevation  $z_b$  is taken from the local values over the cells immediately next to the interface



## 2. Our LTS solution (1D)

### 2.4 Stability

For the linear advection equation

$$u_i^{n+1} = u_i^n + \nu \left( \tilde{u}_{i-1/2,L}^n - \tilde{u}_{i+1/2,L}^n \right)$$

$$\nu = c \frac{\Delta t}{\Delta x}$$

$$\tilde{u}_{i+1/2,L}^n = w_L * u(x, t_n) \approx \sum_{k=1}^K f_k u_{i-k+1}^n \Delta x$$

$$\sum_{k=1}^K f_k \Delta x = 1$$

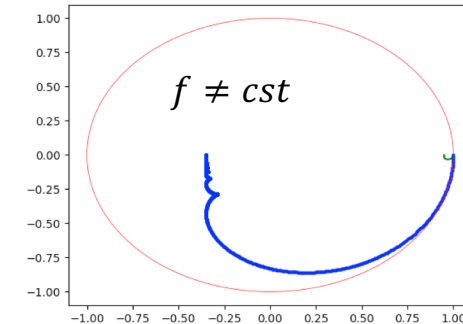
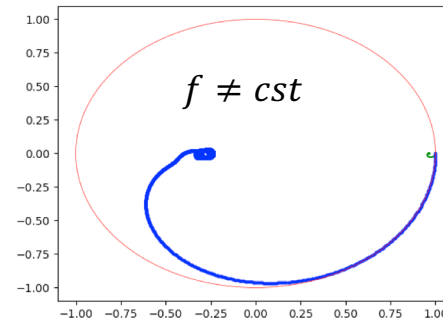
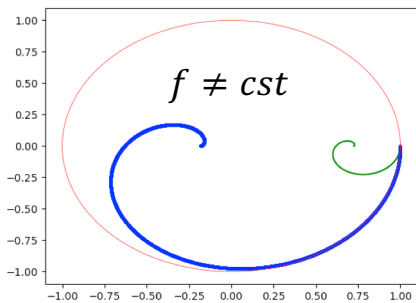
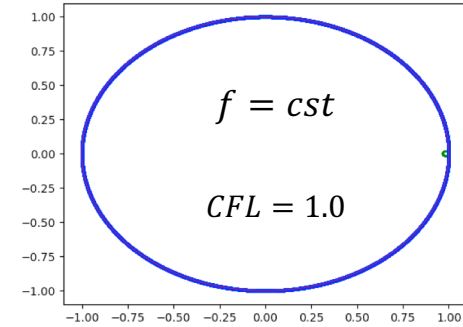
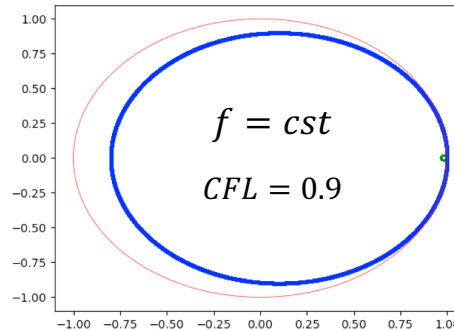
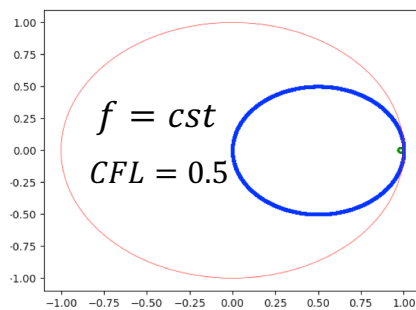
Von Neumann analysis:

$$u_i^n = u_0^0 \exp(\mu t + j\sigma i \Delta x)$$

Need to prove that this amplification factor remains « in the unity disk »

$$A_K \equiv \frac{u_i^{n+1}}{u_i^n} = 1 + \nu (1 - e^{j\sigma \Delta x}) \sum_{k=1}^K f_k e^{-jk\sigma \Delta x} \Delta x$$

## 2.4 Stability



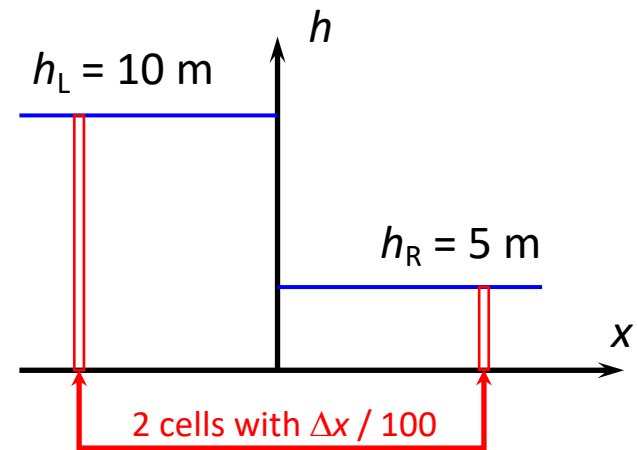
**Theorem 1.** For every  $D > 0$  and every decreasing kernel function  $f$  on  $[0, D]$ ,

$$\exists \nu > 0 : \forall \sigma \geq 0, |A_\infty(\sigma)| \leq 1$$

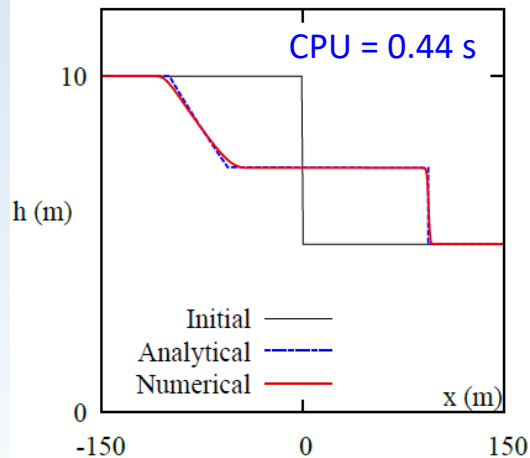
The same result *should* hold for  $A_K(\sigma)$ ...

Application to the shallow water equations

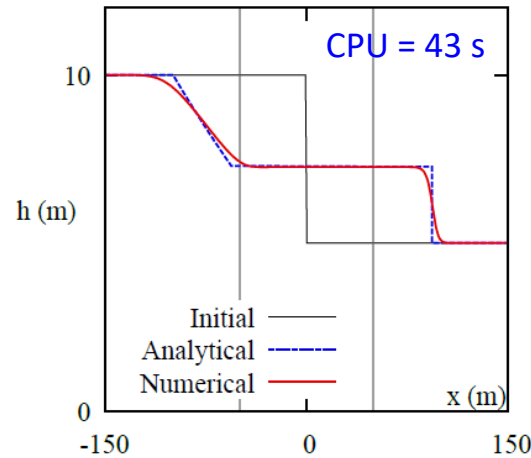
- Dambreak problem, flat bottom ( $\exists$  analytical solution)



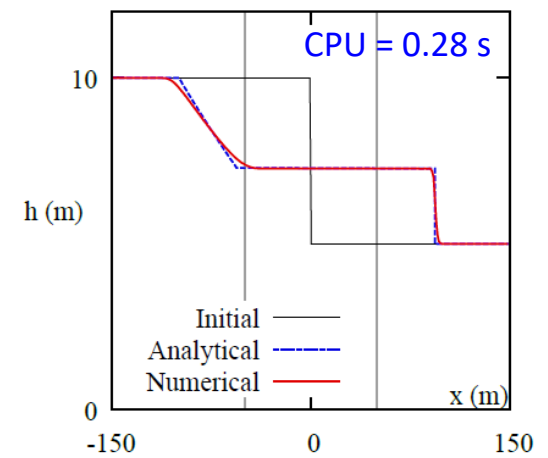
Godunov regular grid, max. CFL = 1.0



Godunov irregular grid, max. CFL = 1.0



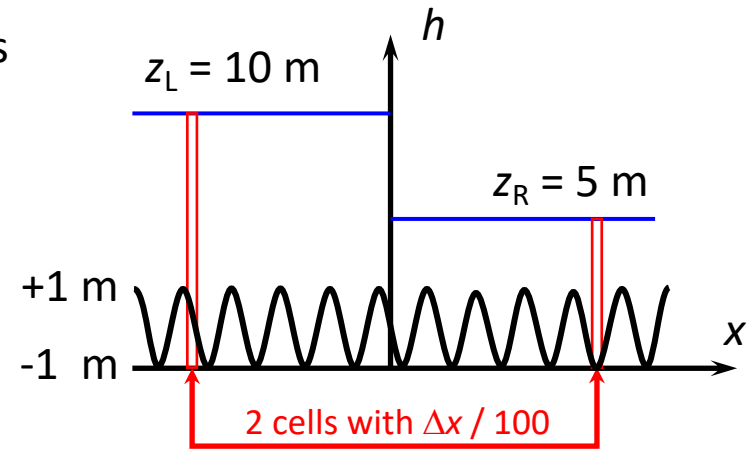
LTS irregular grid, max. CFL = 170



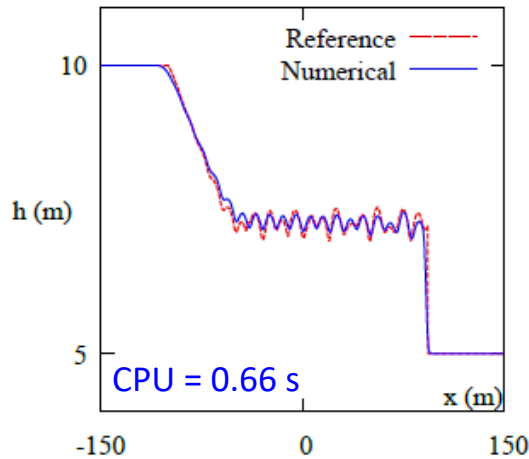
### 3. Computational examples

Application to the shallow water equations

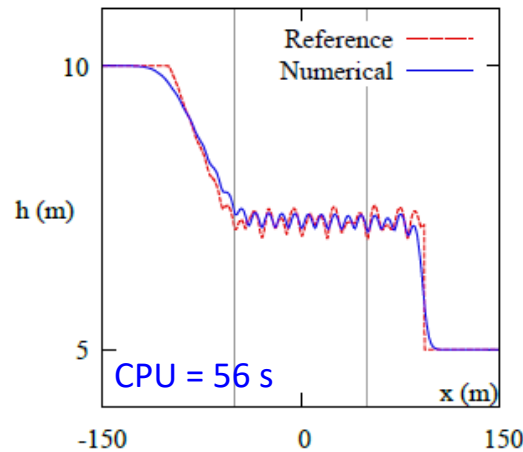
- Dambreak problem, sinusoidal bottom (∄ analytical solution)



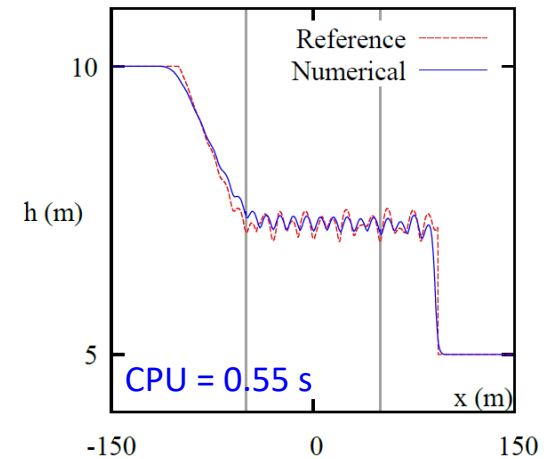
Godunov regular grid, max. CFL = 1.0



Godunov irregular grid, max. CFL = 1.0



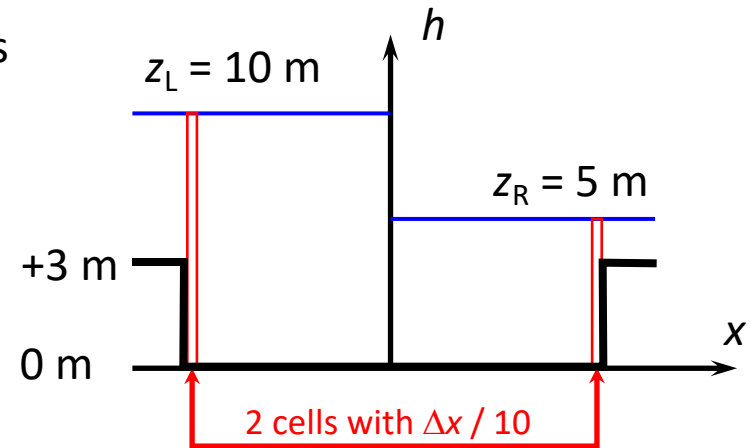
LTS irregular grid, max. CFL = 170



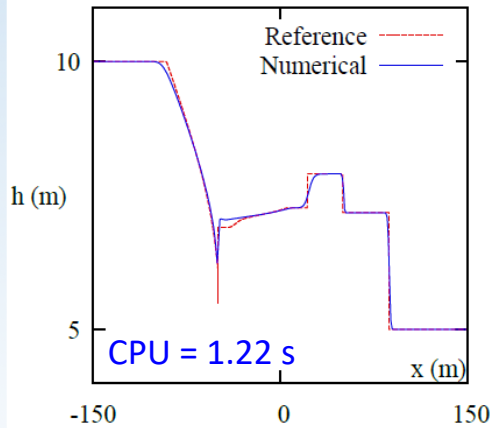


Application to the shallow water equations

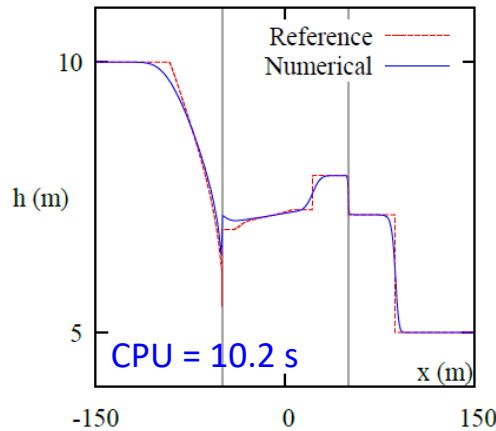
- Dambreak problem, discontinuous bottom (∄ analytical solution)



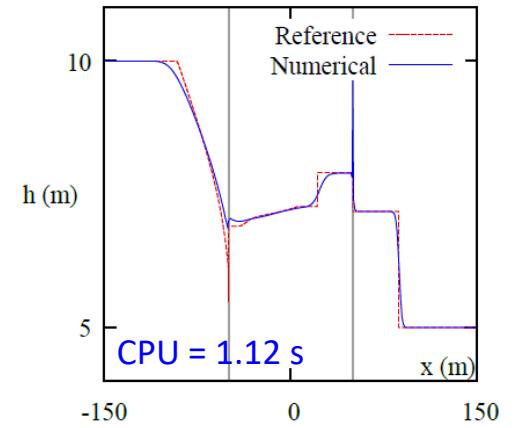
Godunov regular grid, max. CFL = 1.0



Godunov irregular grid, max. CFL = 1.0



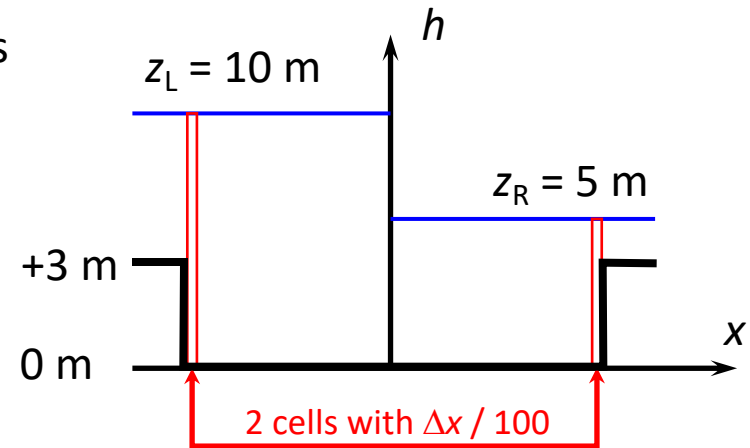
LTS irregular grid, max. CFL = 12



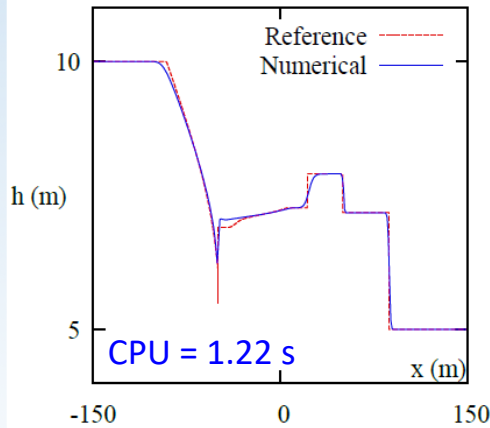
### 3. Computational examples

Application to the shallow water equations

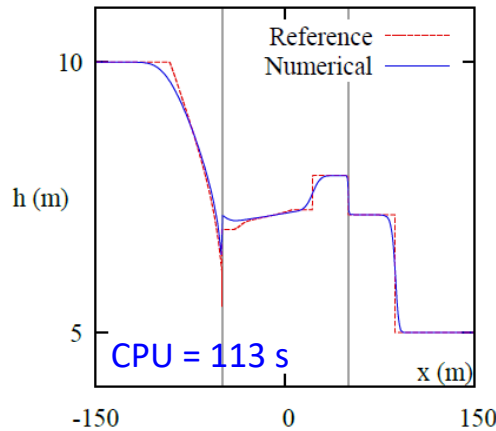
- Dambreak problem, discontinuous bottom (∄ analytical solution)



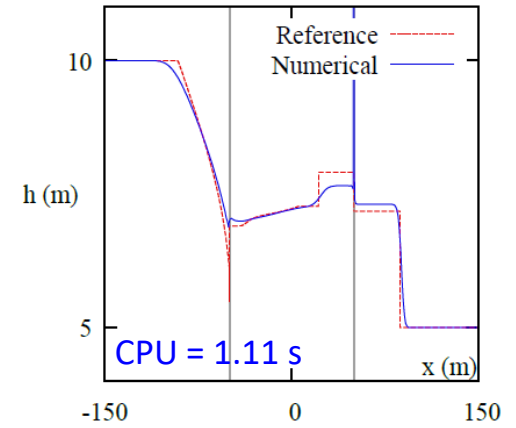
Godunov regular grid, max. CFL = 1.0



Godunov irregular grid, max. CFL = 1.0



LTS irregular grid, max. CFL = 120



So far

- No theoretical CFL restriction for homogeneous PDEs
- Smooth geometrical source terms (continuous topography) : no noticeable balancing issues/instability identified
- Stiff geometrical source terms (discontinuous topography): conservation issues as CFL increases (balancing errors are amplified), no instability identified
- Large  $I$  contrasts may require iterative estimate of  $D$

Pending questions / ongoing developments

- Treatment of boundaries
- How to handle dry/disconnected regions?
- Going 2D

