# HIERACHICAL METHODS FOR SHAPE OPTIMIZATION IN AERODYNAMICS Part II: Additive methods in Aerodynamics

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# 1 Introduction

The development of hierarchical methods for solving Optimum Shape design problems for CFD is motivated by efficiency improvement in the case where the set of parameter involves discretized functions.

Non-gradient methods need reduce the number of unknown.

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Gradient-based methods generally require preconditioners. Indeed, let us consider the application of gradient iteration in a space of functions. The functional analysis tells us  $u^{n+1} = u^n - \rho g^n$ that the continuous iterate is less regular than  $u^n$  when the gradient  $q^n$  involves k-th derivatives of  $u^n$ . In other words the iteration operator is unbounded, with a *regularity loss* equal to k. The continuous iteration then diverges. In the discrete case, this results in the amplification of high frequency error components, except for very small, mesh dependent, step lengths  $\rho$ . This kind of phenomenon is the rule when iterative solution algorithm are applied to discrete PDE systems. Then a solution can be found in preconditioning or multi-levelling. Experience shows that this problem is very frequent in Optimal Shape Design when the iterations are applied in order to find a function describing the shape ("CAD-free" approach) and not only a few shape parameters. In this paper, we propose an explanation of this, related to the lack of smoothness of the Hadamard derivative of a function with respect to the domain shape. As a practical consequence, the discrete gradients to be used for updating shapes are oscillatory unless a very small step length is adopted. To cure this, Reuther and Jameson have proposed in (20) a method of correction smoothing that solves a Laplace-Beltrami system set on the boundary. This method is also used in (16). In (1), Arian and Ta'asan analyse the Fourier symbol of the Hadamard-type functional Hessian. According to the degree of this Hessian, these authors propose to use either a Laplace-Beltrami operator when the regularity loss is 2 or a Neumann-Dirichlet pseudo-differential operator when regularity loss is 1.

In this chapter, we examine how to introduce multi-level principles in shape design algorithms in order to reduce their mesh-dependancy or their number of parameter dependancy of convergence speed.

Two main directions will be presented:

- to extend multi-level principle to CAD parametrization. We define multi-level Beziers parametrization.

- to extend multi-level preconditioning to CAD-free parametrization following the Bramble-Pasciak-Xu (BPX) theory (see (4; 25; 6)).

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# 2 Hadamard formula and functional preconditioning

#### 2.1 The Dirichlet case

Large scale problems coming from Partial Differential Equations generally result in a poor conditioning that further degrades when the number of degrees of freedom is increased. To explain and solve this problem, we can either analyse directly the behavior of discrete eigenvalues when the number of unknowns is increased or analyse the continuous -functional- problem and the continuous version of the algorithm. We concentrate on the second way.

We are interested in minimizing a functional  $j(\gamma)$  with respect to a "shape parameter"  $\gamma$ . More precisely, an initial geometry  $\Omega_0$  is equipped with a vector field  $\vec{n}_0$  normal to its boundary. A family of domains  $\Omega_{\gamma}$  of  $R^d$  is parameterized by a displacement  $\gamma \in C^{l+\alpha}(\partial \Omega_0)$  of the boundary in direction  $\vec{n}_0$ .

The state equation is the Poisson problem:

$$-div \ grad \ z(\gamma) = f \ ; \ z(\gamma) = 0 \ on \ \partial\Omega_{\gamma}$$

Let D be a subdomain of  $\Omega_{\gamma}$  (inside  $\Omega_{\gamma}$  for any  $\gamma$  admissible). The functional j to minimize is defined by:

$$j(\gamma) = \frac{1}{2} ||z(\gamma) - z_{target}||_D^2$$

We are interested in computing the derivative of the functional with respect to the shape parameter. We have to do it by a chain rule involving the state equation. A classical difficulty comes from the variable domain in the state equation. Early solution to this variation calculus dates back to Hadamard. We shall instead invoke some techniques initated by the interior variation of Garabedian, see (10), and extended later on in (17),(19), (?). (8),(9).

We introduce first a family of diffeomorphisms  $(T_{\gamma})_{\gamma}$  smooth with respect to  $\gamma$  such that:

$$\begin{array}{ccccc} T_{\gamma} & maps & \mathcal{O} & on & \mathcal{O}, \\ T_{\gamma} & maps & \Omega_0 & on & \Omega_{\gamma}, \end{array}$$

The problem can be stated in terms of a new state variable:

$$j(\gamma) = \frac{1}{2} ||\tilde{z}(\gamma) \circ T_{\gamma}^{-1} - z_{target}||_{D}^{2}$$
$$\tilde{z}(\gamma) = z(\gamma) \circ T_{\gamma} .$$

The new state  $\tilde{z}(\gamma)$  is solution of a well-posed Dirichlet problem with coefficients and source depending on  $\gamma$  through  $T_{\gamma}$  and  $T_{\gamma}^{-1}$  but formulated in the domain D which does not vary with the control  $\gamma$ . Then  $\tilde{z}(\gamma)$  depends implicitly on  $\gamma$  and under regularity assumptions on f and  $\partial\Omega_0$ , we can prove by applying the implicit function theorem that mapping:

$$\gamma \to \tilde{z}(\gamma)$$

is continuously differentiable from  $\mathcal{C}^{l+\alpha}(\partial\Omega_0)$  in  $\mathcal{C}^{l-1+\alpha}(\Omega_0)$ . It then interesting to introduce a linear extension operator from  $\Omega_0$  to  $\mathcal{O}$ :

$$\mathcal{P}: \mathcal{C}^k(\Omega_0) \to \mathcal{C}^k(\mathcal{O})$$
$$f \mapsto \mathcal{P}f$$
$$\mathcal{P}f|_{\Omega_0} = f .$$

Let:

$$\bar{z}(\gamma) = (\mathcal{P}\tilde{z}(\gamma)) \circ T_{\gamma}^{-1}$$

then, by chain rule and with a possible reduction of regularity l, we deduce that

$$\gamma \to \bar{z}(\gamma) \tag{1}$$

is continuously differentiable from  $\mathcal{C}^{l+\alpha}(\partial\Omega_0)$  in  $\mathcal{C}^{l-1+\alpha}(\Omega_0)$ . Now, by construction,  $\bar{z}(\gamma)$  is an extension of  $z(\gamma)$ :

$$\bar{z}(\gamma)|_{\Omega_{\gamma}} = z(\gamma) . \tag{2}$$

And we can exhibit the derivative of  $\bar{z}(\gamma)$  with respect to  $\gamma$  on  $\Omega_{\gamma}$ , in which it is the solution of the non-homogeneous Dirichlet problem:

$$-divgrad \frac{\partial \bar{z}}{\partial \gamma}(\gamma).\delta \gamma = 0 \text{ in } \Omega_{\gamma}$$
$$\frac{\partial \bar{z}}{\partial \gamma}(\gamma).\delta \gamma = \frac{\partial \bar{z}(\gamma)}{\partial n_{\gamma}}(n_{\gamma}, n_{0})\delta \gamma$$
(3)

where  $\vec{n}_{\gamma}$  is the normal to  $\partial \Omega_{\gamma}$ . One interesting way to state the above result is to compute the *total* derivative of the following state variational residual, which is possible due to the derivability of an extension of  $z(\gamma)$ :

$$\begin{split} \Psi(\gamma, z; \phi_1, \phi_2) &= -\int_{\Omega_{\gamma}} (div \ grad \ z \ - \ f) \ \phi_1 \ dv \ + \ \int_{\partial \Omega_{\gamma}} z \ \phi_2 \ d\sigma \ , \\ \Psi(\gamma, \bar{z}(\gamma); \phi_1, \phi_2) &= \ 0 \ \forall \phi_1, \forall \phi_2 \ . \end{split}$$

Indeed the derivative with respect to  $\bar{z}$  produces the left-hand side of (), and after injection of the state equation, the derivative with respect to  $\gamma$  reduces to:

$$\int_{\partial\Omega_{\gamma}} \frac{\partial \bar{z}(\gamma)}{\partial n_{\gamma}} (n_{\gamma}, n_0) \delta\gamma \ \phi_2 \ d\sigma \ .\Box$$

By chain rule, j is also differentiable and its gradient is expressed as follows:

$$j'(\gamma).\delta\gamma = \int_{\partial\Omega_{\gamma}} \frac{\partial z(\gamma)}{\partial n_{\gamma}} \frac{\partial p(\gamma)}{\partial n_{\gamma}} < \vec{n}_{\gamma}, \vec{n}_{0} > \delta\gamma \quad d\partial\Omega_{\gamma}$$

where  $p(\gamma)$  the following adjoint state:

$$-div \ grad \ p(\gamma) = z_{target} - z(\gamma) \ ; \ p(\gamma) = 0 \ on \ \partial \Omega_{\gamma}$$

Taking the  $L^2$  space as pivot space for a *continuous* gradient method would produce the following iteration ( $\rho$  is a positive step length):

$$\gamma^* = \gamma - \rho g_{L^2}$$
 where  $g_{L^2} = \frac{\partial z(\gamma)}{\partial n_{\gamma}} \frac{\partial p(\gamma)}{\partial n_{\gamma}} < \vec{n}_{\gamma}, \vec{n}_0 > .$ 

We observe that, starting from a previous iterate  $\gamma$  belonging to  $C^{l+\alpha}(\partial\Omega_0)$ , we get a corrected  $\gamma^*$  that is only of regularity  $C^{l-1+\alpha}(\partial\Omega_0)$ . One degree of regularity is lost at each iteration, and iteration cannot continue after a few steps. We need precondition the iteration:

$$\gamma^* = \gamma - \rho B g_{L^2}$$
.

where the self-adjoint invertible operator B is chosen in order to recover the degree of regularity lost by the  $L^2$  gradient. It is a functional preconditioner similar to the one introduced in the Least Square formulation of (5).

The main motivation for doing this is that the continuous algorithm has a convergence rate that evidently does not depend on a mesh size. Then we can try to build some consistent discretization that will hopefully converge with a rate not so different from the continuous rate. This means that *essentially mesh-independent rates* might be obtained. Another way to understand that point is to remember that in the linear periodic case, Fourier's analysis shows that operators involving *p*-th order spatial derivatives will have eigenvalues of the order of  $\left(\frac{1}{\Delta x}\right)^p$ . The smaller this degree, the better the condition number. The strategy which we propose is to build the operator *B* in such a way that **the resulting order of spatial differentiation after preconditioning be equal to zero**.

#### 2.2 The Euler case

It is probably a generic property of shape first derivatives (formulated in terms of shape location) that the shape gradient of the state variable have a derivative loss of 1. The above analysis shows it rigourously for the Dirichlet problem. We consider now a more complex model: the flow around an obstacle of boundary  $\partial \Omega_{\gamma}$  is considered. We denote by  $W_1, W_2, W_3, W_4, W_5, p$  the conserved variables (density, three moments, energy) and the pressure of the flow of a perfect gas. The state equation is now the system of steady Euler equation that we write under a variational form as follows: for all  $\phi = (\phi_1, \phi_2, \phi_3, \phi_4, \phi_5)$ belonging to the appropriate space,

$$(\Psi(\gamma, W), \phi) = -\int_{\Omega_{\gamma}} (F(W)\partial\phi x + G(W)\partial\phi y + H(W)\partial\phi z) \, d\Omega_{\gamma} + \int_{\partial\Omega_{\gamma}} p \left(n_x^{\gamma} \phi_2 + n_y^{\gamma} \phi_3 + n_z^{\gamma} \phi_4\right) \, d\partial\Omega_{\gamma} = 0, \qquad (4)$$

where F(W), G(W) and H(W) hold for the usual Euler fluxes, corresponding respectively to each of the space directions.

The cost functional is denoted in a generic manner as follows:

$$j(\gamma) = J(\gamma, W(\gamma)) .$$
(5)

where  $W(\gamma)$  is the solution of state system for the parameter  $\gamma$ . On the way of a rigorous Hadamard differentiation with respect to shape, an invertibility of Jacobian is missing for the Euler solution and the implicit function theorem is not justified. We follow the lines of (3). The adjoint state  $\Pi$  is solution of:

$$\left(\frac{\partial F}{\partial W}\right)^* \frac{\partial \Pi}{\partial x} + \left(\frac{\partial G}{\partial W}\right)^* \frac{\partial \Pi}{\partial y} + \left(\frac{\partial H}{\partial W}\right)^* \frac{\partial \Pi}{\partial z} = -\frac{\partial j}{\partial W}.$$

in combination with boundary conditions, among which we mention only the boundary condition on the shape boundary controlled by the parameter  $\gamma$ :

$$\Pi_2 n_x^{\gamma} + \Pi_3 n_y^{\gamma} + \Pi_4 n_z^{\gamma} = 0 \quad \text{on} \quad \partial \Omega_{\gamma}$$

Similarly to the previous section, the gradient of cost functional is a product of adjoint state by the -formal, this time- derivative of state residual in which state equation has been re-injected. After simplifications, we get:

$$g_{L^{2}}(\gamma, W, \Pi) = - (F(W)\frac{\partial\Pi}{\partial x} + G(W)\frac{\partial\Pi}{\partial y} + H(W)\frac{\partial\Pi}{\partial z}) (\vec{n}^{\gamma} \cdot \vec{V}) + (\nabla p \Pi + p \nabla\Pi) (\vec{n}^{\gamma} \cdot \vec{V}) .$$
(6)

The gradient correction then writes:

$$\gamma = \gamma - \rho g_{L^2}(\gamma, W, \Pi)$$
.

We observe that this correction is generally less regular than the boundary parameter  $\gamma$ . Indeed, inspired by elliptic smoothness, we can estimate that the state variables are at most as regular as the boundary, but the above correction involves first derivatives of the adjoint state. Additionally, the normal vector  $\vec{n}^{\gamma}$  is a first derivative of the boundary parametrization  $\gamma$ . Therefore the formal loss of regularity is again 1.

### 3 Additive multilevel preconditioners

Additive multilevel preconditioners have been initially derived in a discrete context for solving elliptic Partial Differential systems which are typically of second order or of even order. An extremely rich litterature exists on this topic. The hierarchical basis method was first analysed by Yserentant in his pioneering paper (26). The work of Yserentant was apparently motivated by the famous unpublished technical report of Bank and Dupont (2). A more complete theory was proposed by Bramble, Pasciak and Xu (4), (24). See also the wavelet extensions, for example in (6). An extended theory can be found in (12; 25). The purpose of this section is to recall some results of (25) (7) in a format adapted to our purpose. For simplicity, we state them in the case of Dirichlet boundary conditions.

Let  $\Omega$  be a regular subdomain of  $\mathbb{R}^n$ . Let  $H^1(\Omega)$  the usual Sobolev space and  $V = H_0^1(\Omega)$ its subspace of functions vanishing at boundary  $\partial \Omega$ . It is included in  $H = L^2(\Omega)$ , a pivot space the scalar product and norm of which are denoted by:

$$(u,v) = \int_{\Omega} u v dx; ||u|| = (u,u)^{1/2}$$

Let  $(V_k)_{k=0,1,\dots}$  be a sequence of discretisation subspaces of V :

$$V_0 \subset \cdots \subset V_k \subset \cdots \subset V$$

To fix the ideas, subspaces  $V_k$ 's can be considered as built from nested quasi uniform meshes with mesh size:

$$h_k \quad \tilde{=} \quad 2^{-k}$$

For any k we introduce the fine-to-coarse projection operator  $Q_k : V \longrightarrow V_k$  defined for all  $u \in V$  by

$$Q_{-1}u = 0$$
; and for  $k \ge 0$ ,  $(Q_k u, v) = (u, v) \quad \forall v \in V_k$ .

Let  $a \in [0, 3/2]$  and  $B_a$  defined by:

$$B_a = \sum_{k=0}^{+\infty} \left(\frac{1}{2^a}\right)^k \left(Q_k - Q_{k-1}\right).$$
 (7)

#### Some comments:

In the case a = 1,  $B_1$  is nothing but the identity operator on functions for which the series is convergent. Each term represent a particular scale and smoothing is performed by damping high-frequencies/small scales.  $\Box$ 

#### Theorem (7):

Let s be such that  $-3/2 \leq s \leq s+a \leq 3/2$ . Then the operator B is bounded from  $H_0^s(\Omega)$  into  $H_0^{s+a}(\Omega)$ , in particular, there exists two positive numbers c and C such that for all  $u \in H_0^s(\Omega)$ , we have

$$c \|u\|_{H^{s}(\Omega)} \leq \|B_{a}u\|_{H^{s+a}(\Omega)} \leq C \|u\|_{H^{s}(\Omega)}.$$

This statement shows that the operator B has smoothing properties that can be quantified as a gain in regularity. This gain is exactly a. It can be prescribed by the user in order to precondition an operator having a as regularity loss. We introduce in the next sections the agglomeration-based construction of the  $V_k$  spaces.

### 4 Combination with the BFGS acceleration

Modern optimisation algorithms are equipped with quasi-Newton processes which allow a rather good efficiency for many applications. One of the most efficient quasi-Newton acceleration relies on the BFGS method, see for example (18). The purpose of this section is to demonstrate on a simplified example that even for a very favourable context, a good quasi-Newton optimizer needs the complement of a preconditioner as far as the number of discretization unknowns is large enough.

We consider the following minimisation problem in 1D:

$$u_{opt} = \operatorname{ArgMin} \frac{1}{2} \int_0^1 (|\frac{\partial u}{\partial x}|^2 - fu) \, dv \; ; \; u(0) = u(1) = 0 \; .$$
 (8)

We can precondition this problem with the 1D version of our agglomeration preconditioner. The Hessian of this functional is a second-order differential operator, i.e. derivative loss is 2.

This problem can be solved by a conjugate gradient and in that case we just need, for preconditioning into that algorithm, to multiply the gradient by  $B_2$  defined as in as in (20) (with a = 2).

The use of BGFS is a little more complex. That method builds progressively an approximate Hessian  $\mathcal{M}$  for the functional to minimise, together with its inverse  $\mathcal{W}$ . This is done by *regula falsi* heuristics relying on the knowledge of successive values  $u_{k-1}, u_k$  of the control variable and the corresponding values of the functional gradient  $g_{k-1}, g_k$ . In order to introduce our preconditioner, we define:

$$\hat{u} = B_1 u \text{ and } \hat{g} = B_{-1} g \tag{9}$$

These variables allows the BFGS construction of  $\hat{\mathcal{W}}$  and  $\hat{\mathcal{M}}$ . We can then derive and inverse the matrix to apply to the gradient:

$$\mathcal{W} = B_{-1}\hat{\mathcal{W}}B_{-1} \quad \mathcal{M} = B_1\hat{\mathcal{M}}B_1 \tag{10}$$

We stress that, in contrast to the conjugate gradient, in order to precondition BFGS, we need to handle  $B_1$  and  $B_{-1}$  for the assembly of the BFGS matrix and of its inverse. This is easily done with the proposed multilevel preconditioner: we just choose respectively a = 1, a = -1.

In Fig.1, four experiments are depicted. The number m of unknowns in the discretisation of (8) is first set to 32. As rather classically, the non-preconditioned BFGS algorithm solves exactly the optimum in about 5 + m/2 = 21 iterations. The preconditioned version has a convergence which is more progressive, but not faster. For a larger number of unknowns, m = 256, the unpreconditioned BGFS shows a very slow convergence during at least the 100 first iterations. In large scale shape design, this convergence is not acceptable. In contrast, the preconditioned BFGS convergence is about the same as for the coarse-mesh one (21 iterations).

# 5 Preconditioning by node agglomeration

The usual BPX preconditioner essentially needs embedded meshes, that are not compatible with today's engineering applications, since these applications rely often on unstructured fine meshes. The adaptation of BPX to arbitrary fine meshes was proposed in (15) for equations. It can be in a rather straightforward manner adapted to optimum problems. We present now the main lines of it. The first section describes the case of a 2D or 3D mesh. The second section addresses the particular features of an unknown defined on a non-plane surface discretised by triangles in 3D.

#### 5.1 Multidimensional Agglomeration

We start from a fine triangulation or tetrahedrization. Nodes are located at vertices. The discrete fine level is the subspace of linear combinations of P1 shape functions:  $E_h = span\{\varphi_i, i = 1 \cdots n_h\}.$ 

A dual finite-volume cell  $C_i$  is defined around each vertex *i* by splitting each neighboring element with median plans and keeping subelements containing vertex *i*. The volume of  $C_i$  is denoted  $\mathcal{M}eas$  (*i*).  $E_h$  is equipped by the following weighted scalar product:

$$\forall u_h \text{ and } v_h \in E_h \quad (u_h, v_h)_h = \sum_i^{n_h} (u_h)_i (v_h)_i \mathcal{M}eas (i)$$
(11)

The agglomeration process relies on a partition of the set  $I^f = \{1, \dots, i, \dots, n_h\}$  of fine indices i.

$$I^{f} = I_{1} \cup \dots \cup I_{J} \cup \dots \cup I_{n_{2h}} \quad (n_{2h} \ll n_{h})$$

$$(12)$$

where any  $I_J$  involves the indices of a few neighbouring nodes. An algorithm for building such a partition can be found in (13). For any  $I_J$ , a coarser basis function is defined by:  $\Phi_J = \sum_{i \in I_J} \varphi_i$  and the coarser space is given by:

$$E_{2h} = Span\{\Phi_J, J = 1 \cdots n_{2h}\}$$

The linear prolongation operator,  $\overline{P}$ , from  $E_{2h}$  to  $E_h$  is defined by

$$\forall u_{2h} \in E_{2h} \quad \overline{P}u_{2h} = u_{2h} \in E_h .$$
(13)

Its adjoint  $\overline{P}^*$  is the restriction operator from  $E_h$  to  $E_{2h}$ , and it is defined via  $\overline{P}$  as its adjoint with respect to scalar product ??). by :

$$(\overline{P}^* u_h)_J = \frac{\sum_{j_m \in J}^{n_h} (u_h)_{j_m} \mathcal{M}eas(j_m)}{\mathcal{M}eas(J)}$$
(14)

where J is a coarse cell,  $j_m$  are fine cells included in J.  $\mathcal{M}eas(J)$  represents the measure of coarse cell J. These operators where also introduced in (14) for a different multi-level approach, of multiplicative type. In that paper, it is noted that, in contrast with the usual Galerkin nested finite element sequence, the sequence of spaces generated with the transfers  $\bar{P}$  and  $\bar{P}^*$  defined in (13) and (14) does not enjoy enough regularity according Sec.3 theory. This is related to the fact that the orders of accuracy of transfers have a sum equal to 1 + 1 = 2, not strictly larger than 2, see (11).

We define define **smoother transfers**. The necessary smoothness is recovered by combining the above transfer operator with a "smoothing operator" L, and with  $L^*$ , its adjoint with respect to the discrete  $L^2$  scalar product. This smoothing operator is an average between a node and its neighbors:

$$(Lu)_{i} = (1 - \theta)u_{i} + \theta \frac{\sum_{j \in \mathcal{N}(i) \cup \{i\}} \mathcal{M}eas(j) \ u_{j}}{\sum_{j \in \mathcal{N}(i) \cup \{i\}} \mathcal{M}eas(j)}$$
(15)

where  $\mathcal{N}(i)$  represents the set of neighbors of cell *i* and  $\theta$  is a parameter defined later on. The adjoint  $L^*$  of *L* is defined by:

$$(L^*u)_i = (1-\theta)u_i + \theta \sum_{j \in \mathcal{N}(i) \cup \{i\}} \frac{u_j \ \mathcal{M}eas(j)}{\sum_{k \in \mathcal{N}(j) \cup \{j\}} \mathcal{M}eas(k)} \ .$$
(16)

The coarse level  $V_k$  is built from  $V_{k+1}$  as follows:

$$V_k = L_k \bar{P}_k (V_{k+1}) \tag{17}$$

Starting from the initial fine space that we denote  $V_N$ , we define the analog of the BPX projection operator (3) to the directly coarser level as follows:

$$Q_N = L_N \bar{P}_N \bar{P}_N^* L_N^* \tag{18}$$

and for any coarser level:

$$Q_k = \prod_{i=k+1}^{i=N} L_i \ \bar{P}_i \ \prod_{j=N}^{j=k+1} \ \bar{P}_j^* \ L_j^*.$$
(19)

From this we get:

$$B_a = \sum_{k=0}^{N} \left( \left( \frac{1}{2^a} \right)^k - \left( \frac{1}{2^a} \right)^{k-1} \right) Q_k .$$

$$(20)$$

In practice, for any x,  $B_a x$  is computed within a unique cycle from fine to coarse and back.

#### 5.2 Agglomeration for a surface in 3D

We return to the notations of Sec. 2.2.

Let  $\Sigma_0$  be the initial 3D surface, made of triangles. The generic discrete surface  $\Sigma_{\gamma}$ , is defined by the translation of length  $\gamma$ , of the vertices of  $\Sigma_0$  along an approximate unit normal vector  $\vec{n}$  to  $\Sigma_0$  defined at vertices.

$$\vec{x}_i^{\gamma}$$
 is a vertex of  $\Sigma_{\gamma} \iff \vec{x}_i^{\gamma} = \vec{x}_i^o + \gamma(i) \vec{n}$  (21)

where *i* is the index of the vertex,  $\vec{x}_i^o$  is the physical position of the vertex of  $\Sigma_0$  with same index *i*. In order to precondition a correction on  $\gamma$ , we can construct a sequence of spaces and operators following the same process as in (12),(13), (14), but restricted to the surface and with the area of surfacic cells  $\mathcal{A}rea(j)$  instead of cell measures. In order to adapt our operators to irregular surfaces, the smoothing operator *L*, is now weighted by a scalar product of normals to the surface:

$$(L \ \vec{x})_i = (1 - \theta)\vec{x}_i + \theta \ \frac{\sum_{j \in \mathcal{N}(i) \cup \{i\}} w_{ij}\vec{x}_j}{\sum_{j \in \mathcal{N}(i) \cup \{i\}} w_{ij}}$$
(22)

where  $w_{ij}$  are the weights defined by :

$$w_{ij} = \max \left( \mathcal{A}rea(i)\vec{n}_i \cdot \mathcal{A}rea(j)\vec{n}_j , 0 \right) \quad \|\vec{n}_i\| = 1 \quad \forall \ i \ . \tag{23}$$

With this weighting, the smoothing is avoided on shape's dihedra (e.g. on wing trailing edges) and then allows changes in the angle. The rest of preconditioner definition is the same as in previous section. The smoothing parameter  $\theta$  is set equal to  $\frac{1}{2}$  according to the analysis of (14).

# 6 Application to an optimal shape design problem

We now study the adaptation of the proposed multi-level preconditioner to an optimal shape design loop.

#### 6.1 The aerodynamical shape design problem

We now make more precise the statement of the aerodynamical application introduces in Sec.2.2. In (21), the authors propose a very simplified model for measuring the "sonic boom downwards emission" (SBDE). It consists in evaluating the volume integral of the squared pressure gradient in an "observation box"  $\Omega^B$  (as shown in Figure 4) below the object. The cost functional is therefore the following:

$$j(\gamma) = J(\gamma, W(\gamma)) . \tag{24}$$

where  $W(\gamma)$  is the solution of state system (4) for the parameter  $\gamma$ , and

$$J(\gamma, W) = \alpha_1 (C_D(W) - C_D^t)^2 + \alpha_2 (C_L(W) - C_L^t)^2 + \alpha_3 \int_{\Omega^B} |\nabla \tilde{p}(W)|^2 dV.$$
(25)

Here  $\alpha_1, \alpha_2$  and  $\alpha_3$  are constants that prescribe the weights of three sub-criteria, related to lift, drag and sonic boom emission.  $C_L^t$  and  $C_D^t$  hold for target lift and drag. Lift coefficient  $C_L(W)$ , drag coefficient  $C_D(W)$ , and pressure field p(W) are computed from the design variable  $\gamma$  by solving the state equation (4) and obtaining the state variable W. The notation  $\tilde{p}(W)$  indicates that the pressure p(W) has been smoothed in order to integrate its gradient, even when shocks arise.

Practically, the observation box  $\Omega^B$  is a part of the computational domain placed below the airplane. Its upper boundary is a plane close below the aircraft.

#### 6.2 Discretized problem

The discrete CFD model uses an upwind Euler solver applying to unstructured tetrahedrizations. The shape is changed by moving the nodes on the boundary of the mesh along normals to that boundary. Their displacement is taken into account by a transpiration condition in order to avoid costly remeshings. A gradient of the discrete functional is computed with the help of an adjoint system, built using the Automatic Differentiation tool TAPENADE, (Hascoet and Pascual). Since this gradient is defined as a scalar field defined on the nodes of the boundary, we can apply to it the surfacic multilevel preconditioner introduced in Sec. 5.2.

### 6.3 Numerical evidence of best smoothing index

We consider the shape optimisation starting from an ONERA M6 wing. A coarse discretisation of this geometry, together with a pressure distribution on it is presented in Fig.5. After the necessary number of gradient iterations, the shape optimisation loop produces a new shape, that is quite different from the initial one, as illustrated in Fig.6. The resulting sonic boom reduction can be appreciated from the comparisons of the pressure variations under the wing for both geometries, which are depicted in Fig. 7. This can be obtained by various optimisation algorithms. The question is to obtain it in the most efficient way, and in particular to analyse the impact of our preconditioner.

First we want to show that the optimal parameter a = 1 predicted by the theory is numerically verified. We first study this with a coarse geometry, involving 2203 nodes for the 3D mesh, but yet 780 shape parameters. Two optimisation algorithms are applied, a gradient (Fig.8) and a conjugate gradient (Fig.9). Divergence of the Euler algorithm for distorted shapes did not allow a pertinent study of the BFGS option.

Convergence of the gradient iteration can be evaluated with the evolution of the gradient norm. The case without preconditioning is the case where a = 0. Then convergence is the slowest of the different options tested. Values like a = 0.5, a = 1.5, a = 1., a = 2. provide good speedups, in particular for the 10 first iterations. This point is important for the shape design loops which, in practice, are too computer consuming to allow for more optimisation iterations. But the theoretical value a = 1. appears as numerically optimal, for both cases of a pure gradient optimiser as well as a conjugate gradient one.

Let us verify that the efficiency of the preconditioner is good for finer discretisations. We consider a second mesh with 15463 nodes. The number of shape parameters is 3222. See Fig.10 and Fig.11.

The value a = 1. appears again as numerically optimal. In the case of the conjugate gradient, a good shape is obtained in about 10 iterations.

The case of a = 2. deserves some comments. Indeed, it is equivalent to a Laplace-Beltrami smoothing. All the above results show that that option is not bad, but is clearly not as good as the optimal one.

We complete the above convergence curves by Tables. 1 and 2 which compare the efficiency of the preconditioned algorithm to the standard one in the practical case where convergence of optimisation is not continued more than 10 iterations. If equivalent levels of convergence are asked from both algorithms, then an acceleration of a factor 5 is put in evidence.

### 6.4 Example fo applications

In (23) the proposed preconditioner is applied to more complex geometries. An example is given in Fig.12. Although the functional gradient depicted in Fig.12-middle is a smooth function, an important improvement of the convergence has also been experimented in applying the hierarchical preconditioner. The bottom of this figure shows the change in pressure field in a horizontal plane under the aircraft. This change is obtain in 15 iterations of global optimization. In a multi-discipline problem as shape optimization taking into account structure deformation at cruise, the CAD-free approach is extremely useful since the shape structural deformation and the shape design improvement are

described with the same data format. In that case, efficiency is much improved by the introduction of the above hierarchical preconditioner in the multi-disciplinary optimization loop. The interested reader is invited to examine the paper (22).

# 7 Conclusion

This paper presents a new preconditioning strategy for optimal shape design. An additive multilevel preconditioner is built from (a) the classical Bramble-Pasciak-Xu principle, and (b) the agglomeration principle. Functional analysis considerations show the central role of the loss of regularity in the gradient iteration. This loss can be compensated by a regularity gain carried by the preconditioner. This regularity gain is easily prescribed by the user at any real value.

This helps in particular to combine the preconditioner with a BFGS acceleration. The interest in combining both is demonstrated.

In the case of shape design, we exhibit a simplified example for which the loss of derivative can be rigorously evaluated and is equal to 1.

The final shape design application is a pre-industrial one already addressed by different methods. The numerical results confirm the *a priori* analysis of the regularity loss. They demonstrate that the proposed method improves notably the convergence of the shape design iteration.

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Case	Opt. Iter.	funct.	gradient
GC, N=780, M=11,000	10	0.02	0.3
GC preconditioned, N=780, M=11,000 $$	10	0.04	0.04
GC, N= $3222$ , M= $77,000$	10	0.06	0.3
GC preconditioned, N=3222, M=77,000 $$	10	0.03	0.1

 Table 1: Sonic boom optimisation: result of 10 iterations of CG

Case	Opt. Iter.	gradient	Speed up
GC, N=780, M=11,000	15	0.1	
GC preconditioned, N=780, M=11,000 $$	3	0.1	5
GC, N=3222, M=77,000	100	0.03	
GC preconditioned, N=3222, M=77,000 $$	20	0.04	5

 Table 2: Sonic boom optimisation: effort for dividing the cost by 8

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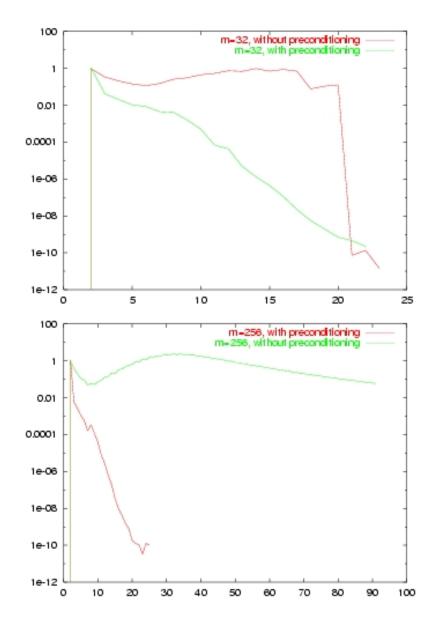


Figure 1: Resolution of a second order optimization problem with a preconditioned BFGS, behavior of the different residuals as functions of iteration number, for problem sizes m = 32 and m = 256.

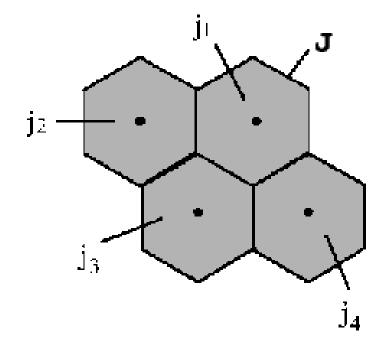


Figure 2: Sketch of the agglomeration of four fine (2D) cells,  $C_{j1}, C_{j2}, C_{j3}, C_{j4}$  into a coarser one  $\bar{C}_J = C_{j1} \cup C_{j2} \cup C_{j3} \cup C_{j4}$ 

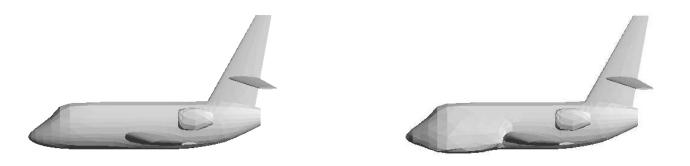


Figure 3: Hierarchical parametrization by agglomeration: second level basis function

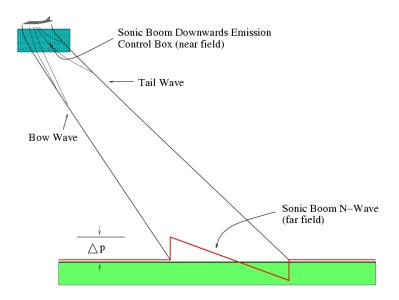


Figure 4: The sonic boom. Sketch of near and far field shock wave patterns of a supersonic aircraft. SBDE's control box  $\Omega^B$ .

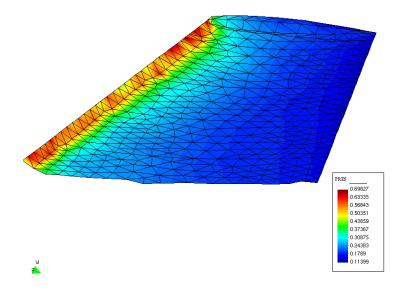


Figure 5: ONERA M6 shape optimisation: pressure contours on initial shape

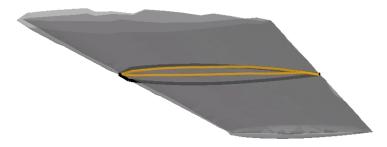


Figure 6: ONERA M6 shape optimisation: sketch of geometry, of initial and final shape.

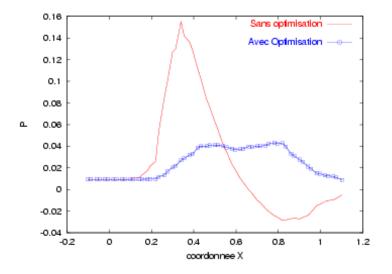


Figure 7: ONERA M6 shape optimisation: pressure signal under the wing, before, after optimisation

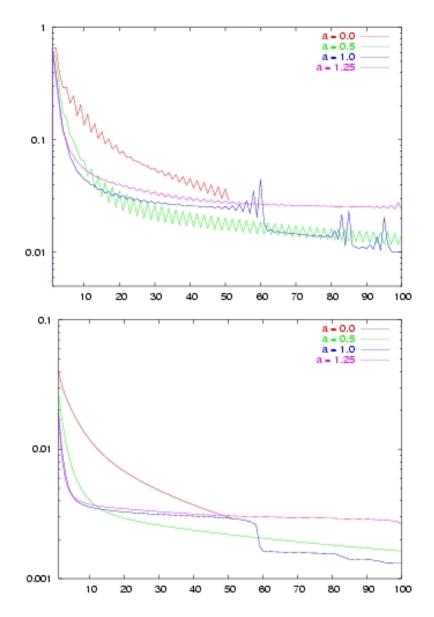


Figure 8: ONERA M6 shape optimization with a gradient method: Upper Gradient convergence, Lower Cost function,  $\mathrm{ns}=2203$ 

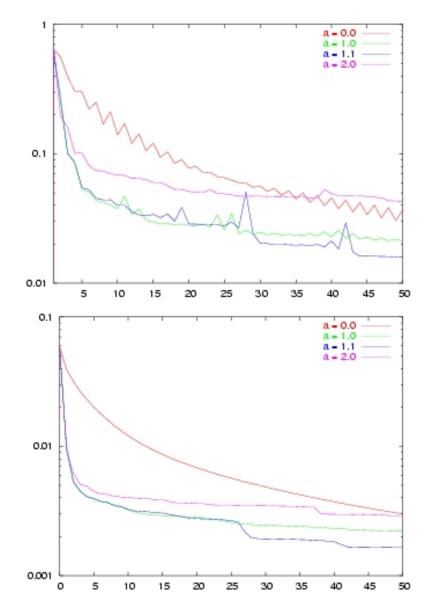


Figure 9: ONERA M6 shape optimization with a conjugate gradient method: Upper Gradient convergence, Lower Cost function, ns = 2203

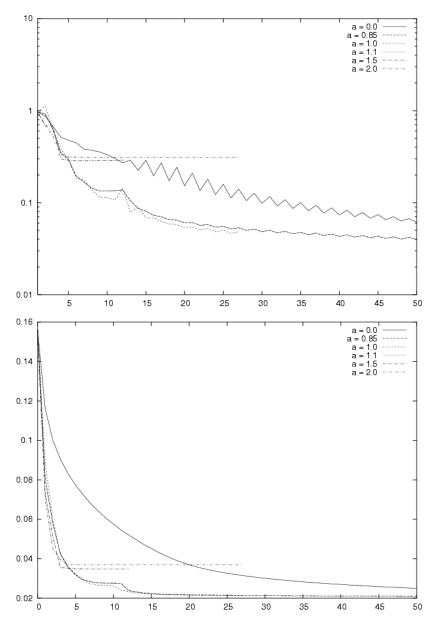


Figure 10: ONERA M6 shape optimization with a gradient method: Upper Gradient convergence, Lower Cost function, ns = 77315

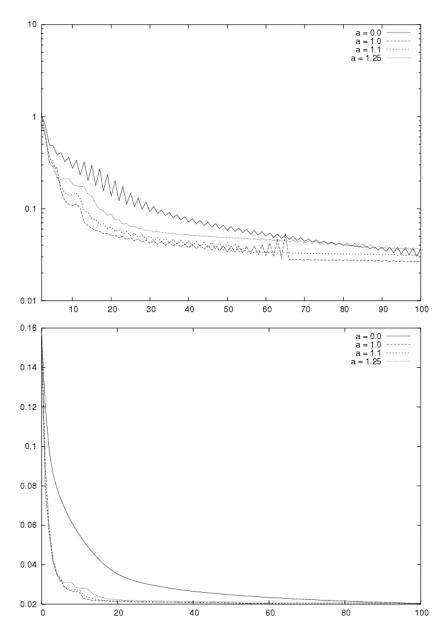


Figure 11: ONERA M6 shape optimization with a conjugate gradient method: Upper Gradient convergence, Lower Cost function, ns = 77315

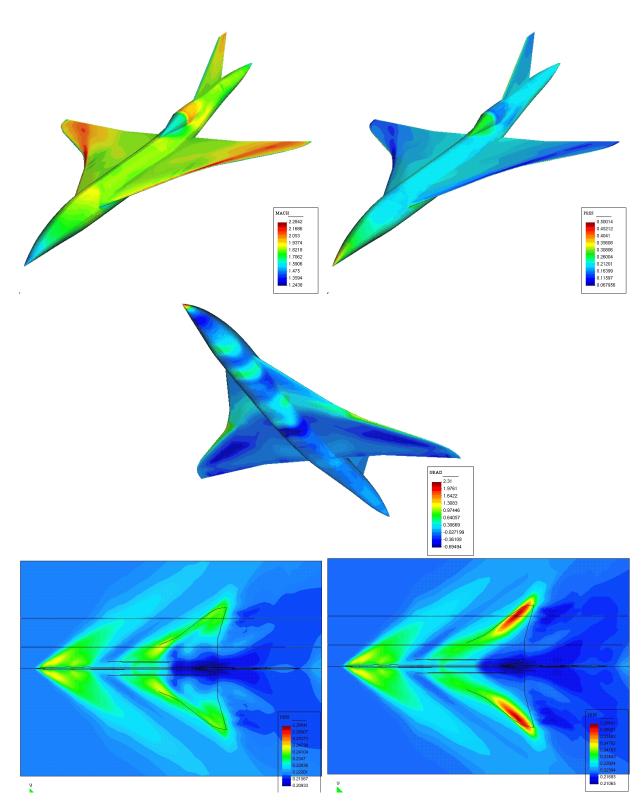


Figure 12: Shape optimization for reducing near field sonic boom: *Upper line:* Contour levels. Left, Mach number. Right, pressure. *Middle line:*Cost functional gradient distribution for the complete aircraft's surface. *Bottom line:*Pressure distribution in a plane below the aircraft. Left, original geometry. Right, optimized geometry.